

Plans for meson spectroscopy

Marc Wagner

Humboldt-Universität zu Berlin, Institut für Physik

Theorie der Elementarteilchen – Phänomenologie/Gittereichtheorie

mcwagner@physik.hu-berlin.de

<http://people.physik.hu-berlin.de/~mcwagner/>

ETMC meeting – Bern

March 24, 2011

Introduction

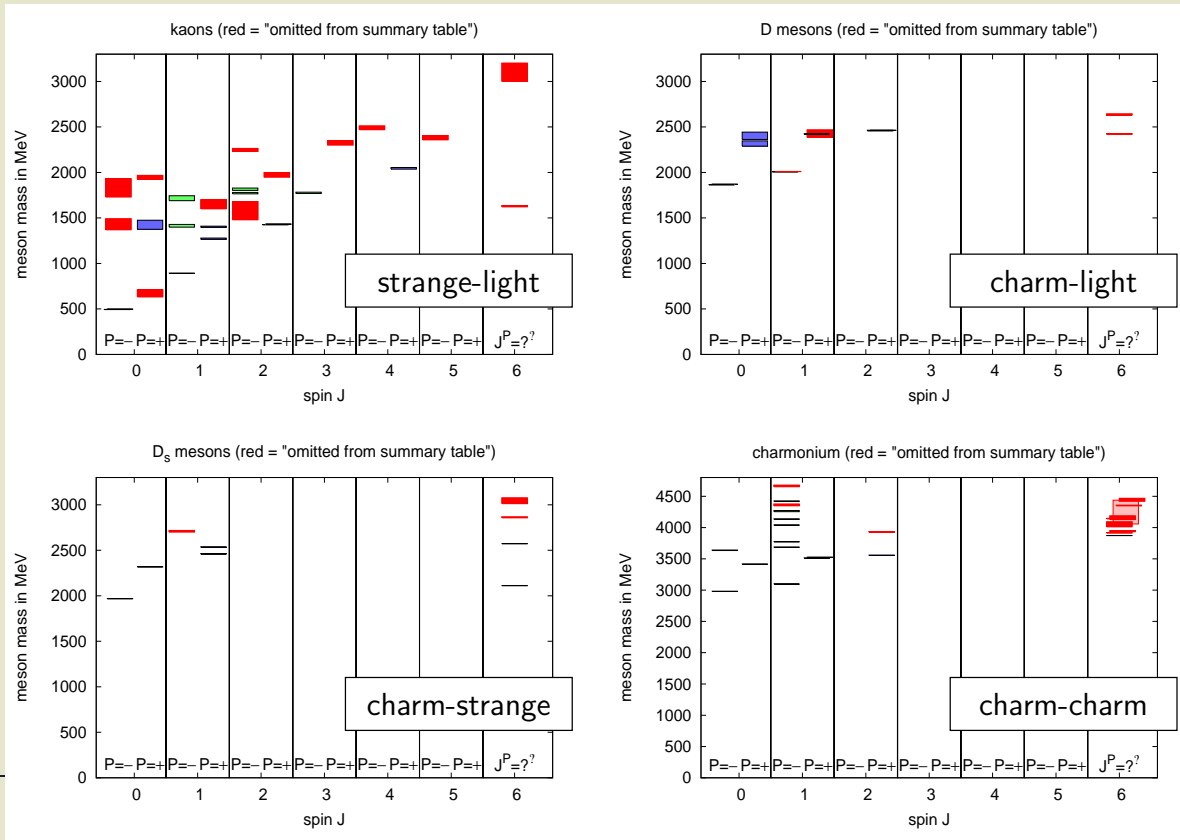
- In May I will start a five year position as leader of a junior research group (Emmy Noether Programme) at Humboldt University Berlin.
- The position is associated with two scientific projects:
 - **Project 1: computation of the spectrum of s and c mesons.**
 - * Goal: Compute the s and c meson spectrum (kaons, D mesons, D_s mesons, charmonium) as fully as possible with $N_f = 2 + 1 + 1$ ETMC gauge field configurations.
 - **Project 2: (mixed) action simulations at fixed topology.**
- Project 1 requires investigation, implementation and understanding of certain lattice techniques, which might be of general interest for members of ETMC:
 - I would like to let you know, what we plan to do in the next months/years such that you can possibly profit from our experience.
 - If some of you are interested, any form of active collaboration will be very welcome.

Physical goals, s/c meson spectrum (1)

- Compute the s and c meson spectrum as fully as possible:
 - Consider all mesons, which have at least one s or c quark, i.e.
 - * kaons (strange-light mesons), [light = up or down]
 - * D mesons (charm-light mesons), [light = up or down]
 - * D_s mesons (charm-strange mesons),
 - * charmonium (charm-charm mesons),
 - * possibly strangeonium (strange-strange mesons).
 - Consider parity \pm , charge conjugation \pm , radial and orbital excitations.
 - Lattice setup:
 - * $N_f = 2 + 1 + 1$ flavor ETMC gauge field configurations.
 - * s/c quarks via an Osterwalder-Seiler mixed action setup (no flavor breaking, only parity is explicitly broken).

Physical goals, s/c meson spectrum (2)

- Experimental status (Particle Data Group): 73 known states.



Physical goals, s/c meson spectrum (3)

- Why is a lattice computation of the s and c meson spectrum important?
 - Some mesons, e.g. D_s , η_c , J/ψ , have been measured experimentally with high precision and can also be computed on the lattice very accurately
→ ideal candidates to test QCD by means of lattice QCD.
 - Some mesons are only poorly understood
→ lattice QCD is the perfect tool to clarify the situation:
 - * 31 meson states labeled with “omitted from summary table” (states colored red), i.e. vague experimental signals, experimental contradictions, states not well established.
 - * Example $X(3872)$ ($\bar{c}c$ state): mass not as expected from quark models; could be a $D^0-\bar{D}^*(2007)^0$ molecule, a bound diquark-antidiquark, ... or models could yield wrong answers.
 - * Example $D_{s0}^*(2317)$, $D_{s1}(2460)$: masses significantly lower than expected from quark models, almost equal or even lower than the corresponding D mesons; could be D - K molecules, tetraquarks, ...

Physical goals, s/c meson spectrum (4)

- Why is a lattice computation of the s and c meson spectrum important?
 - Lattice QCD predictions of meson masses give valuable input for future experiments.
 - Comprehensive new information expected from existing and new facilities (BABAR, Belle, CEBAF, CLEO, upgraded BES, CDF, D0, FAIR (PANDA), LHC, ...) ... i.e. a “hot topic”.
 - Lattice results for the s and c meson spectrum exist, but no comprehensive picture available at the moment (different discretizations, scale setting methods, numbers of quark flavors, sometimes rather coarse lattice spacings, unphysically heavy u/d quarks, no extrapolations).

Technical aspects, overview

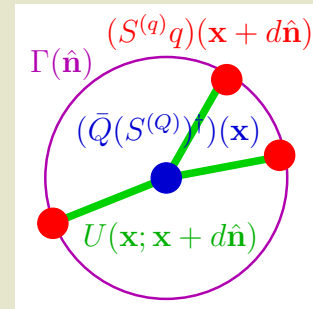
- Construction and selection of suitable meson creation operators, application of the “generalized eigenvalue problem” .
- Efficient computation of quark propagators and correlation functions (stochastic sources, one-end-trick, distillation).
- How to deal with multiparticle states?

Construction/selection of operators (1)

- Construction of suitable meson creation operators:
 - Goal: construct a set of operators, which almost exclusively excites the states you want to compute; then the masses of these states can be extracted from the corresponding correlation functions/matrices at rather small temporal separation, where statistical errors are also small.
 - General form of a meson creation operator in a continuum-like notation:

$$\mathcal{O}(\mathbf{x}) = \left(\bar{Q}(S^{(Q)})^\dagger \right)(\mathbf{x}) \int d\hat{\mathbf{n}} \Gamma(\hat{\mathbf{n}}) U(\mathbf{x}; \mathbf{x} + d\hat{\mathbf{n}}) \left(S^{(q)} q \right)(\mathbf{x} + d\hat{\mathbf{n}}).$$

- Degrees of freedom:
 - * Spin structure, parity (γ matrices).
 - * Angular momentum structure (displacement of valence quarks, lattice versions of spherical harmonics).
 - * Width of the operator (smearing techniques).
 - * Nodes in the generated wavefunctions (application of derivatives).



Construction/selection of operators (2)

- Selection of suitable meson creation operators:
 - Optimize operators by minimizing effective masses at small temporal separations.
 - Choose a small set of operators, which are sufficiently different:
 - * **What are suitable criteria? Can this selection be automated in an effective way?**
 - * The HS Collaboration recommends minimizing the condition number of “normalized” correlation submatrices at small temporal separation (condition number = largest eigenvalue/smallest eigenvalue).
[C. Morningstar, PoS [LATTICE2008](#), 009 (2008)]

Construction/selection of operators (3)

- Application of the “generalized eigenvalue problem”:

$$C(t)v_n(t, t_0) = \lambda_n(t, t_0)C(t_0)v_n(t, t_0) \quad ,$$

$$E_n = \lim_{t \rightarrow \infty} E_n^{\text{eff}}(t) \quad , \quad E_n^{\text{eff}}(t) = \ln \left(\frac{\lambda_n(t, t_0)}{\lambda_n(t+1, t_0)} \right) \quad , \quad n = 1, \dots, N.$$

- Is it advisable to determine ground states from correlation matrices solving a generalized eigenvalue problem or is the effective mass of a single optimized correlation function sufficient?
- My current experience with static-light mesons, static-light baryons, the static potential (Wilson loops), ... indicates that there are no practical benefits in using correlation matrices.
- However B. Blossier et al. showed that the difference between the mass E_n and the effective mass $E_n^{\text{eff}}(t)$ is proportional to $e^{-(E_{N+1}-E_n)t}$, i.e. decreases exponentially with respect to t , where E_{N+1} is the mass of the first state “out of the basis”.

[B. Blossier, M. Della Morte, G. von Hippel, T. Mendes and R. Sommer, JHEP 0904, 094 (2009)]

Efficient computation of propagators (1)

- Computation of all-to-all quark propagators would be ideal, is, however, computationally prohibitively expensive.

Efficient computation of propagators (2)

- Strategies to maximize efficiency:
 - Unbiased stochastic estimation of all-to-all quark propagators (stochastic sources, possibly diluted).
 - One-end-trick:
 - (+) Eliminates statistical noise on “one end” of the correlation function.
 - (–) Requires spin diluted sources.
 - (–) Requires separate inversions for each operator.
 - Distillation:
 - (+) A specific type of smearing recently proposed by the HS Collaboration.
 - (–) Computationally very expensive (requires low lying eigenmodes of the lattice Laplacian).

[M. Peardon *et al.* [HSC], *Phys. Rev. D* **80**, 054506 (2009)]

- Implement and compare these approaches with respect to efficiency.

How to deal with multiparticle states? (1)

- At (close to) realistic pion masses ($m_{\text{PS}} \lesssim 300 \text{ MeV}$) most excited states have the same quantum numbers as lighter multiparticle states (e.g. ground state + pion(s)), i.e. are resonances.

How to deal with multiparticle states? (2)

- How to deal with contamination by such multiparticle states?
 - Usually multiparticle states are just ignored (one assumes e.g. that a state created by a two-quark meson operator has negligible overlap to multiparticle states) ... easy, but questionable/dangerous.
 - Lüscher's method to extract resonances from the volume dependence of the spectrum ... theoretically sound, but computationally very demanding/not applicable in practice.
[\[M. Lüscher, Nucl. Phys. B 364, 237 \(1991\)\]](#)
 - Compute the overlap of e.g. a two-quark meson operator and a four-quark multiparticle operator to demonstrate that the effect of this particular multiparticle state is indeed negligible.
[\[C. McNeile and C. Michael \[UKQCD\], Phys. Rev. D 63, 114503 \(2001\)\]](#)
[\[C. McNeile, C. Michael and P. Pennanen \[UKQCD\], Phys. Rev. D 65, 094505 \(2002\)\]](#)
[\[C. McNeile, C. Michael and G. Thompson \[UKQCD\], Phys. Rev. D 70, 054501 \(2004\)\]](#)
 - Further approaches/ideas?