

# Heavy hybrid mesons and tetraquarks from lattice QCD

## II – Tetraquarks

Effective Field Theory Seminar – Technische Universität München, Germany

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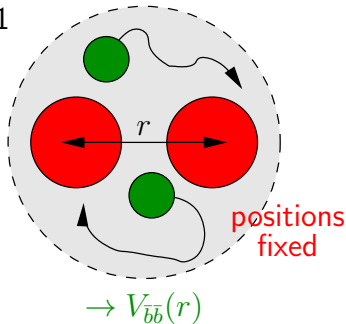
September 14, 2018



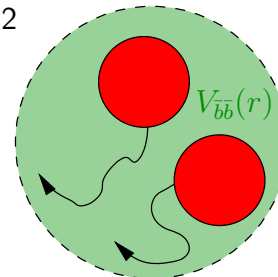
# Basic idea to study $bbqq$ tetraquarks (1)

- Study heavy tetraquarks  $\bar{b}\bar{b}qq$  or  $\bar{b}\bar{b}\bar{q}q$  in two steps.
    - (1) Compute potentials of two static quarks ( $\bar{b}\bar{b}$  or  $\bar{b}b$ ) in the presence of two lighter quarks ( $qq$  or  $\bar{q}q$ ,  $q \in \{u, d, s, c\}$ ) using lattice QCD.
    - (2) Explore, whether these potentials are sufficiently attractive to host bound states or resonances ( $\rightarrow$  tetraquarks) by using techniques from quantum mechanics and scattering theory.
- ((1) + (2)  $\rightarrow$  Born-Oppenheimer approximation).

step 1



step 2



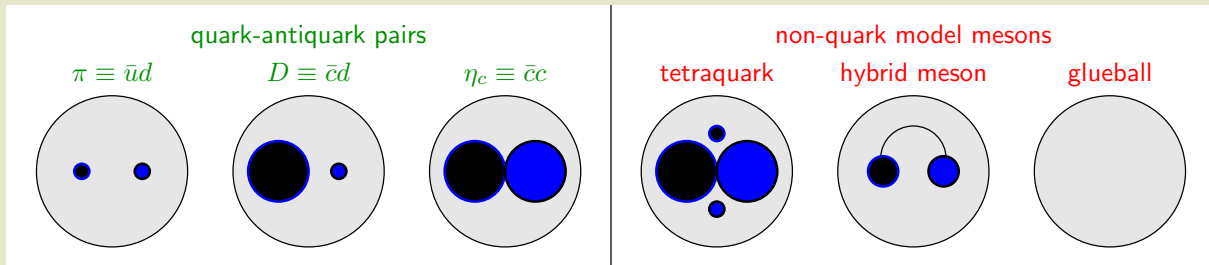
$\rightarrow$  existence of a tetraquark ... or not

# Basic idea to study $bbqq$ tetraquarks (2)

- The talk summarizes
  - [P. Bicudo, M.W., Phys. Rev. D **87**, 114511 (2013) [arXiv:1209.6274]]
  - [P. Bicudo, K. Cichy, A. Peters, B. Wagenbach, M.W., Phys. Rev. D **92**, 014507 (2015) [arXiv:1505.00613]]
  - [P. Bicudo, K. Cichy, A. Peters, M.W., Phys. Rev. D **93**, 034501 (2016) [arXiv:1510.03441]]
  - [P. Bicudo, J. Scheunert, M.W., Phys. Rev. D **95**, 034502 (2017) [arXiv:1612.02758]]
  - [P. Bicudo, M. Cardoso, A. Peters, M. Pflaumer, M.W., arXiv:1704.02383].
- For recent work from other groups using a similar approach cf. e.g.
  - [W. Detmold, K. Orginos, M. J. Savage, Phys. Rev. D **76**, 114503 (2007) [arXiv:hep-lat/0703009]]
  - [G. Bali, M. Hetzenegger, PoS **LATTICE2010**, 142 (2010) [arXiv:1011.0571 [hep-lat]]
  - [Z. S. Brown and K. Orginos, Phys. Rev. D **86**, 114506 (2012) [arXiv:1210.1953 [hep-lat]]

# Why are such studies important? (1)

- **Meson**: system of quarks and gluons with integer total angular momentum  $J = 0, 1, 2, \dots$
- Most mesons seem to be **quark-antiquark pairs**  $\bar{q}q$ , e.g.  $\pi \equiv \bar{u}d$ ,  $D \equiv \bar{c}d$ ,  $\eta_s \equiv \bar{c}c$  (quark-antiquark model calculations reproduce the majority of experimental results).
- Certain mesons are poorly understood (significant discrepancies between experimental results and quark model calculations), could have a more complicated structure, e.g.
  - **2 quarks and 2 antiquarks (tetraquark)**,
  - **a quark-antiquark pair and gluons (hybrid meson)**,
  - **only gluons (glueball)**.



# Why are such studies important? (2)

- Indications for tetraquark structures:

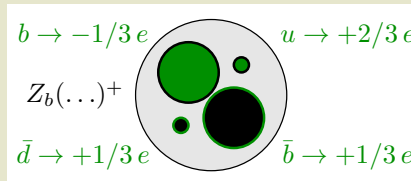
- Electrically charged mesons  $Z_b(10610)^+$  and  $Z_b(10650)^+$ :

- \* Mass suggests a  $b\bar{b}$  pair ...

- \* ... but  $b\bar{b}$  is electrically neutral ...?

- \* **Easy to understand, when assuming a tetraquark structure:**

$Z_b(\dots)^+ \equiv b\bar{b}u\bar{d}$  ( $u \rightarrow +2/3 e$ ,  $\bar{d} \rightarrow -1/3 e$ ).



- Electrically charged  $Z_c$  states:

- \* Similar to  $Z_b$ .

- Mass ordering of light scalar mesons:

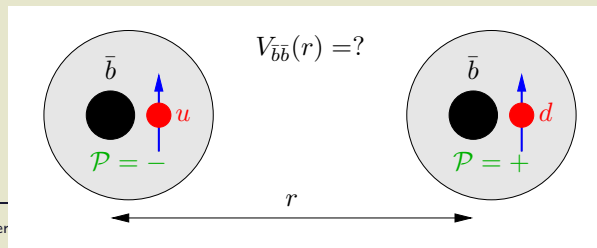
- \* E.g.  $m_{\kappa} > m_{a_0(980)}$  ...?

# Outline

- $\bar{b}bqq$  /  $BB$  potentials.
- Lattice setup.
- $\bar{b}bqq$  tetraquarks.
- Inclusion of heavy spin effects.
- $\bar{b}b\bar{q}q$  /  $\bar{B}B$  potentials

# $\bar{b}\bar{b}qq$ / $BB$ potentials (1)

- From now on  $\bar{b}\bar{b}qq$  ( $\bar{b}\bar{b}q\bar{q}$  technically more difficult, will be discussed at the end of this talk).
  - Spins of static antiquarks  $\bar{b}\bar{b}$  are irrelevant (they do not appear in the Hamiltonian).
  - At large  $\bar{b}\bar{b}$  separation  $r$ , the four quarks will form two static-light mesons  $\bar{b}q$  and  $\bar{b}q$ .
  - Consider only pseudoscalar/vector mesons ( $j^P = (1/2)^-$ , PDG:  $B, B^*$ ) and scalar/pseudovector mesons ( $j^P = (1/2)^+$ , PDG:  $B_0^*, B_1^*$ ), which are among the lightest static-light mesons ( $j$ : spin of the light degrees of freedom).
  - Compute and study the dependence of  $\bar{b}\bar{b}$  potentials in the presence of  $qq$  on
    - the “light” quark flavors  $q \in \{u, d, s, c\}$  (isospin, flavor),
    - the “light” quark spin (the static quark spin is irrelevant),
    - the type of the meson  $B, B^*$  and/or  $B_0^*, B_1^*$  (parity).
- Many different channels: attractive versus repulsive, different asymptotic values ...

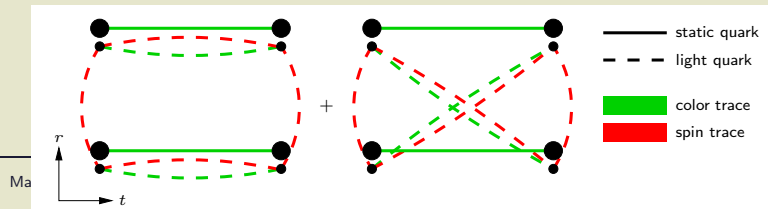
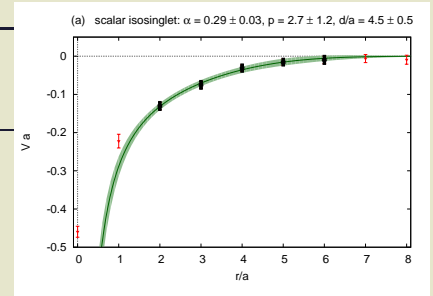


# $\bar{b}\bar{b}qq / BB$ potentials (2)

- Rotational symmetry broken by static quarks  $\bar{b}\bar{b}$ .
- Remaining symmetries and quantum numbers:
  - Rotations around the separation axis (e.g.  $z$  axis), quantum number  $j_z$ .
  - $P$ .
  - $P_x$  (reflection along an axis perpendicular to the separation axis, e.g.  $x$  axis).
- To extract the potential(s) of a given sector ( $I, I_z, |j_z|, P, P_x$ ), compute the temporal correlation function of the trial state

$$(C\Gamma)_{AB} (C\tilde{\Gamma})_{CD} \left( \bar{Q}_C(-\mathbf{r}/2) q_A^{(1)}(-\mathbf{r}/2) \right) \left( \bar{Q}_D(+\mathbf{r}/2) q_B^{(2)}(+\mathbf{r}/2) \right) |\Omega\rangle.$$

- $q^{(1)}q^{(2)} \in \{ud - du, uu, dd, ud + du, ss, cc\}$  (isospin  $I, I_z$ , flavor).
- $\Gamma$  is an arbitrary combination of  $\gamma$  matrices (spin  $|j_z|$ , parity  $P, P_x$ ).
- $\tilde{\Gamma} \in \{(1 - \gamma_0)\gamma_5, (1 - \gamma_0)\gamma_j\}$  (irrelevant).





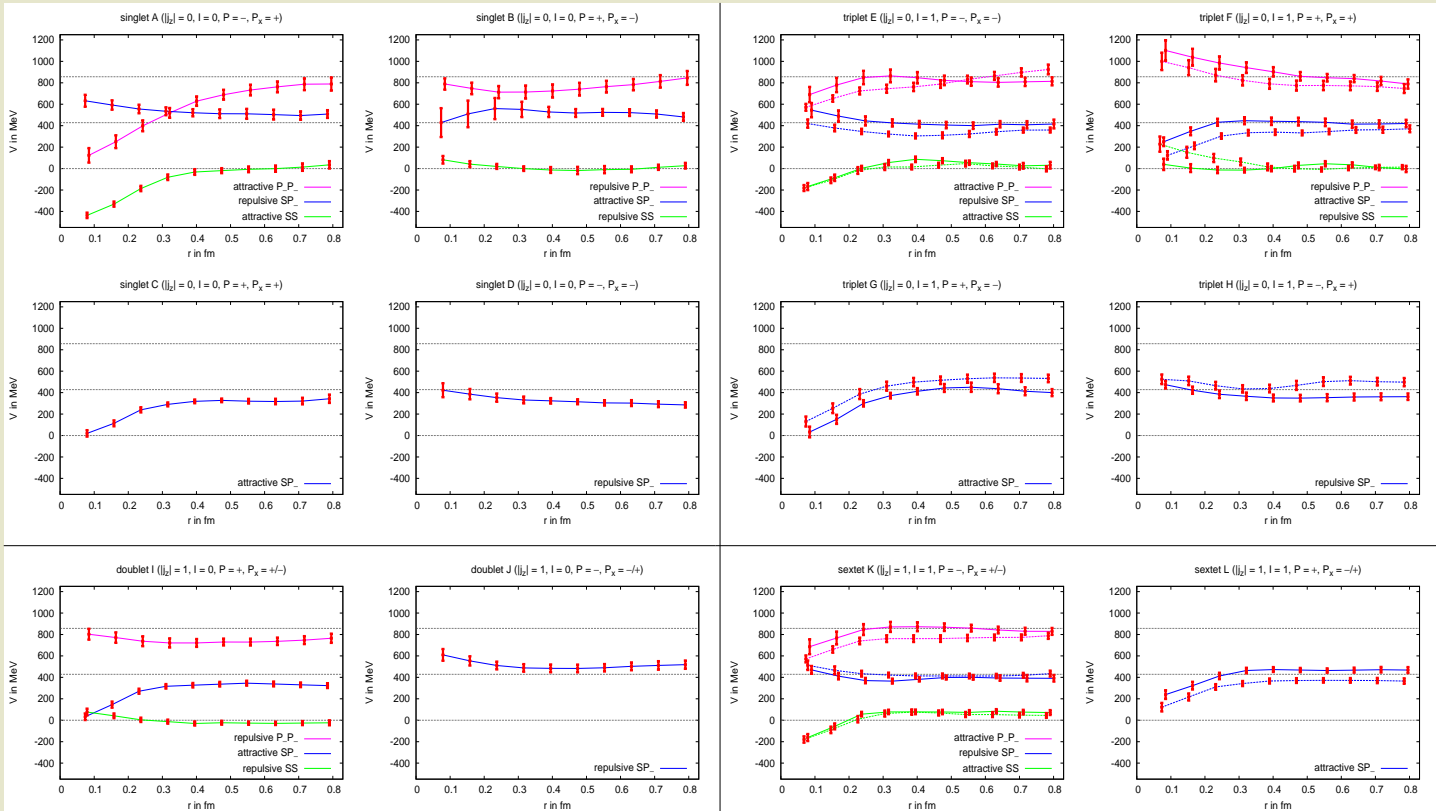
# Lattice setup

- ETMC gauge link ensembles:
  - $N_f = 2$  dynamical quark flavors.
  - Lattice spacing  $a \approx 0.079$  fm.
  - $24^3 \times 48$ , i.e. spatial lattice extent  $\approx 1.9$  fm.
  - Three different pion masses  $m_\pi \approx 340$  MeV,  $m_\pi \approx 480$  MeV,  $m_\pi \approx 650$  MeV.

[R. Baron *et al.* [ETM Collaboration], JHEP 1008, 097 (2010) [arXiv:0911.5061 [hep-lat]]

# $\bar{b}\bar{b}qq$ / $BB$ potentials (3)

- $I = 0$  (left) and  $I = 1$  (right);  $|j_z| = 0$  (top) and  $|j_z| = 1$  (bottom).



# $\bar{b}\bar{b}qq$ / $BB$ potentials (4)

## Why are there three different asymptotic values?

- Differences  $\approx 400$  MeV, approximately the mass difference of  $B_{0,1}^*$  ( $P = +$ ) and  $B^{(*)}$  ( $P = -$ ).
- Suggests that the three different asymptotic values correspond to  $B^{(*)}B^{(*)}$  potentials, to  $B^{(*)}B_{0,1}^*$  potentials and  $B_{0,1}^*B_{0,1}^*$  potentials.
- Can be checked and confirmed, by rewriting the  $\bar{b}\bar{b}qq$  creation operators in terms of meson-meson creation operators (Fierz transformation).
- Example:  $uu$ ,  $\Gamma = \gamma_3$  (attractive, lowest asymptotic value),

$$\begin{aligned}
 & (C\gamma_3)_{AB} \left( \bar{Q}_C(-\mathbf{r}/2)q_A^{(u)}(-\mathbf{r}/2) \right) \left( \bar{Q}_D(+\mathbf{r}/2)q_B^{(u)}(+\mathbf{r}/2) \right) \propto \\
 & \propto (B^{(*)})_{\uparrow}(B^{(*)})_{\downarrow} + (B^{(*)})_{\downarrow}(B^{(*)})_{\uparrow} - (B_{0,1}^*)_{\uparrow}(B_{0,1}^*)_{\downarrow} - (B_{0,1}^*)_{\downarrow}(B_{0,1}^*)_{\uparrow}.
 \end{aligned}$$

- Example:  $uu$ ,  $\Gamma = 1$  (repulsive, medium asymptotic value),

$$\begin{aligned}
 & (C1)_{AB} \left( \bar{Q}_C(-\mathbf{r}/2)q_A^{(u)}(-\mathbf{r}/2) \right) \left( \bar{Q}_D(+\mathbf{r}/2)q_B^{(u)}(+\mathbf{r}/2) \right) \propto \\
 & \propto (B^{(*)})_{\uparrow}(B_{0,1}^*)_{\downarrow} - (B^{(*)})_{\downarrow}(B_{0,1}^*)_{\uparrow} + (B_{0,1}^*)_{\uparrow}(B^{(*)})_{\downarrow} - (B_{0,1}^*)_{\downarrow}(B^{(*)})_{\uparrow}.
 \end{aligned}$$

# $\bar{b}\bar{b}qq$ / $BB$ potentials (5)

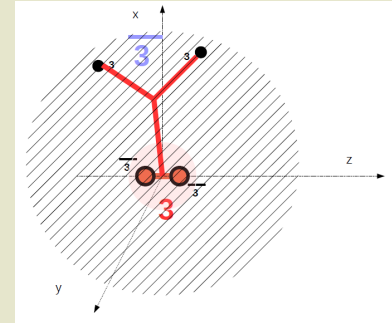
## Why are certain channels attractive and others repulsive? (1)

- Fermionic wave function must be antisymmetric (Pauli principle); in quantum field theory/QCD automatically realized.
- $qq$  isospin:  $I = 0$  antisymmetric,  $I = 1$  symmetric.
- $qq$  angular momentum/spin:  $j = 0$  antisymmetric,  $j = 1$  symmetric.
- $qq$  color:
  - $(I = 0, j = 0)$  and  $(I = 1, j = 1)$ : must be antisymmetric, i.e., a triplet  $\bar{3}$ .
  - $(I = 0, j = 1)$  and  $(I = 1, j = 0)$ : must be symmetric, i.e., a sextet  $6$ .
- The four quarks  $\bar{b}\bar{b}qq$  must form a color singlet:
  - $qq$  in a color triplet  $\bar{3}$  → static quarks  $\bar{b}\bar{b}$  also in a triplet  $3$ .
  - $qq$  in a color sextet  $6$  → static quarks  $\bar{b}\bar{b}$  also in a sextet  $\bar{6}$ .

# $\bar{b}\bar{b}qq$ / $BB$ potentials (6)

## Why are certain channels attractive and others repulsive? (2)

- Assumption: attractive/repulsive behavior of  $\bar{b}\bar{b}$  at small separations  $r$  is mainly due to 1-gluon exchange,
  - color triplet  $\bar{3}$  is attractive,  $V_{\bar{b}\bar{b}}(r) = -2\alpha_s/3r$ ,
  - color sextet  $\bar{6}$  is repulsive,  $V_{\bar{b}\bar{b}}(r) = +\alpha_s/3r$
 (easy to calculate in LO perturbation theory).



- Summary:
  - $(I = 0, j = 0)$  and  $(I = 1, j = 1)$  → attractive  $\bar{b}\bar{b}$  potential  $V_{\bar{b}\bar{b}}(r)$ .
  - $(I = 0, j = 1)$  and  $(I = 1, j = 0)$  → repulsive  $\bar{b}\bar{b}$  potential  $V_{\bar{b}\bar{b}}(r)$ .
- Expectation consistent with the obtained lattice results.
- Pauli principle and assuming “1-gluon exchange” at small  $r$  explains, why certain channels are attractive and others repulsive.**

# $\bar{b}\bar{b}qq$ / $BB$ potentials (7)

- Summary of  $\bar{b}\bar{b}qq$  /  $BB$  potentials:

$B^{(*)}B^{(*)}$ potentials:	attractive:	$1 \oplus 3 \oplus 6$	(10 states).
	repulsive:	$1 \oplus 3 \oplus 2$	( 6 states).
$B^{(*)}B_{0,1}^*$ potentials:	attractive:	$1 \oplus 1 \oplus 3 \oplus 3 \oplus 2 \oplus 6$	(16 states).
	repulsive:	$1 \oplus 1 \oplus 3 \oplus 3 \oplus 2 \oplus 6$	(16 states).
$B_{0,1}^*B_{0,1}^*$ potentials:	attractive:	$1 \oplus 3 \oplus 6$	(10 states).
	repulsive:	$1 \oplus 3 \oplus 2$	( 6 states).

- 2-fold degeneracy due to spin  $j_z = \pm 1$ .
- 3-fold degeneracy due to isospin  $I = 1, I_z = -1, 0, +1$ .

→ 24 **different**  $\bar{b}\bar{b}qq$  /  $BB$  potentials.

# $\bar{b}\bar{b}qq$ / $BB$ potentials (8)

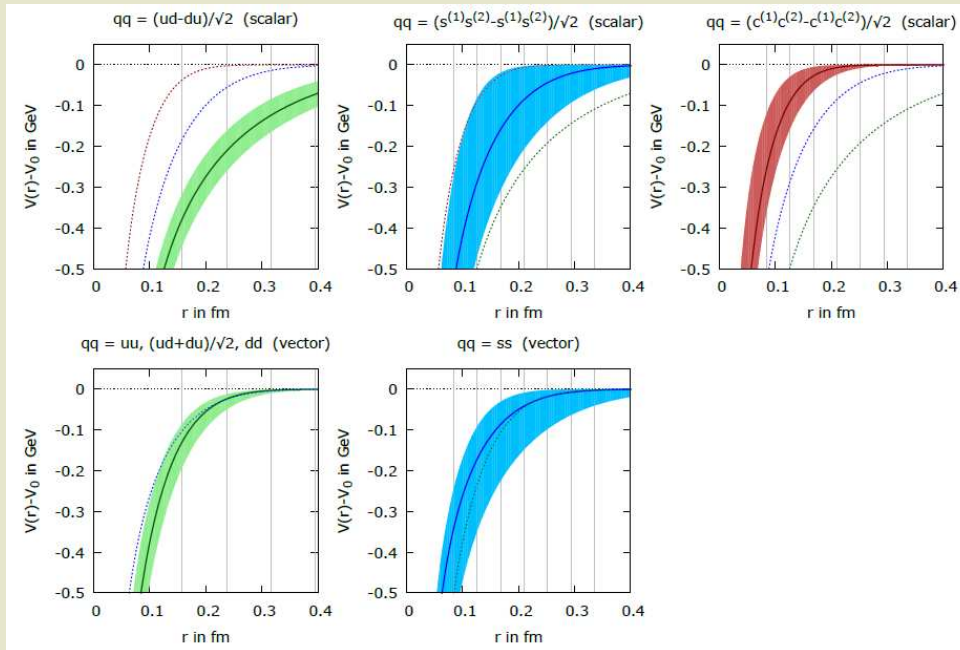
- Focus on the two attractive channels between  $B$  and  $B^*$ :
  - Scalar isosinglet ( $(I = 0, j = 0)$ , more attractive):  
 $qq = (ud - du)/\sqrt{2}$ ,  $\Gamma = (1 + \gamma_0)\gamma_5$ .
  - Vector isotriplet ( $(I = 1, j = 1)$ , less attractive):  
 $qq \in \{uu, (ud + du)/\sqrt{2}, dd\}$ ,  $\Gamma = (1 + \gamma_0)\gamma_j$ .
- Computations for  $qq = ll, ss, cc$  ( $l \in \{u, d\}$ ) to study the mass dependence.
- Parameterize lattice potential results by continuous functions obtained by  $\chi^2$  minimizing fits of

$$V_{\bar{b}\bar{b}}(r) = -\frac{\alpha}{r} \exp\left(-\left(\frac{r}{d}\right)^p\right) + V_0 :$$

- $1/r$ : 1-gluon exchange at small  $\bar{b}\bar{b}$  separations.
- $\exp(-(r/d)^p)$ : color screening at large  $\bar{b}\bar{b}$  separations due to meson formation.
- Fit parameters  $\alpha$ ,  $d$  and  $V_0$ ;  $p = 2$  from quark models.

# $\bar{b}bqq$ / $BB$ potentials (9)

- Potentials for  $qq = ll$ ,  $l \in \{u, d\}$  are wider and deeper than potentials for  $qq = ss, cc$ .  
 → **Good candidates to find tetraquarks are systems of two very heavy and two very light quarks, i.e.,  $\bar{b}bll$ .**



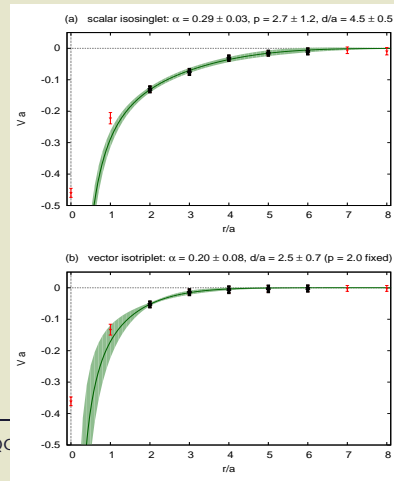
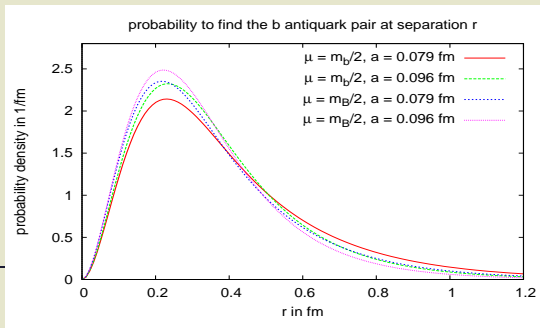


# $\bar{b}\bar{b}qq$ tetraquarks (1)

- Solve the Schrödinger equation for the relative coordinate of the heavy quarks  $\bar{b}\bar{b}$  using the previously computed  $\bar{b}\bar{b}qq / BB$  potentials,

$$\left( -\frac{1}{2\mu}\Delta + V_{\bar{b}\bar{b}}(r) \right) \psi(\mathbf{r}) = E\psi(\mathbf{r}) \quad , \quad \mu = m_b/2.$$

- Possibly existing bound states, i.e.,  $E < 0$ , indicate stable  $\bar{b}\bar{b}qq$  tetraquarks.
- There is a bound state for  $qq = (ud - du)/\sqrt{2}$  (i.e., the scalar isosinglet potential) and orbital angular momentum  $l = 0$  of  $\bar{b}\bar{b}$ , binding energy  $E = -90^{+43}_{-36}$  MeV with respect to the  $B + B^*$  threshold, i.e. confidence level  $\approx 2\sigma$ .
- No further bound states, in particular not for  $qq = ss, cc$  (i.e.,  $B_s B_s, B_c B_c$ ).



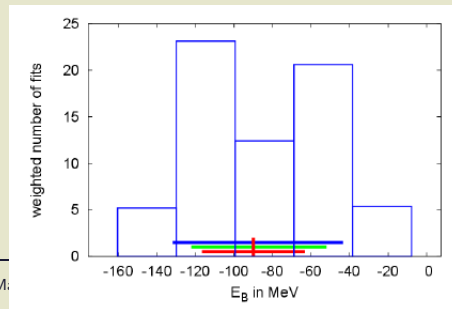
# $\bar{b}\bar{b}qq$ tetraquarks (2)

- Estimate the systematic error by varying input parameters:

- the  $t$  fitting range to extract the potential from effective masses,
- the  $r$  fitting range for

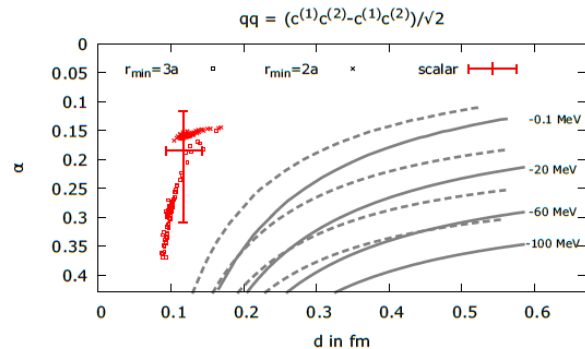
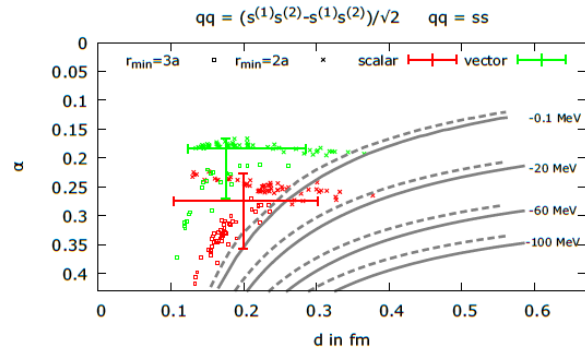
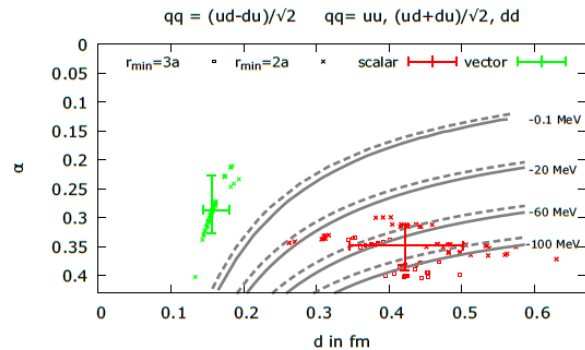
$$V_{\bar{b}\bar{b}}(r) = -\frac{\alpha}{r} \exp\left(-\left(\frac{r}{d}\right)^p\right) + V_0.$$

- Right: isoline plots of the binding energy  $E$  for  $l = 0$ .
- Bottom: histogram for the binding energy  $E$  for  $qq = (ud - du)/\sqrt{2}$  and  $l = 0$ .



M:

QCD, II



# $\bar{b}\bar{b}qq$ tetraquarks (3)

- To quantify “no binding”, we list for each channel the factor, by which the reduced mass  $\mu$  in the Schrödinger equation has to be multiplied, to obtain a tiny but negative energy  $E$  (again for  $l = 0$ ).

$qq$	spin	factor
$(ud - du)/\sqrt{2}$	scalar	0.46
$uu, (ud + du)/\sqrt{2}, dd$	vector	1.49
$(s^{(1)}s^{(2)} - s^{(2)}s^{(1)})/\sqrt{2}$	scalar	1.20
$ss$	vector	2.01
$(c^{(1)}c^{(2)} - c^{(2)}c^{(1)})/\sqrt{2}$	scalar	2.57

- Factors  $\ll 1$  indicate strongly bound states, while for values  $\gg 1$  bound states are essentially excluded.
- Light quarks ( $u/d$ ) are unphysically heavy (correspond to  $m_\pi \approx 340$  MeV); physically light  $u/d$  quarks yield similar results.
- Mass splitting  $m(B^*) - m(B) \approx 50$  MeV, neglected at the moment, is expected to weaken binding (will be discussed below).

# $\bar{b}\bar{b}qq$ tetraquarks (4)

What are the quantum numbers of the predicted  $\bar{b}\bar{b}qq$  tetraquark?

- $I(J^P) = 0(1^+)$ .

- Light scalar isosinglet:  $qq = (ud - du)/\sqrt{2}$ ,  $I = 0$ ,  $j = 0$  in a color  $\bar{3}$ ,  $\bar{b}\bar{b}$  in a color  $3$  (antisymmetric) ... as discussed above.
- Wave function of  $\bar{b}\bar{b}$  must also be antisymmetric (Pauli principle).
  - \*  $\bar{b}\bar{b}$  is flavor symmetric.
  - \*  $\bar{b}\bar{b}$  spin must also be symmetric, i.e.,  $j_b = 1$ .
- **The predicted  $\bar{b}\bar{b}qq$  tetraquark has isospin  $I = 0$ , spin  $J = 1$ .**
- We study a state, which correspond for large  $\bar{b}\bar{b}$  separations to a pair of  $B^{(*)}$  mesons in a spatially symmetric s-wave.
- **The predicted  $\bar{b}\bar{b}qq$  tetraquark has parity  $P = +$**  (the product of the parity quantum numbers of the two mesons, which are both negative).

# Inclusion of heavy spin effects

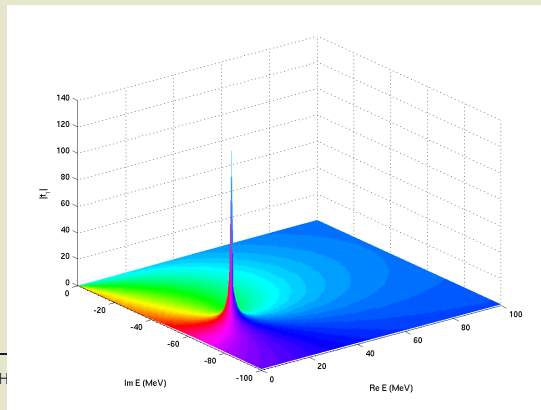
- Heavy spin effects have been neglected so far, e.g. mass splitting  $m_{B^*} - m_B \approx 46$  MeV.
- Mass splitting  $m_{B^*} - m_B$  is, however, of the same order of magnitude as the previously obtained binding energy  $E = -90_{-36}^{+43}$  MeV.
- Moreover, two competing effects:
  - The attractive  $\bar{b}b u d$  channel corresponds to a linear combination of  $BB^*$  and/or  $B^*B^*$ .
  - The  $BB^*$  interaction is a superposition of attractive and repulsive  $\bar{b}b u d$  potentials.
- **Will there still be a bound state, when heavy spin effects are taken into account?**
  - Yes.
  - We include heavy spin effects by solving a coupled channel Schrödinger equation. [P. Bicudo, J. Scheunert, M.W., Phys. Rev. D **95**, 034502 (2017) [arXiv:1612.02758]]
  - Binding energy  $E = -59_{-30}^{+38}$  MeV.
  - Tetraquark is approximately a 50%/50% superposition of  $BB^*$  and  $B^*B^*$  (strong attraction more important than light constituents).

# $\bar{b}\bar{b}qq$ tetraquark resonances (1)

- Most hadrons are unstable, i.e., resonances.
- If a  $\bar{b}\bar{b}qq$  potential  $V_{\bar{b}\bar{b}}(r)$  is not sufficiently attractive to host a bound state, there could still be a clear resonance.
- Comparatively easy to investigate within our approach (since we have potentials  $V_{\bar{b}\bar{b}}(r)$ , no Lüscher method etc. necessary).
- Use standard methods from scattering theory:
  - Solve Schrödinger equation with potential  $V_{\bar{b}\bar{b}}(r)$  and appropriate boundary conditions (incident plane wave, outgoing spherical wave)
    - partial wave amplitudes  $f_l(E)$ .
  - Use partial wave amplitudes  $f_l(E)$  to ...
    - \* ... determine phase shifts and contributions of partial waves to total cross section
      - peak indicates resonance mass.
    - \* ... determine poles of the S or the T matrix in the complex energy plane (correspond to poles of  $f_l(E)$ )
      - real part of a pole  $\equiv$  resonance mass
      - imaginary part of a pole  $\equiv$  resonance width.

# $\bar{b}\bar{b}qq$ tetraquark resonances (2)

- Exploratory study mostly for  $qq = (ud - du)/\sqrt{2}$  (i.e., the scalar isosinglet potential) and orbital angular momentum  $l = 1$  of  $\bar{b}\bar{b}$ :
  - There is a resonance for  $qq = (ud - du)/\sqrt{2}$  and  $l = 1$ :
    - Resonance mass  $E = +17_{-4}^{+4}$  MeV above the  $BB$  threshold.
    - Decay width  $\Gamma_{\rightarrow B+B} = 112_{-103}^{+90}$  MeV.
    - Quantum numbers  $I(J^P) = 0(1^-)$ .
- [P. Bicudo, M. Cardoso, A. Peters, M. Pflaumer, M.W., arXiv:1704.02383].
- There do not seem to be resonances in other channels ( $l > 1$ , vector isotriplet potential, heavier quarks  $qq$ ).



# $\bar{b}b\bar{q}q$ / $\bar{B}B$ potentials

- Exploring the existence of  $\bar{b}b\bar{q}q$  tetraquarks in the same way is more difficult:
  - $\bar{b}bq\bar{q}$  (discussed on previous slides) can decay into:
    - \*  $\bar{B} + \bar{B}$ .  
“Easy” ... on the level of the Schrödinger equation for the relative coordinate of the two  $\bar{b}$  quarks (step (2) of the BO approximation).
  - $\bar{b}b\bar{q}q$  can decay into:
    - \*  $\bar{B} + B$ .  
“Easy” ... on the level of the Schrödinger equation for the relative coordinate of the  $\bar{b}$  quark and the  $b$  quark (step (2) of the BO approximation).
    - \*  $\bar{b}b + \bar{q}q$  (“bottomonium + pion”).  
“Rather hard” ... on the level of lattice QCD, when computing the  $\bar{b}b$  potentials in the presence of  $\bar{q}q$  (step (1) of the BO approximation).
      - A potential can be relevant for a  $\bar{b}b\bar{q}q$  tetraquark (if  $\bar{q}q$  is close to  $\bar{b}b$ ) ...
      - ... or just a  $\bar{b}b$  potential shifted by the mass of a  $\bar{q}q$  meson.
- Work in progress.
  - [A. Peters, P. Bicudo, L. Leskovec, S. Meinel and M.W., PoS LATTICE 2016, 104 (2016) [arXiv:1609.00181]]
  - [A. Peters, P. Bicudo and M.W., EPJ Web Conf. 175, 14018 (2018) [arXiv:1709.03306]]



# Summary and outlook

- Computation of  $3 \times 24$  different  $\bar{b}bqq / BB$  potentials.
- Prediction of a stable  $\bar{b}bqq$ ,  $qq = (ud - du)/\sqrt{2}$  tetraquark.
  - Quantum numbers  $I(J^P) = 0(1^+)$ .
  - Binding energy  $E = -59^{+38}_{-30}$  MeV with respect to the  $B + B^*$  threshold.
- Prediction of a  $\bar{b}bqq$ ,  $qq = (ud - du)/\sqrt{2}$  tetraquark resonance.
  - Quantum numbers  $I(J^P) = 0(1^-)$ .
  - Resonance mass  $E = +17^{+4}_{-4}$  MeV above the  $B + B$  threshold.
  - Decay width  $\Gamma_{\rightarrow B+B} = 112^{+90}_{-103}$  MeV.
- Future plans:
  - Explore  $\bar{b}bqq$  tetraquark resonances in detail.
  - Investigate the structure of the predicted  $I(J^P) = 0(1^+)$  tetraquark ... is it a mesonic molecule or rather a diquark-antidiquark?
  - Study  $\bar{b}b\bar{q}q / BB$ , which is experimentally more relevant ( $Z_b(10610)^+$ ,  $Z_b(10650)^+$ , ...), but theoretically much harder.