

Electric conductivity of hadronic matter

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In collaboration with

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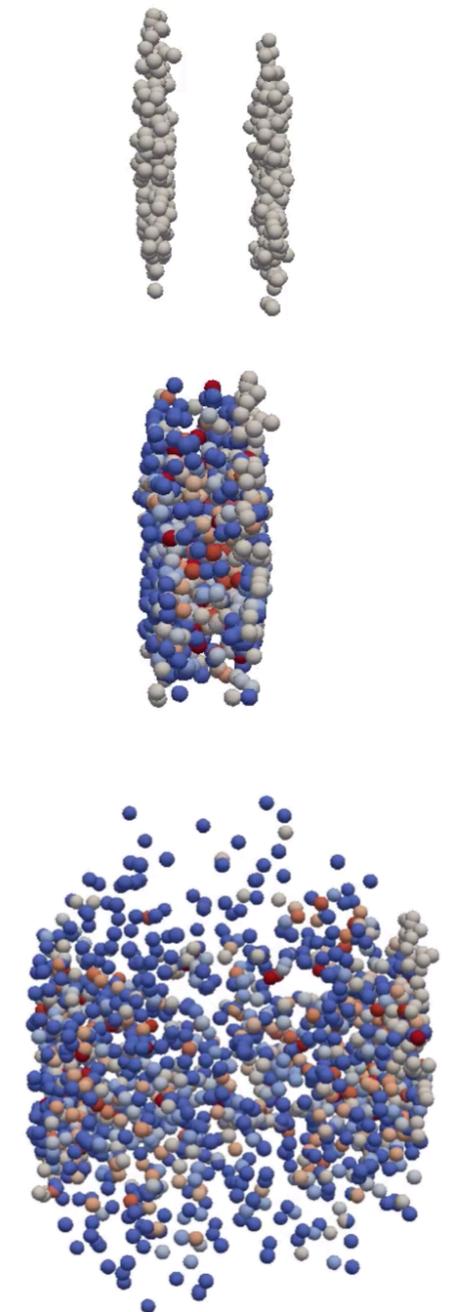
See also:

arXiv:1810.12527



Introduction

- Studying heavy ion collisions in order to understand properties of QCD matter
- Transport models give unique insights on dynamical evolution of heavy ion collisions
- Characterize the interacting hadron gas by its linear transport coefficients



Pb-Pb collision with 17.3 GeV center-of-mass energy
(by J. Mohs)



Definition: electric conductivity σ_{el}

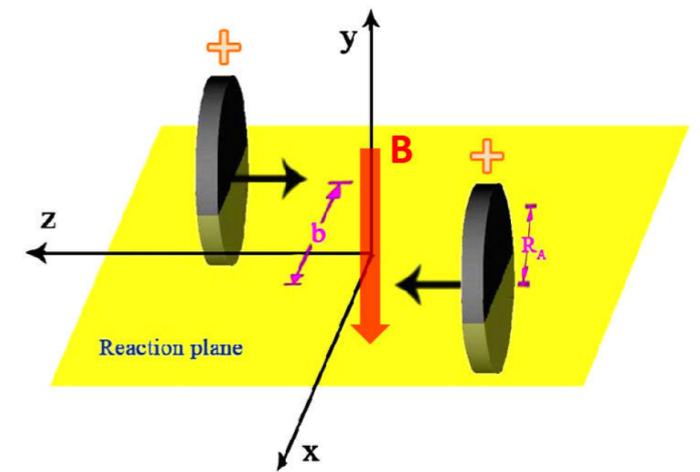
Ohm's law: $\vec{J}_Q(t, \vec{x}) = \sigma_{el} \vec{E}(t, \vec{x})$

- Electric current \vec{J}_Q
- Electric conductivity σ_{el}
- Electric field \vec{E}

Conservation of $U(1)$ electric charge leads to a transport of electric charge \rightarrow transport coefficient

Applications of σ_{el}

- High velocities of charged particle will lead to big electric and magnetic fields $\sim m_\pi^2$.
-> Time evolution depends on σ_{el}
- Soft dilepton and photon rates are related to the electric conductivity
- Magneto-hydrodynamic description of the expanding plasma after the collision needs the values of transport coefficients



X-G Huang Rept. Prog. Phys. 79, 076302, [2016]

Model: SMASH

- Simulating **M**any **A**ccelerated **S**trongly-interacting **H**adrons
- Hadronic transport approach, effectively solving Boltzmann equation

$$p^\mu \partial_\mu f_i(x, p) + m_i F^\alpha \partial_\alpha^p f_i(x, p) = C_{coll}^i$$

- Uses a geometric collision criterion

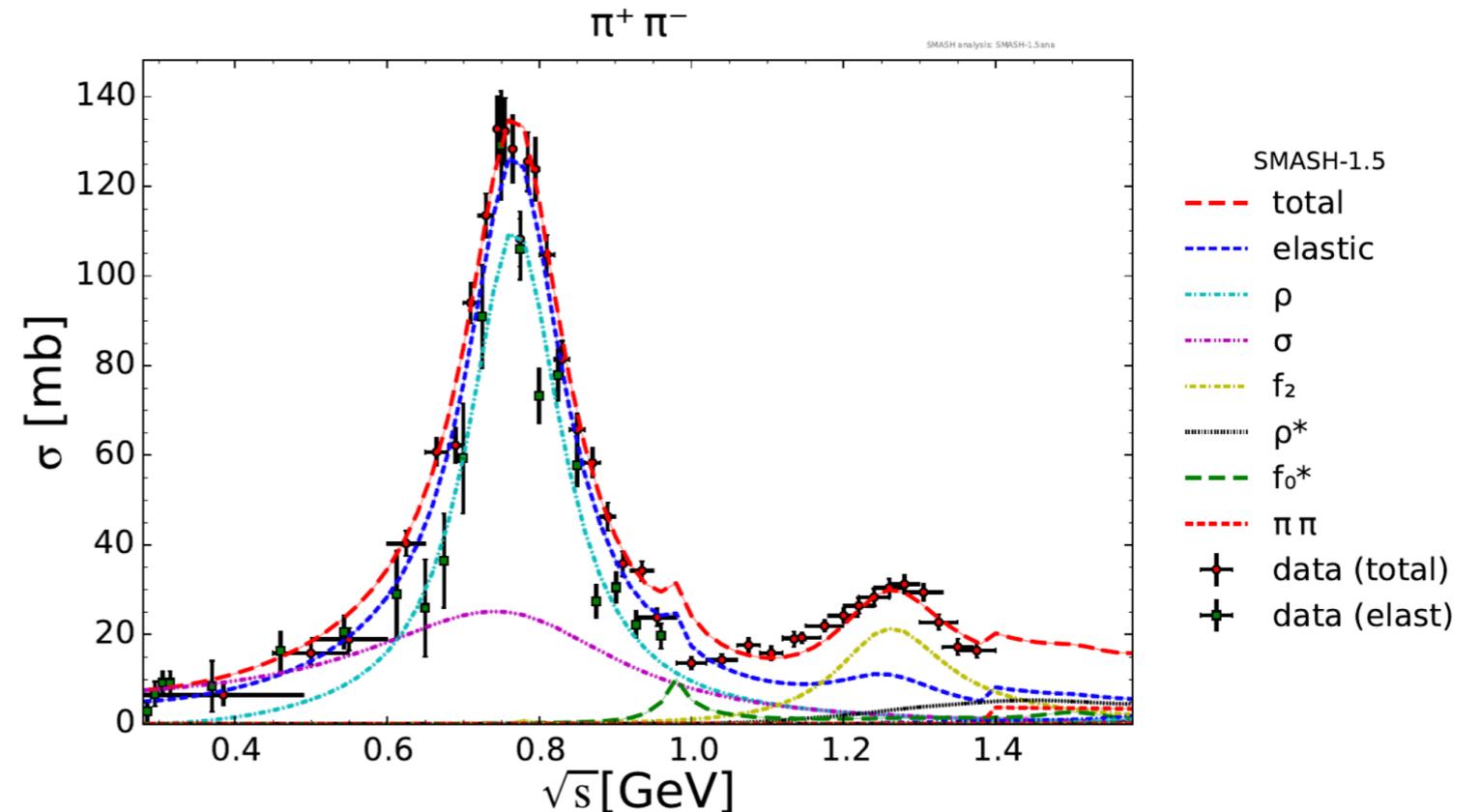
$$d_{coll} < \sqrt{\frac{\sigma_{tot}}{\pi}}$$

- Good description of low energy heavy ion collisions and of late stages of high energy heavy ion collisions

Model: SMASH

Weil et al PRC.94.054905 [2016]

- SMASH contains all baryons and mesons with masses up to $\sim 2\text{GeV}$
- $2 \leftrightarrow 1$ and $2 \leftrightarrow 2$ collision modeled through resonance formation and decay
- Resonance properties obtained from experiment



Particle	Mass (MeV)	Decay width (MeV)
π	138	0
ρ	776	149
K	494	0
K^*	892	50.8
N	938	0
Δ	1232	117

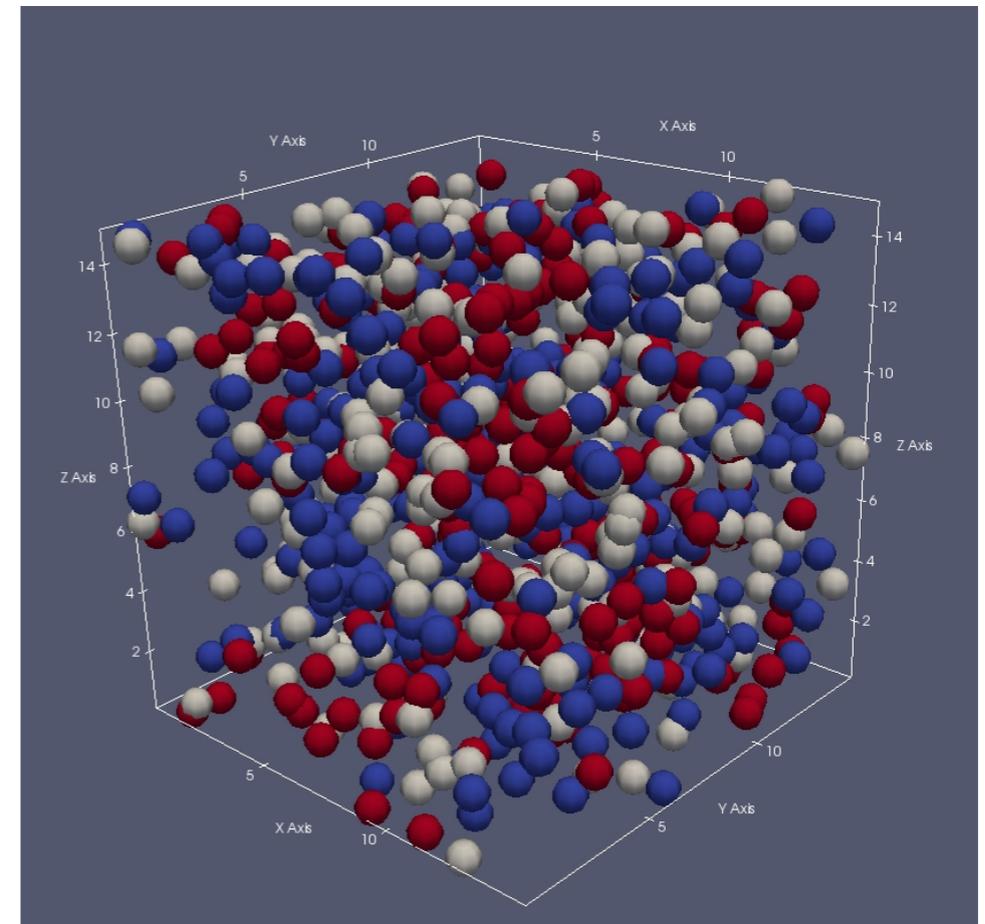
Box setup

- Box with periodic boundary conditions, simulating infinite matter

- Initial momenta distribution:

$$\frac{d^3N}{dp^3} \sim \exp\left(-\sqrt{\vec{p}^2 + m^2}/T\right)$$

- One has to check for thermal/chemical equilibrium

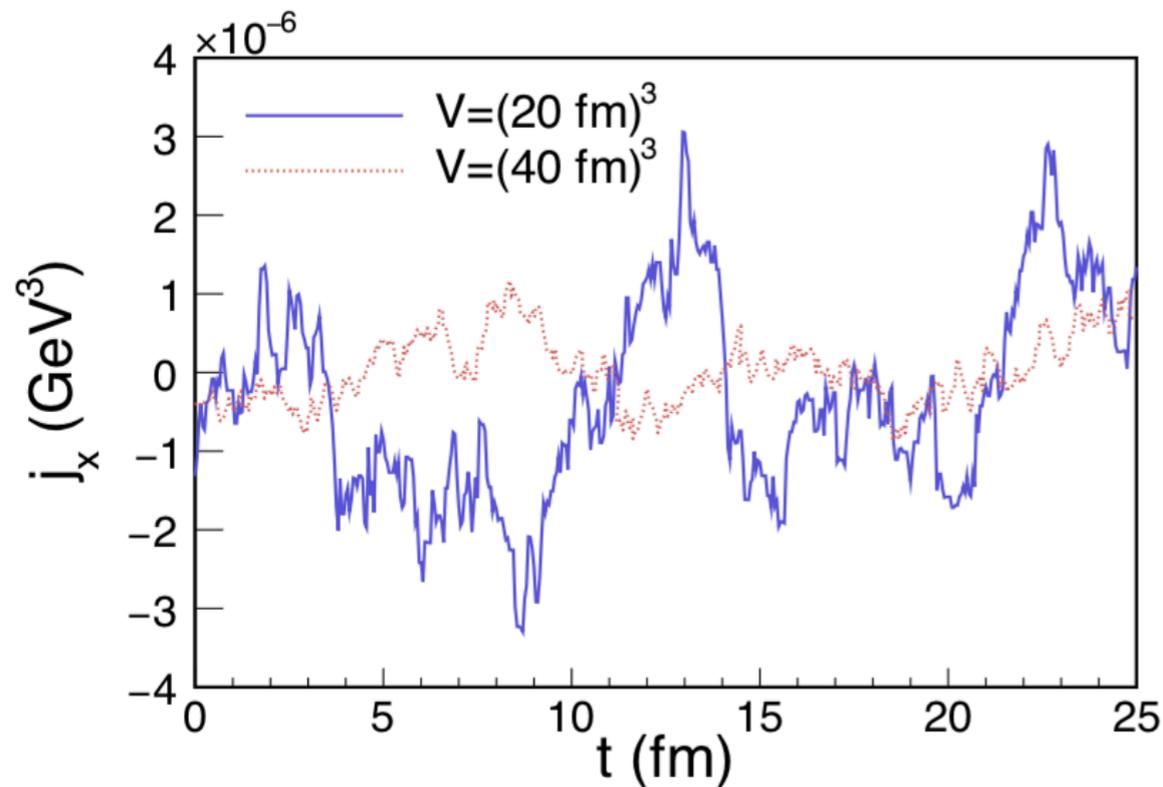


Box filled with π , $m_\pi = 0.138 GeV$
interacting via constant elastic cross
section $\sigma_{tot} = 30 mb$

Electric current

For a set of particles:

$$j_Q^i(t) = \frac{1}{V} \sum_{k=1}^N \frac{q_k p_k^i}{p_k^0} \Big|_t \sim \frac{1}{V} \int d^3x J_Q^i(t, \vec{x})$$



Box filled with π , interacting via constant elastic cross section $\sigma_{tot} = 30mb$ at $T = 125MeV$

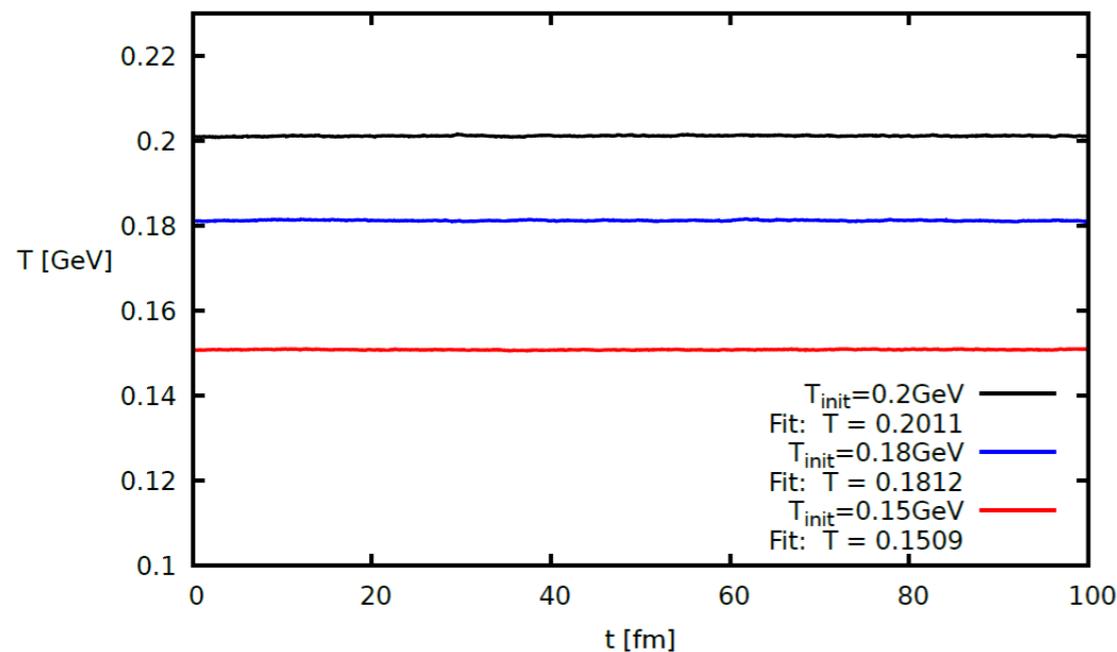
- Pions interacting as hard spheres
- Since no electric field is applied, the current fluctuates around zero

- Strength of fluctuation $\sim \frac{1}{\sqrt{V}}$

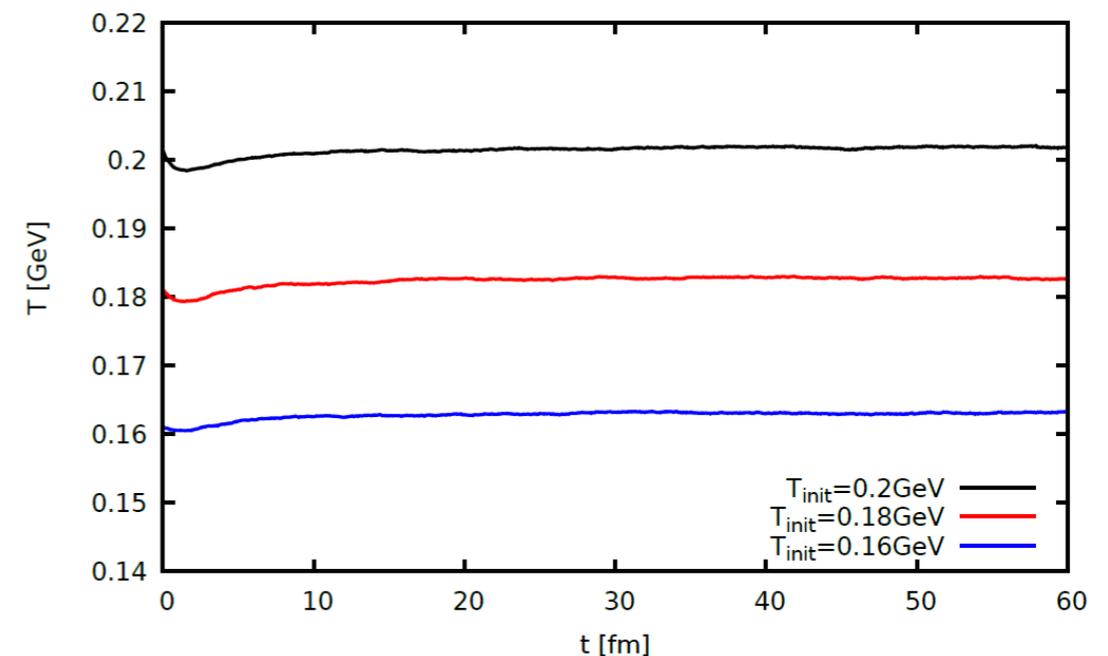
Calculation of the temperature

- Thermal equilibrium has to be reached in the box
- Calculate temperature by fitting the pion momenta spectrum

$$\frac{d^3N}{dp^3} \propto V e^{-\sqrt{\vec{p}^2 + m^2}/T}$$



Box filled with π , $m_\pi = 0.138 \text{ GeV}$ interacting via constant elastic cross section $\sigma_{\text{tot}} = 30 \text{ mb}$



Box filled with π and ρ mesons interaction via $2 \leftrightarrow 1$ resonance formation/decay

Green-Kubo formalism

- Well-known formalism to extract transport coefficients
- Based on linear response theory
- Possible to extract transport properties by looking at fluctuations of currents around equilibrium

$$\sigma_{el} = \frac{V}{3T} \int_0^{\infty} \langle \vec{j}_Q(t) \vec{j}_Q(0) \rangle dt$$

Correlation function

- How does $\langle \vec{j}_Q(t) \vec{j}_Q(0) \rangle$ look like?

- Since Ohm's law is inconsistent with causality it is modified to:

$$J_Q^i(t, \vec{x}) = \int_V d^3\vec{x}' \int_{t_0}^t dt' \Sigma^{ij}(t-t'; \vec{x} - \vec{x}') E^i(t', \vec{x}') + \Xi_Q^i(t, \vec{x})$$

where Σ^{ij} is the memory kernel of the electric conductivity and $\Xi_Q^i(t, \vec{x})$ a fluctuation term

- With an exponential Ansatz for the memory kernel:

$$\Sigma^{ij}(t-t'; \vec{x} - \vec{x}') = \frac{\sigma_{el} \delta^{ij}}{\tau_Q} \delta^{(3)}(\vec{x} - \vec{x}') \exp\left(-\frac{|t-t'|}{\tau_Q}\right),$$

It follows with the fluctuation dissipation theorem that the correlation function of the spatial average of the current-current correlator has an exponential decay

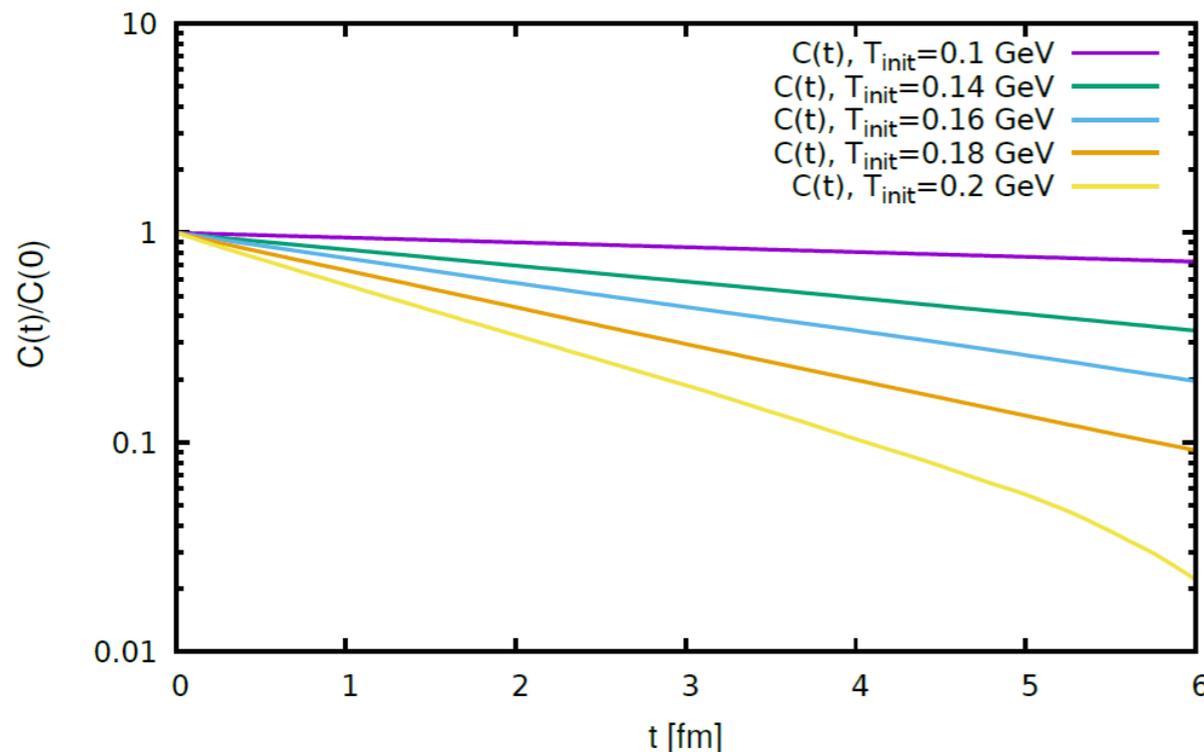
$$\langle j_Q^i(t) j_Q^j(t') \rangle = \frac{\sigma_{el} \delta^{ij} T}{\tau_Q V} \exp\left(-\frac{|t-t'|}{\tau_Q}\right) \quad \rightarrow \text{Ansatz: } \langle \vec{j}_Q(t) \vec{j}_Q(0) \rangle = C(0) e^{-t/\tau_Q}$$

Murase, Hirano
arXiv:1304.3243 [nucl-th]

Correlation function

Correlation function:
$$C(t) \equiv \langle j_Q^i(t) j_Q^i(0) \rangle = \frac{1}{K-t} \sum_{s=0}^{K-t} j_Q^i(s\Delta t) j_Q^i(s\Delta t + t)$$

Where K is the total number of time steps



Correlation function of a box filled with pions interacting with a constant cross section $\sigma_{tot} = 30mb$ for different initialization temperatures.

- Correlation function follows an exponential decay
 -> extract values of τ_Q and $C(0)$ by fitting an exponential function

$$C(t) = C(0)e^{-\frac{t}{\tau_Q}}$$

$$\longrightarrow \sigma_{el} = \frac{C(0)V\tau_Q}{T}$$

- Error grows for later times
 -> find cut-off criterion

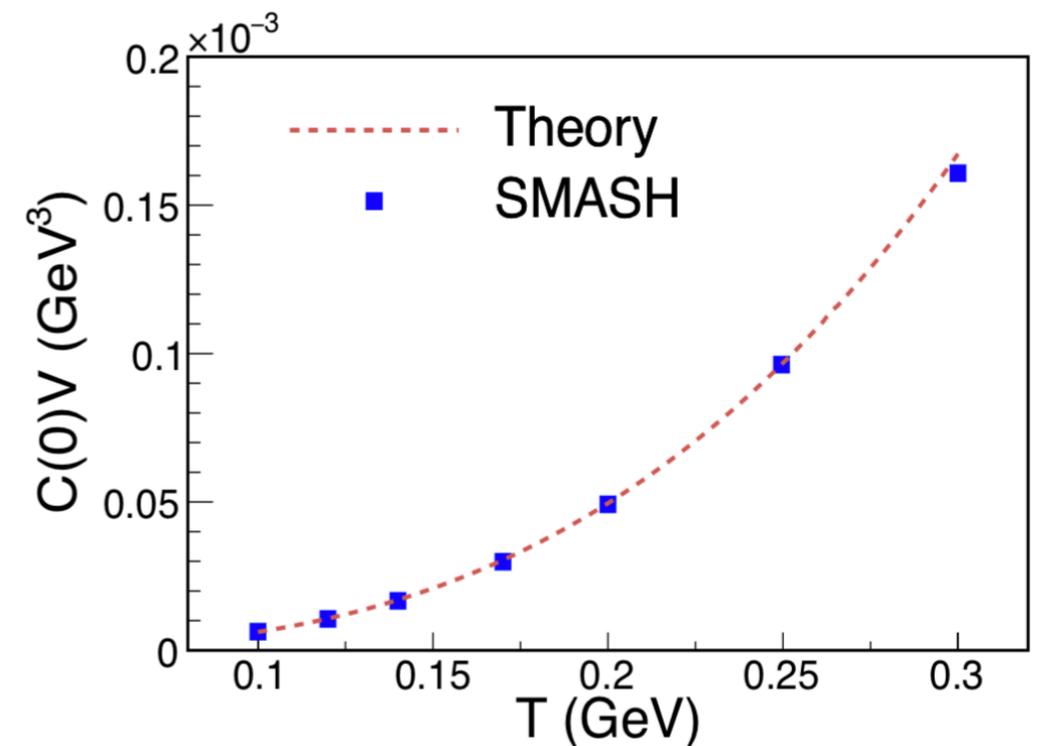
Correlation function at t=0

- Can be computed analytically
- Well defined in thermal equilibrium
- Good cross check if calculation is done correctly

For a massless gas:
$$C(0) = \frac{2(qe)^2 n}{3V}$$

$$C(0)V = \sum_{a=1}^{N_s} \frac{g_a q_a^2 e^2}{6\pi^2} I(m_a, T, \mu_Q)$$

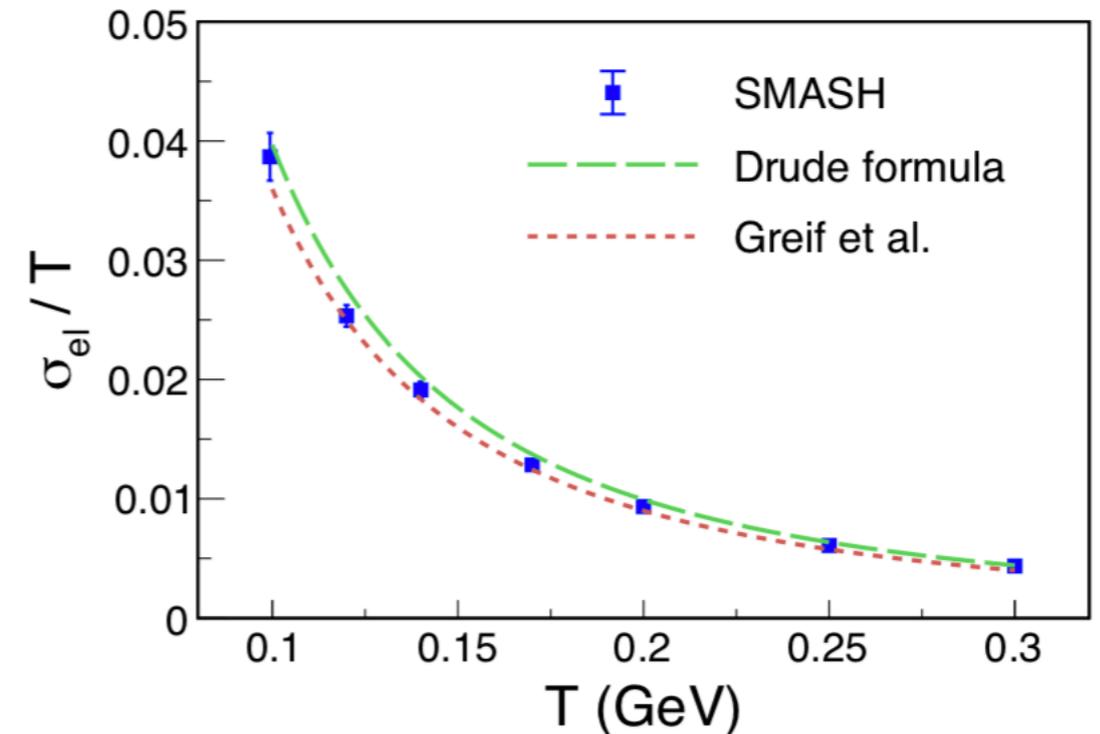
$$I(m_a, T, \mu_Q) = \int_0^\infty dp \frac{p^4}{m^2 + p^2} e^{-\frac{\sqrt{m_a^2 + p^2} - q_a e \mu_Q}{T}}$$



$C(0)V$ of massless gas with three species
 $q_a = \pm 1, 0$ constant cross section $\sigma_{tot} = 30mb$

Test Case: Massless gas

Massless gas with three species
 $q_a = +1, 0, -1$
 interacting via constant cross section
 $\sigma_{el} = 30mb$



- Drude Formula (classical kinetic theory)

$$\sigma_{el}^{Drude} = \frac{1}{2} \frac{2e^2}{3\sigma T} = \frac{3.9686 \cdot 10^{-4}}{T} (GeV^2)$$

Greif et al. Phys. Rev. D90, 094014 (2014)

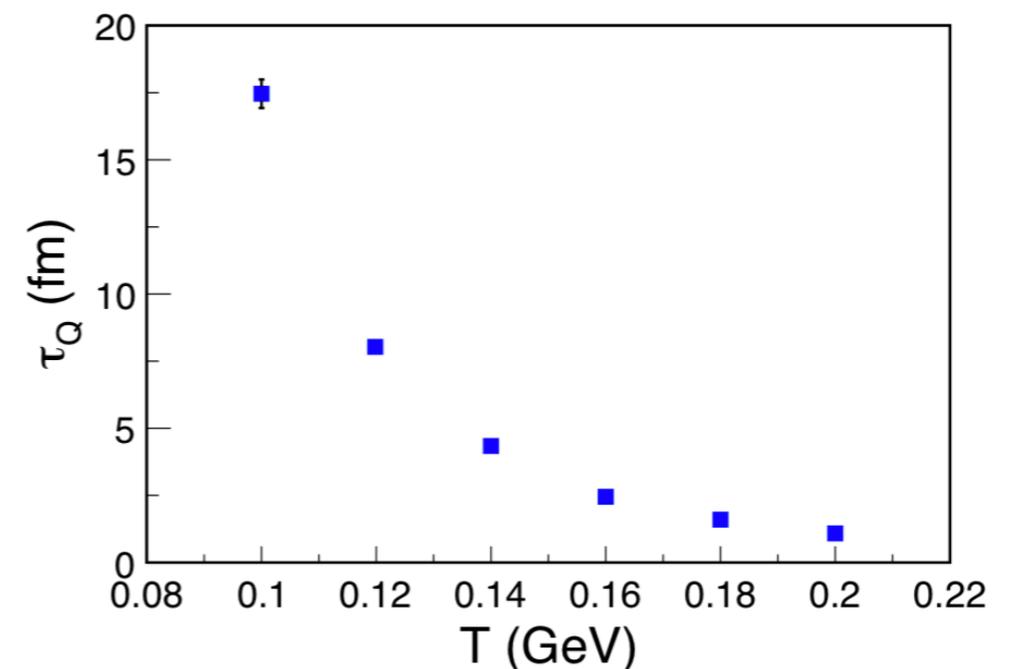
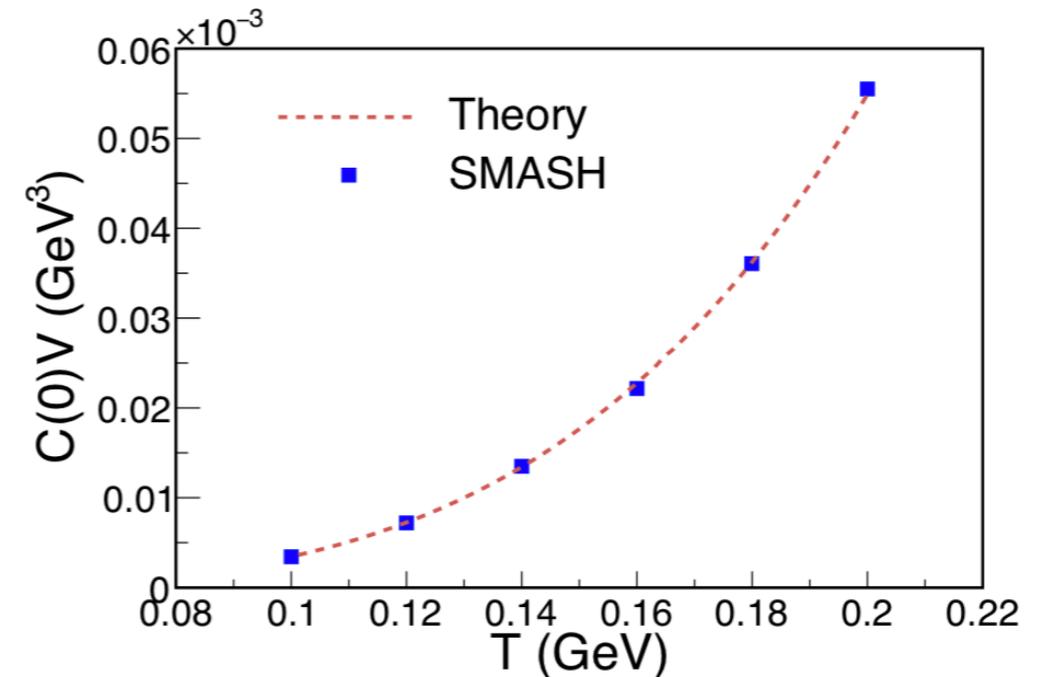
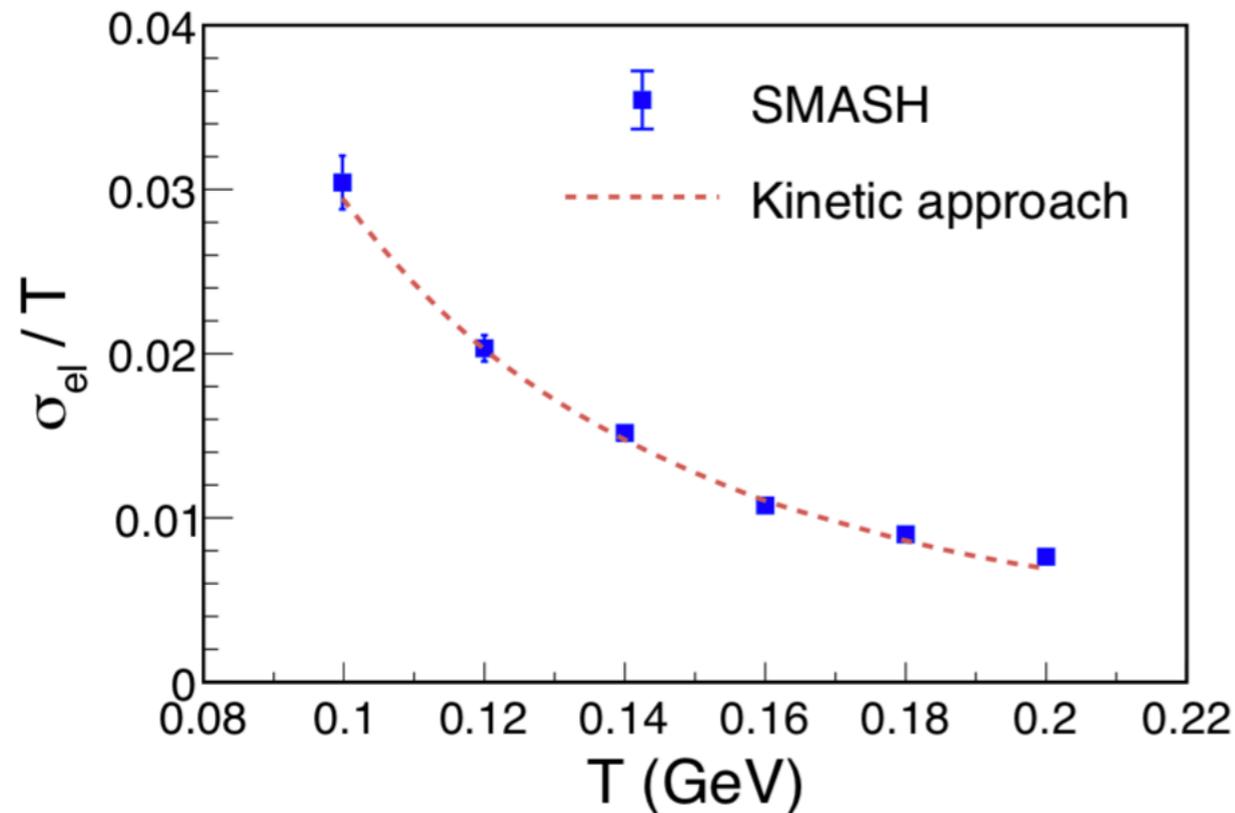
- Kinetic theory

$$\sigma_{el}^{Drude} = \frac{2/3}{3/13} \frac{3}{30} \frac{0.000832737}{T} = \frac{3.5967 \cdot 10^{-4}}{T} (GeV^2)$$

Greif et al. Phys. Rev. D93, 096012 (2016)

Test case: Hadron mixture

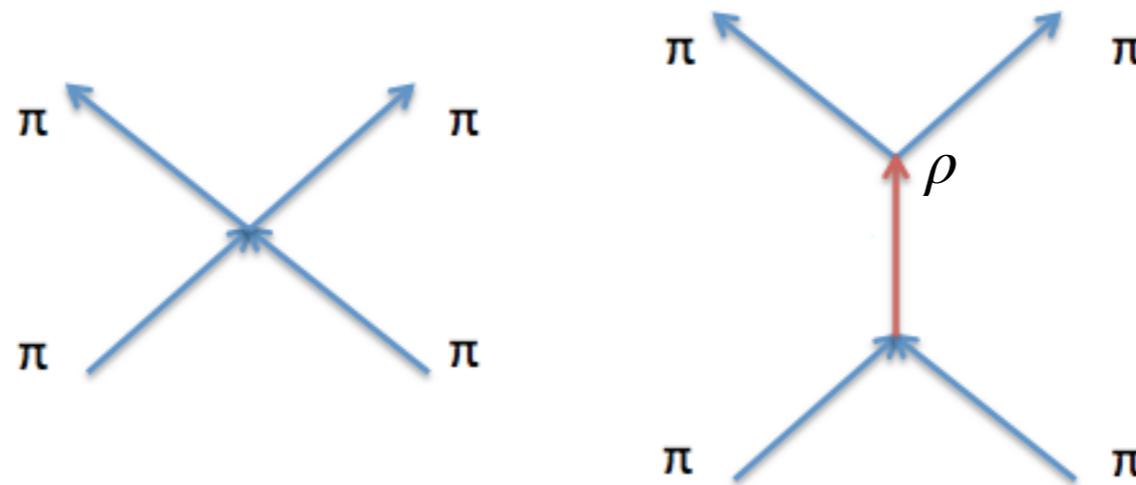
Three types of massive hadrons (π, K, N) interacting with a constant cross section
 $\sigma_{el} = 30mb$



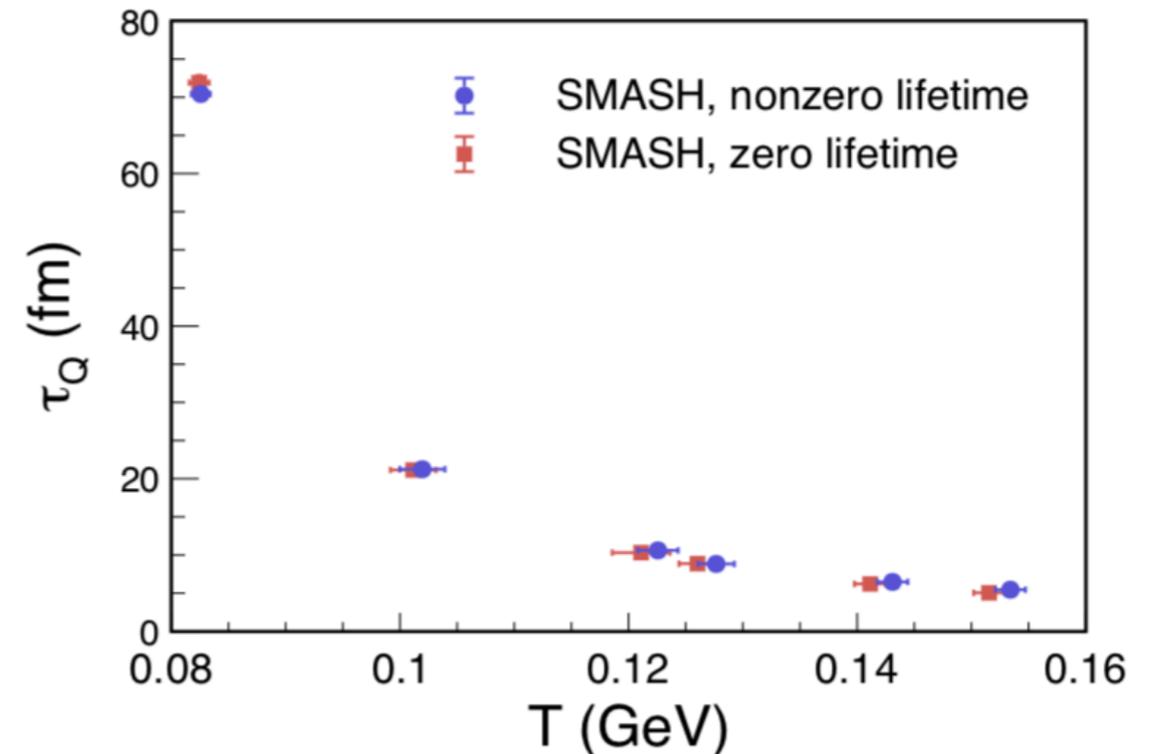
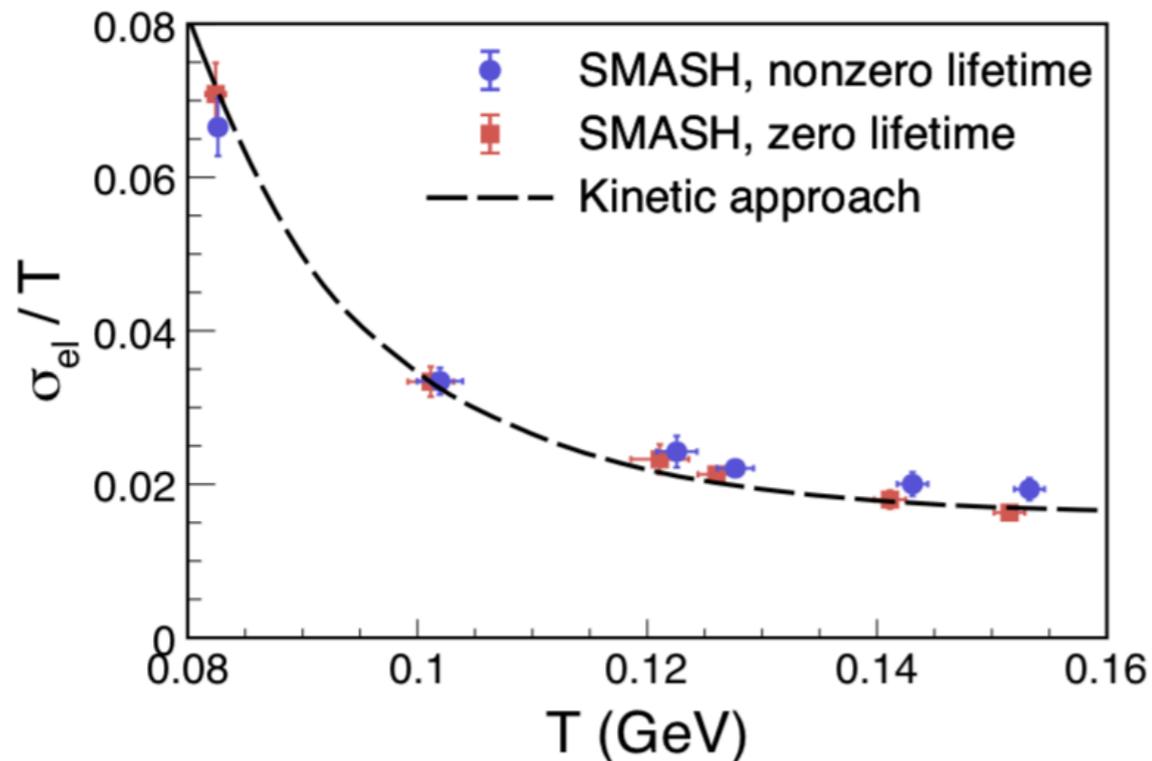
Greif et al.: Phys. Rev. D **96**, 059902

Test case: $\pi\rho$ - gas

- Pions interacting via energy dependent elastic cross section by forming a ρ meson.
- Analytic calculation only models the interaction with an energy dependent cross section, ρ does not propagate.
- How is the transport coefficient affected by resonance lifetime?



Test case: $\pi\rho$ - gas

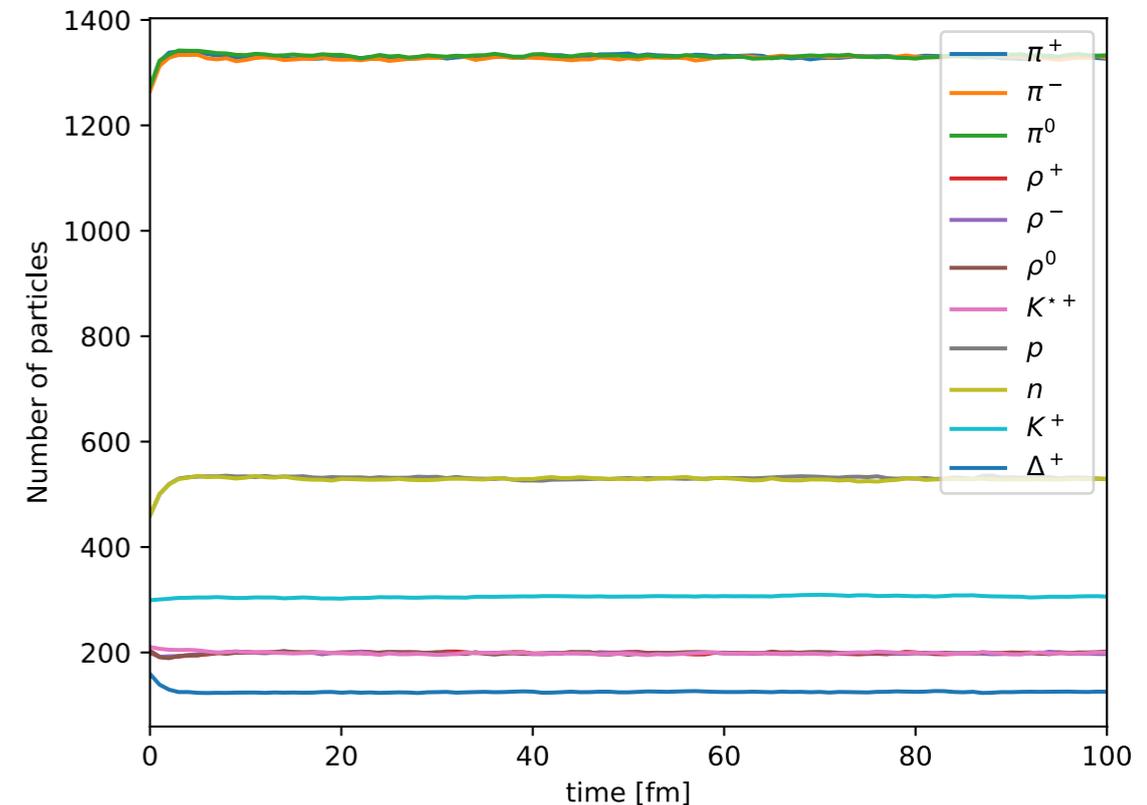


- Excellent agreement between analytic calculation and SMASH
- Resonance lifetime does not affect significantly the relaxation time τ_Q
- Difference in the electric conductivity comes from $C(0)$

Greif et al. :
Phys. Rev. D **93**, 096012 (2016)

Hadron gas

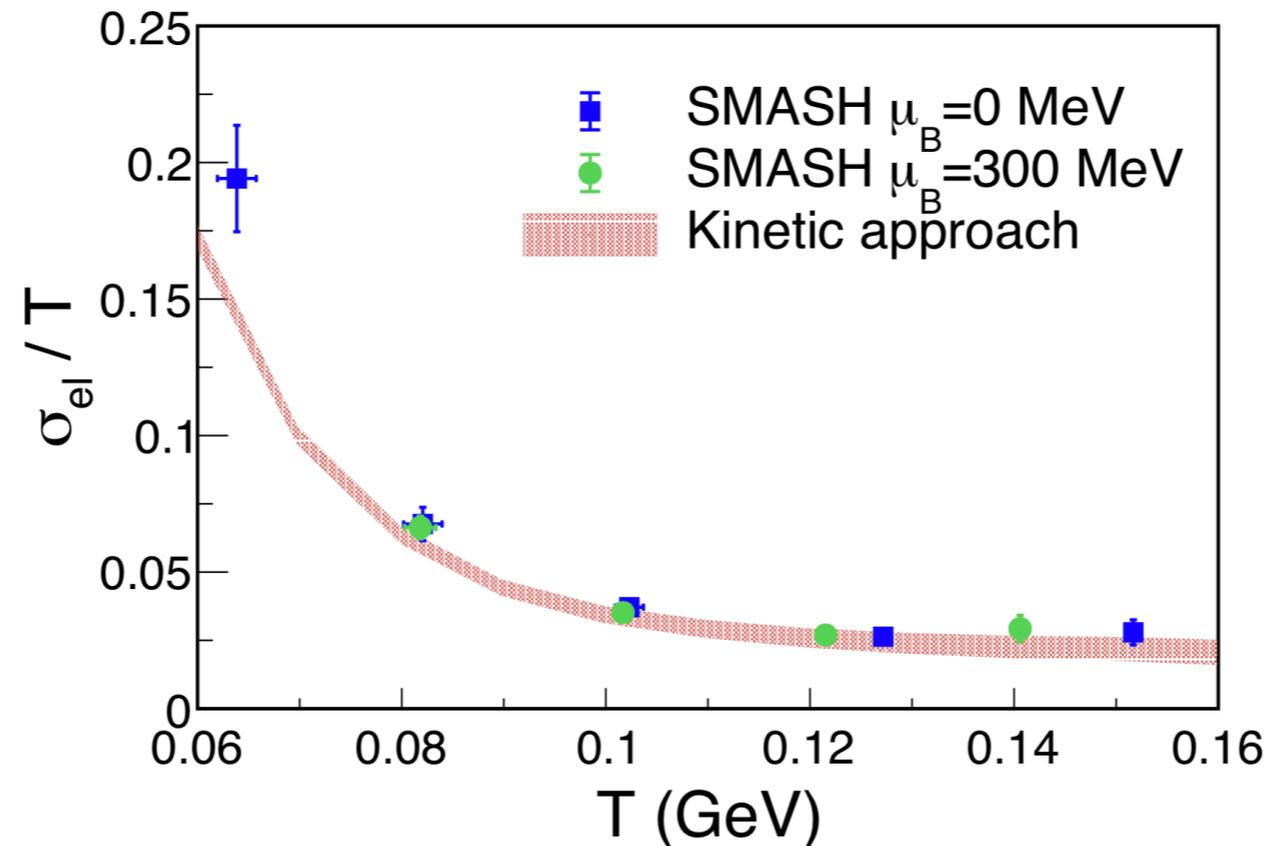
- Relatively simple hadron gas model with pions, kaons and nucleons, which represents a physical hadron gas below T_c .
- Both thermal and chemical equilibrium has to be reached for Green-Kubo method to be valid



Multiplicities of individual species in the hadron gas at $T = 140MeV$ and $\mu_B = 300MeV$

	π	K	N
π	ρ	K^\star	Δ
K		0	0
N			0

Hadron gas



- Slightly different result since the analytic calculation includes constant cross section + zero resonance lifetime
- Similar to the shear viscosity, baryochemical potential does not affect the electric conductivity as long as $\mu_B \ll m_N$

Greif et al.: Phys. Rev. D **96**, 059902 Rose et al. Phys. Rev. C **97**, 055204

Shear viscosity

- Linear transport coefficient which measures how momentum is transferred in transverse direction
- Same methodology as for the electric conductivity with different current
- Shear viscosity is calculated with

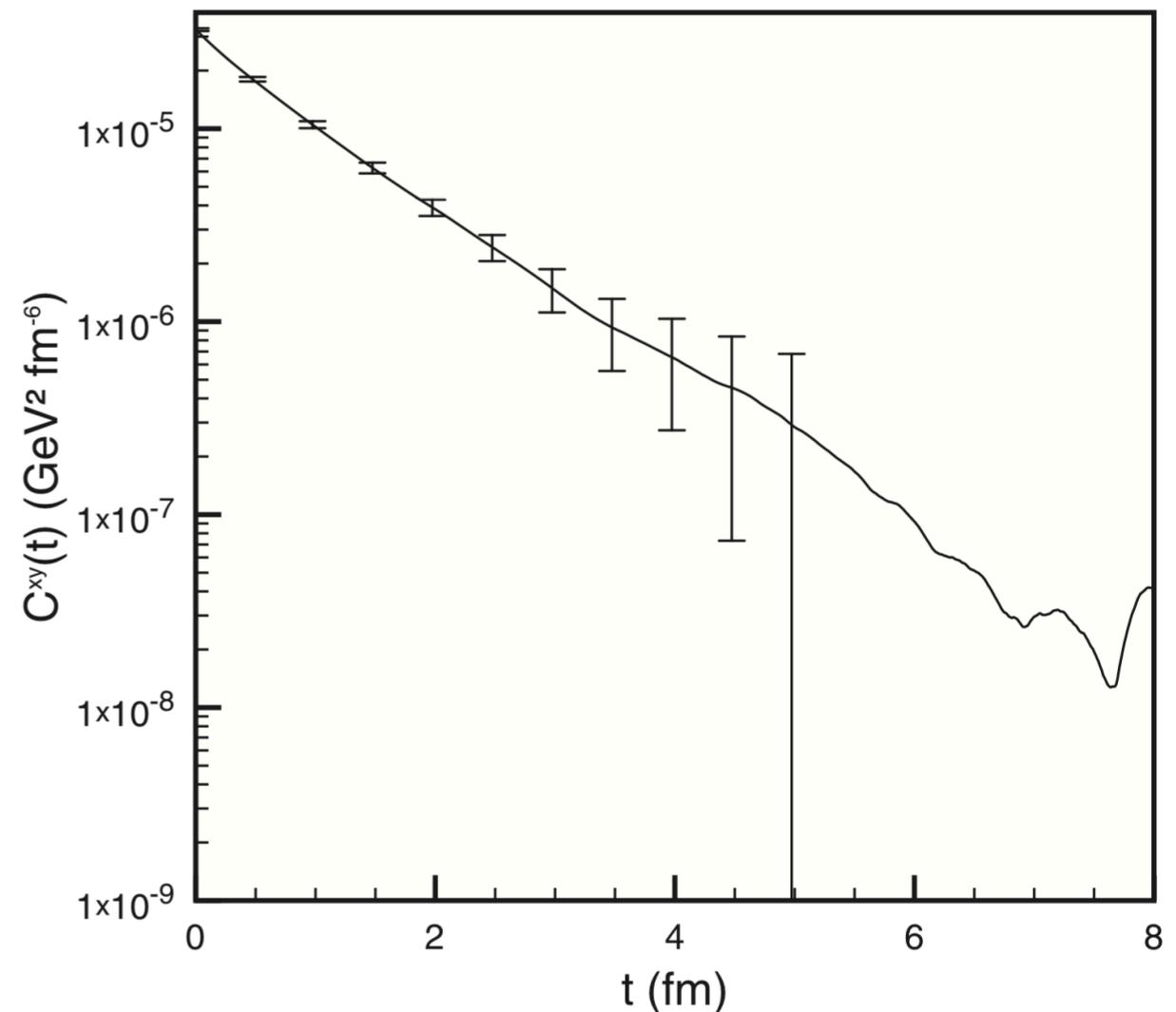
$$\eta = \frac{V}{T} \int_0^\infty \langle T^{xy}(t) T^{xy}(0) \rangle$$

with the energy-momentum tensor

$$T^{\mu\nu} = \frac{1}{V} \sum_i^{N_{part}} \frac{p_i^\mu p_i^{\nu}}{p_i^0}$$

and by fitting the correlation function with

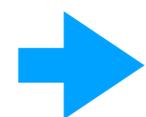
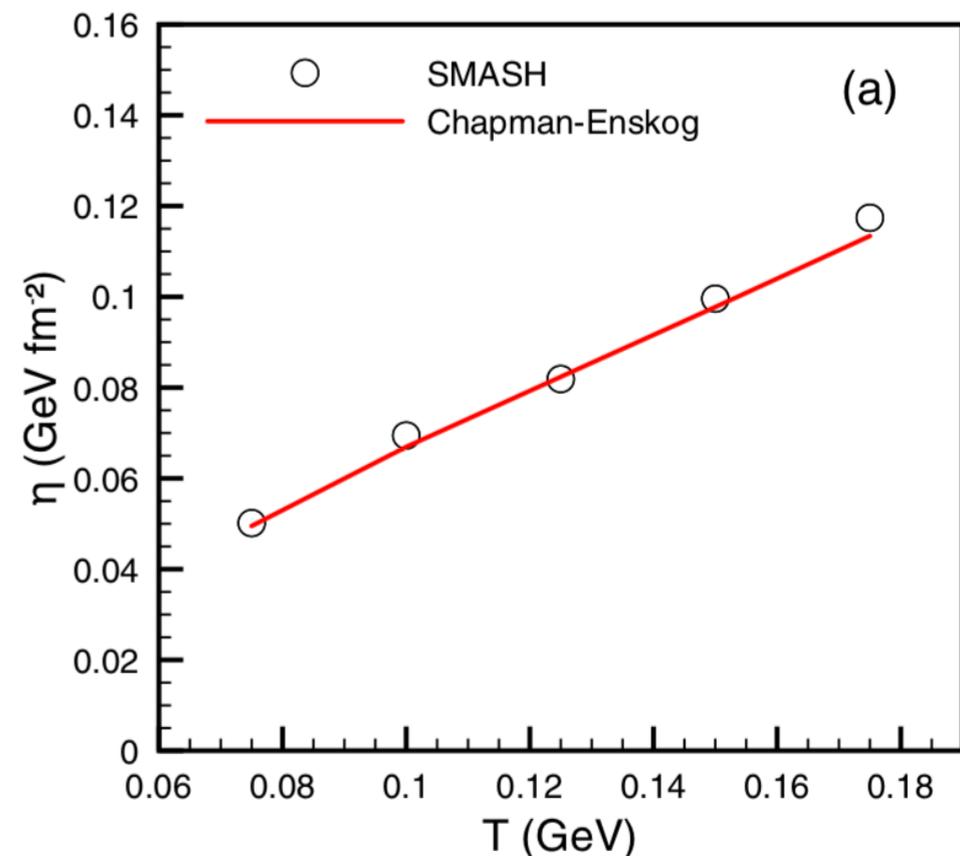
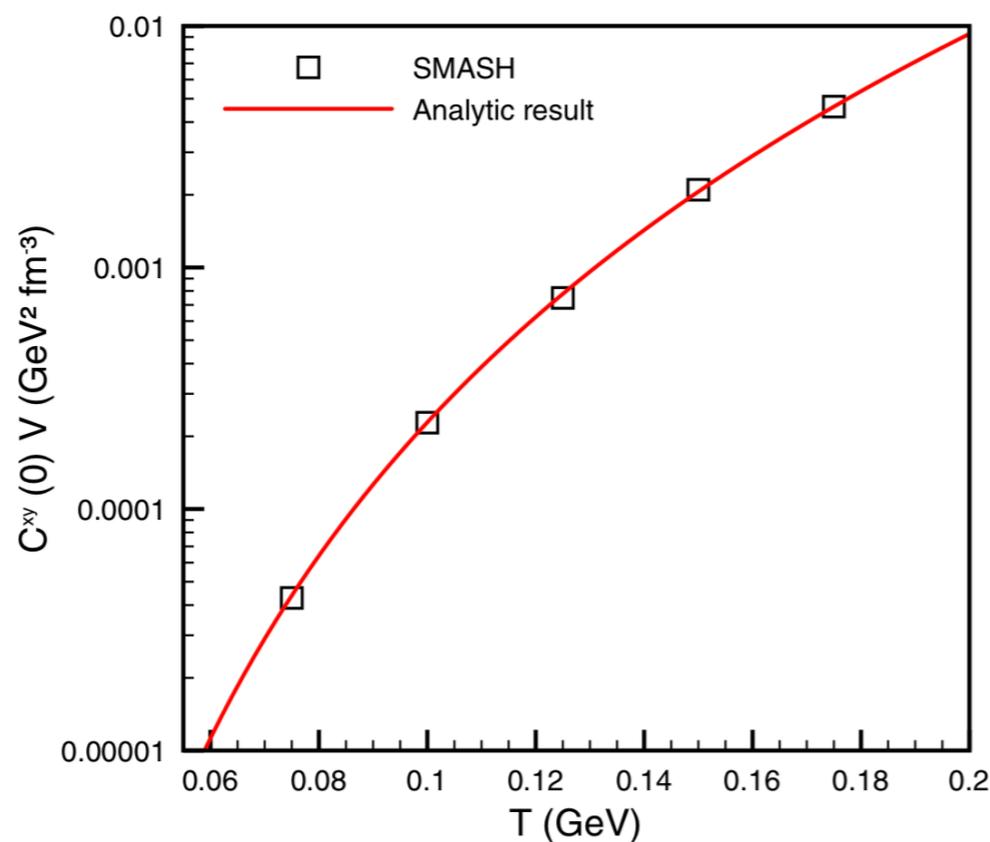
$$C^{xy} \equiv \langle T^{xy}(t) T^{xy}(0) \rangle = C^{xy}(0) e^{-t/\tau}$$



Rose et al. Phys. Rev. C **97**, 055204

Shear viscosity of a π - gas

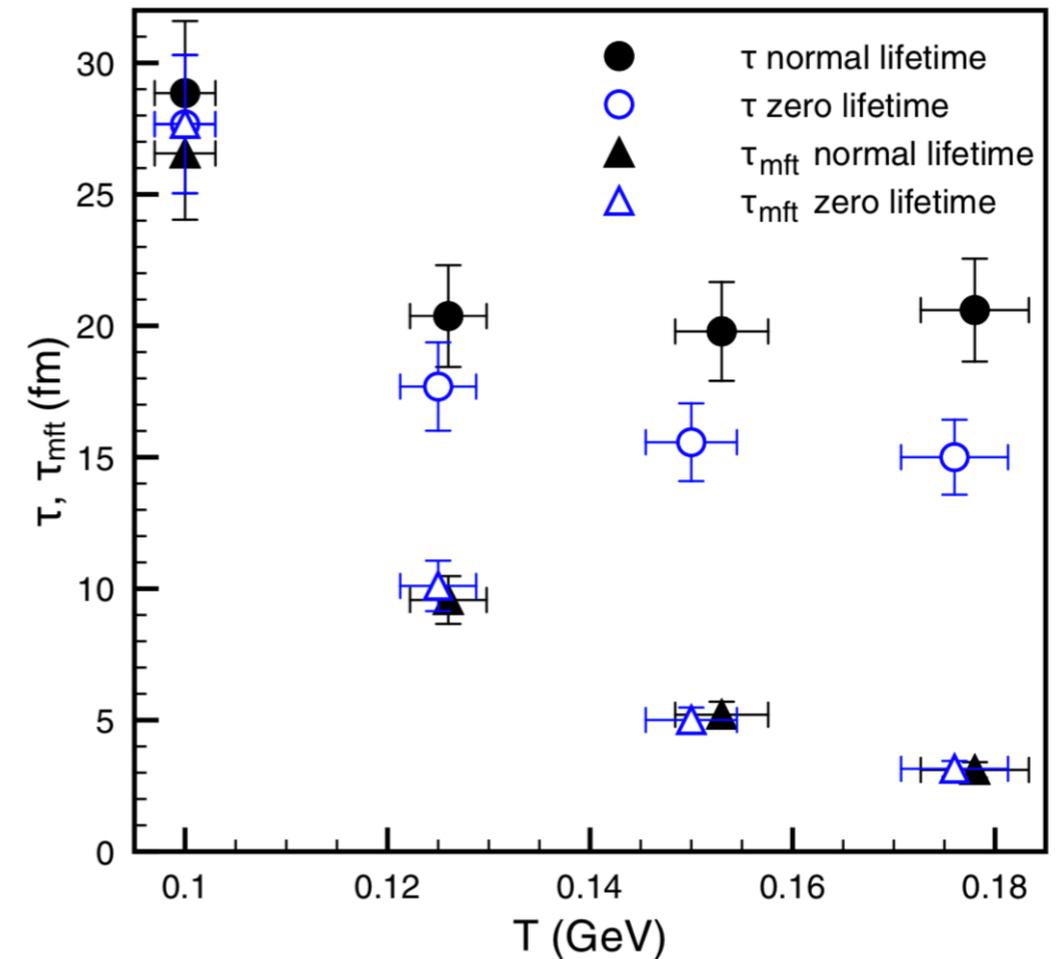
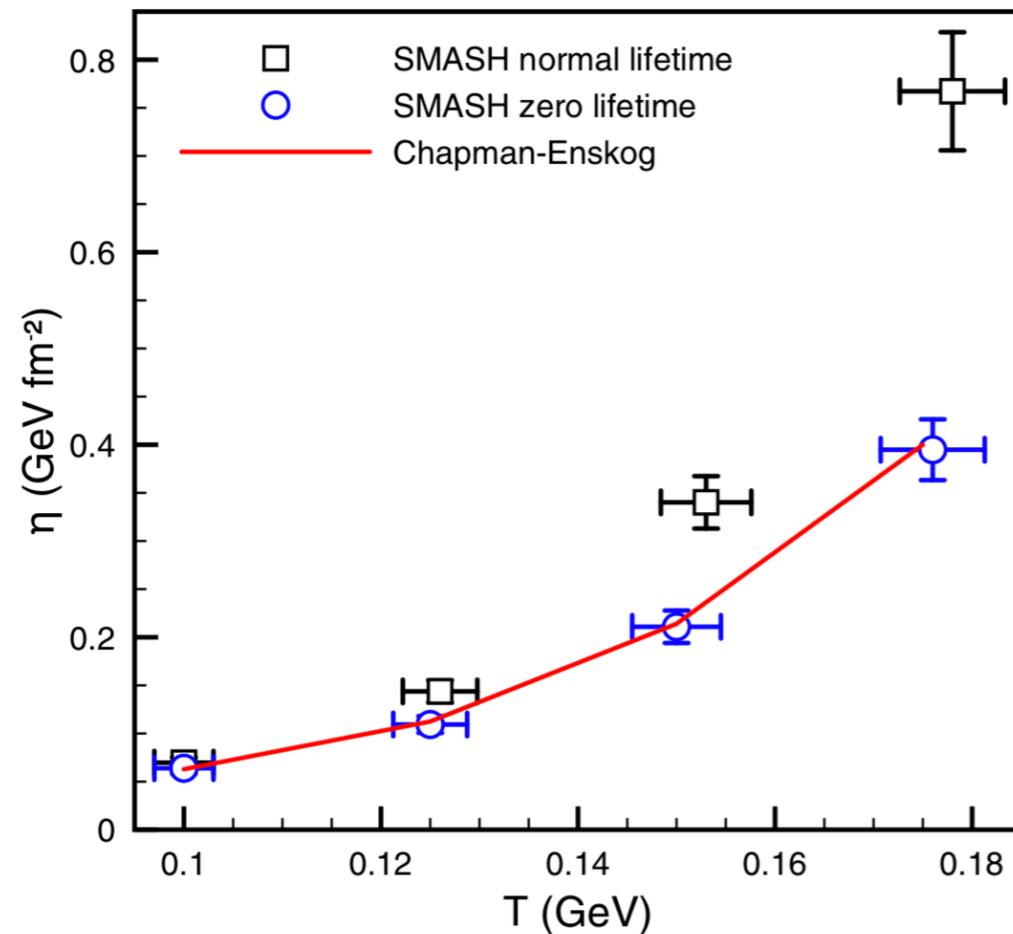
Test case: simple π - gas interacting with a constant cross section



Similarly to the electric conductivity there is a perfect agreement between SMASH and analytic calculations for the simple test cases!

Rose et al. Phys. Rev. C **97**, 055204

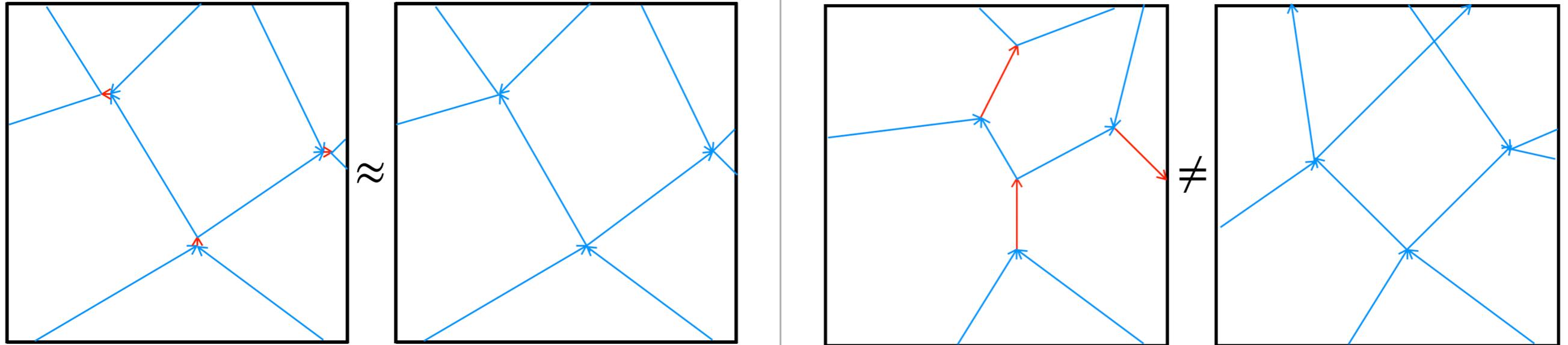
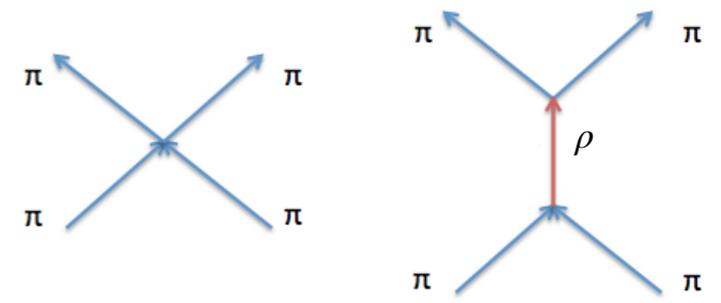
Shear viscosity of a $\pi\rho$ - gas



- Perfect agreement between SMASH with zero lifetime and analytic calculation
- Difference in τ between the zero and normal resonance lifetime cases

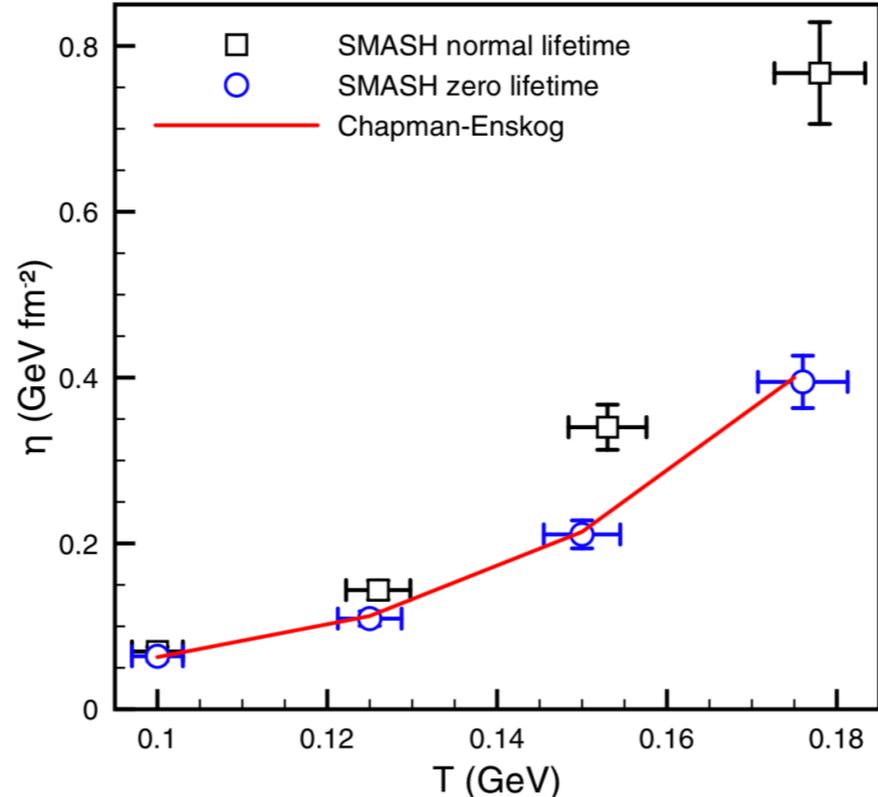
Rose et al. Phys. Rev. C **97**, 055204

Impact of lifetime of the resonance: Shear viscosity



Low temperature
 $\tau_{mft} \gg \tau_{res}$

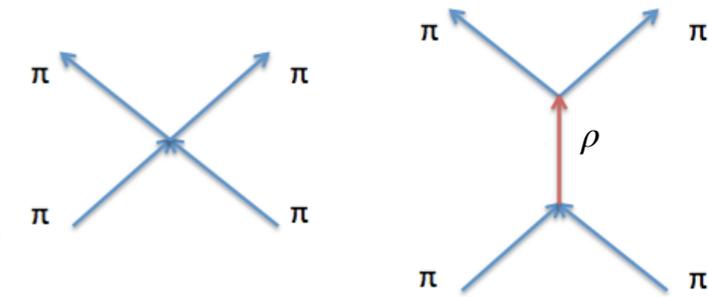
High temperature
 $\tau_{mft} \sim \tau_{res}$



τ_{mft} = mean free time
 τ_{res} = resonance lifetime

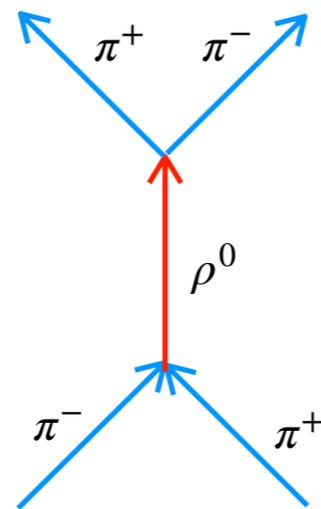
Rose et al. Phys. Rev. C **97**, 055204

Impact of lifetime of the resonance: electric conductivity



- There are several processes in which the equilibration has already happened at the first interaction
- Since $m_\rho \gg m_\pi$, the contribution of ρ mesons to current is small compared to that of pions
- Rate of formation/decays is not affected, when changing the resonance lifetime

Take for example the following reaction:



- Rho does not contribute to the electric current
- Flux already disappears in the first interaction
- value of resonance life time does not play a significant role

Conclusions and Outlook

▶ Conclusion

- Computed the electric conductivity of a hadron gas using the Green-Kubo formalism
- Found excellent agreement with analytic calculation
- Resonance lifetimes have to be taken into account when comparing transport coefficients from analytic and numerical calculations

▶ Outlook

- Add more species to the calculation
- Go to higher values of μ_B
- Compute diffusion coefficient matrix of conserved charges (S,Q,B) numerically within SMASH, as in [Greif et al. Phys. Rev. Lett. 120, 242301]

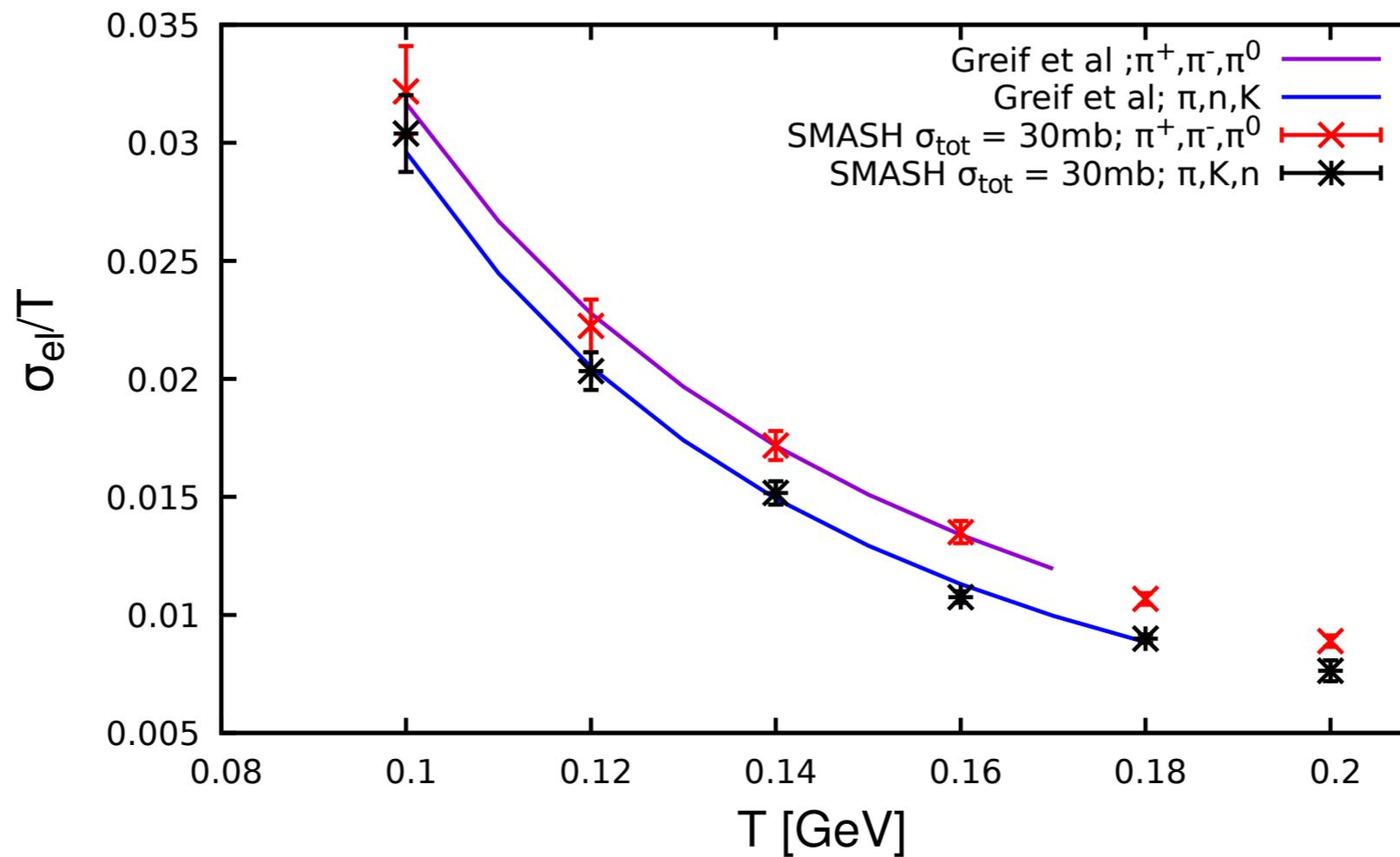
Backup slides

Hadron gas

Interaction channels from Greif et al.: Phys. Rev. D **96**, 059902

	π^+	π^-	π^0	K^+	K^-	K^0	\bar{K}^0	p	n	\bar{p}	\bar{n}
π^+	10	ρ	ρ	10	10	K^*	10	Δ	10	10	Δ
π^-		10	ρ	K^*	10	10	K^*	10	Δ	Δ	10
π^0			5	K^*	10	K^*	K^*	Δ	Δ	Δ	Δ
K^+				10	10	10	50	6	10	20	10
K^-					10	50	10	20	10	6	10
K^0						10	50	6	6	20	20
\bar{K}^0							10	8	20	6	6
p								20	20	100	20
n									20	20	100
\bar{p}										10	10
\bar{n}											10

Electric conductivity



$C(0) \pi\rho - gas$

