# Collective effects in O-O collisions from a hybrid approach

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**Transport Meeting** 

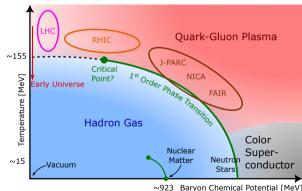
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Motivation for small collision systems Model Descriptions SMASH-vHLLE Hybrid Angantyr Results Initial State Nuclear Modification Factor Anisotropic Flow Conclusions

# Motivation for small collision systems

• At high temperature and pressure: transition from confined quarks (hadrons) to deconfined quark gluon plasma

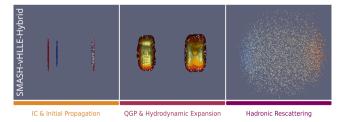
- Lattice QCD: smooth phase transition at  $\mu_b\approx 0$
- Where is the critical endpoint?
- **Goal:** Studying the limits of QGP formation in small collision systems



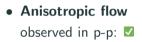
#### Signatures of QGP formation

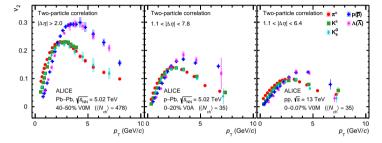
- QGP behaves like an almost ideal fluid
- $\rightarrow$  **Anisotropic Flow:** initial anisotropic density profile (almond shape)  $\rightarrow$  particles are pushed outwards anisotropically
- $\rightarrow$  Jet Quenching: high  $p_T$  particles from initial hard scatterings travel through the created medium
  - Data is well described by hybrid approaches:

Pre-Equilibrium, Hydrodynamics, hadronic evolution



#### Collectivity in small systems





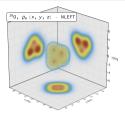
Observation of partonic flow in proton-proton and proton-nucleus collisions, ALICE Collaboration

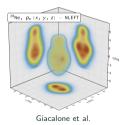
• Jet quenching observed in p-p: ×

$$R_{AA}(p_T) = \frac{1}{\langle N_{coll} \rangle} \frac{d^2 N_{AA}/dp_T dy}{d^2 N_{pp}/dp_T dy}\Big|_{y=0}$$

#### Intermediate small systems: O and Ne

- Bridging region between high multiplicity p-p and low multiplicity Pb-Pb
- Sharper energy gradients and higher event-by-event fluctuations compared to heavy-ion collisions
- Nuclear clustering:
  - $\rightarrow$  structure leaves imprint on energy profile
  - $\rightarrow$  small timescales allow us to take snapshots
- Configurations can be obtained by Nuclear Lattice Effective Field Theory (NLEFT)





# **Model Descriptions**

• SMASH solves the relativistic Boltzmann equation numerically

$$p^{\mu}\partial_{\mu}f_{i}(\vec{x},\vec{p})+m_{i}F^{\alpha}\partial^{p}_{\alpha}f_{i}(\vec{x},\vec{p})=C^{i}_{coll}$$

- For the initial conditions, a hypersurface of constant proper time  $au_0$  is defined
- Particles that cross this hypersurface get removed from the evolution and stored for the hydrodynamical evolution

#### Hydrodynamic Evolution - vHLLE

• The particles are smeared according to

$$\Delta P^{\alpha}_{ijk} = P^{\alpha} C \exp\left(-\frac{\Delta x_i^2 + \Delta y_j^2}{R_{\perp}^2} - \frac{\Delta \eta_k^2}{R_{\eta}^2} \gamma_{\eta}^2 \tau_0^2\right)$$

• The system is governed by the conservation of energy and momentum, as well as the net-baryon, net-charge and net-strangeness.

$$\partial_{\mu}T^{\mu\nu} = 0 \qquad \partial_{\mu}j^{\mu}_{c} = 0$$

• Equation of state needed!

 $\rightarrow$  Chiral mean field EoS from QCD that is fitted to HRG at lower temperatures

#### Particlization and Hadronic Afterburner

- When the system reaches  $\varepsilon_{switch}$ : Construct hypersurface of constant energy density
- Each surface element is particlized individually in 2 steps:
  - 1. Sample number of particles of each species using a Poisson distribution
  - 2. Particles' momenta are sampled according to the Cooper-Frye formula

$$\frac{\mathrm{d}N}{\mathrm{d}\vec{p}} = \frac{\mathrm{g}}{(2\pi)^3} \int_{\sigma} \left[ f_0(x,\vec{p}) + \delta f(x,\vec{p}) \right] \frac{p^{\mu}}{E} \mathrm{d}\sigma_{\mu}$$

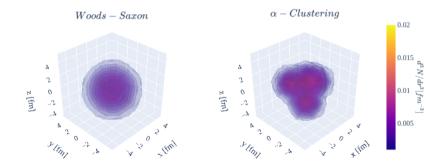
- $\rightarrow\,$  Conservation laws only satisfied on average!
  - Resulting particles are put into SMASH for hadronic evolution until kinetic freezeout

- $\bullet\,$  Extrapolation of  $\operatorname{Pythia}\,$  p-p events to A-A collisions
- Advanced Monte Carlo Glauber model to determine wounded nucleons
- $\rightarrow\,$  Sub-collisions are combined to obtain full heavy-ion event
- $\rightarrow~$  No collective effects

#### Implementation of $\alpha$ -Clustered Oxygen

- Sample 4 helium nuclei with Woods-Saxon distribution
- Place each on the vertex of a regular tetrahedron





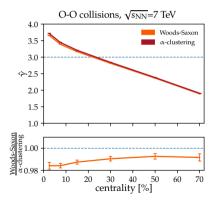
## Results

### Applicabillity of hydrodynamics

- Assumption of local thermal equilibrium not necessarily true especially in small systems!
- Assess degree of equilibration: Opacity

$$\hat{\gamma} = (5\eta/s)^{-1} \left(\frac{1}{a\pi} R \frac{dE_T}{d\eta_s}\right)^{\frac{1}{4}}$$

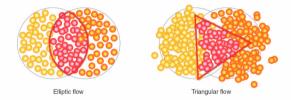
- Measure for the interaction rate in a medium
- Hydrodynamics found to be accurate to kinetic theory if  $\hat{\gamma}>3$
- $\rightarrow\,$  Applicable in central collisions up to 20% in  $\,$  O-O  $\,$



Werthman et al.

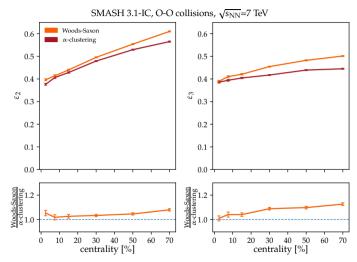
#### Eccentricity

- The Eccentricity describes the shape of the reaction zone in the transverse plane
- *ϵ*<sub>2</sub> and *ϵ*<sub>3</sub> are the ellipticity and triangularity, they measure how close to an ellipse or a triangle the shape is



$$|\epsilon_n| = \frac{\sqrt{\langle r^n \sin(n\varphi) \rangle^2 + \langle r^n \cos(n\varphi) \rangle^2}}{\langle r^n \rangle}$$

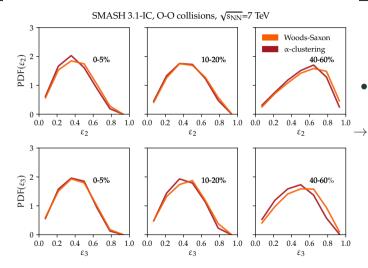
#### Eccentricity



- As expected, both ε<sub>2</sub> and ε<sub>3</sub> increase with centrality
- Woods–Saxon generates higher eccentricities across all centrality classes
- No significant ε<sub>3</sub>
   enhancement from clustering

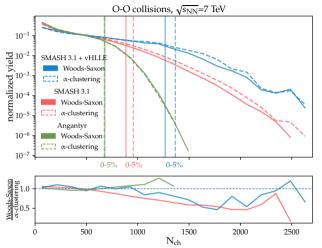
visible in central collisions

#### **Eccentricity Distribution**



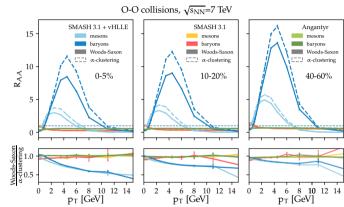
 Similar between centrality classes and distributions
 Random orientations and event-by-event fluctuations dilute geometric effects

#### **Centrality Selection**



- Centrality determined via final state charged particle multiplicity  $N_{ch}$  in rapidity region |y| < 2.5
- Results reflect entropy production of the 3 models:
  - $\rightarrow\,$  Angantyr has no hadronic rescatterings
  - $\rightarrow~$  Hybrid has viscous effects
- $\alpha$ -clustered configuration yields slightly higher multiplicities

#### **Nuclear Modification Factor**



$$R_{AA}(p_T) = \frac{1}{\langle N_{coll} \rangle} \frac{d^2 N_{AA}/dp_T dy}{d^2 N_{pp}/dp_T dy} \Big|_{y=0}$$

- Angantyr and SMASH transport nearly constant
- Hybrid shows expected result for thermal spectra over vacuum spectra

- Mass ordering: Baryons show higher  $R_{AA}$
- $\rightarrow$  Interpretation: radial flow pushes particles,  $\alpha\text{-clustered}$  configurations create denser medium

Goal: Fourier-decomposition of azimuthal particle distribution

$$\frac{dN}{d\phi} = \frac{1}{2\pi} \left( 1 + 2\sum_{n=1}^{\infty} v_n \cos(n(\phi - \Psi_n)) \right)$$

**Problem:** Non-flow  $\rightarrow$  correlations that do not originate from collective flow

- Sources: particle decays, Coulomb interactions, back-to-back jets
- ightarrow Mostly short ranged (small  $\Delta y$  and  $\Delta arphi$ )
- $\rightarrow$  Scaling with  $\frac{1}{N-1}$

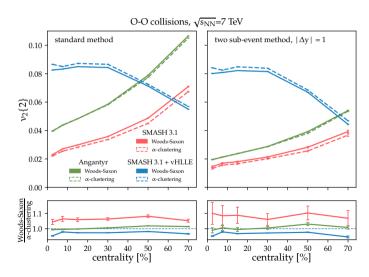
- Flow coefficients  $v_n$  are extracted from multi-particle azimuthal correlations
- Cumulants capture the genuinely correlated part of a distribution function:

$$egin{aligned} c_n\{2\} &= \langle \langle e^{in(\phi_1 - \phi_2)} 
angle 
angle \ c_n\{4\} &= \langle \langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} 
angle 
angle - 2 \langle \langle e^{in(\phi_1 - \phi_2)} 
angle 
angle^2 \end{aligned}$$

• Corresponding flow coefficients:

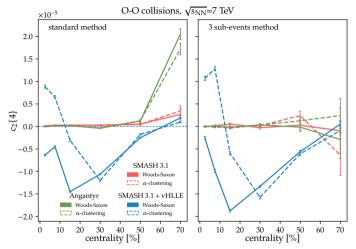
$$v_n\{2\} = \sqrt{c_n\{2\}}$$
  $v_n\{4\} = \sqrt[4]{-c_n\{4\}}$ 

 $\rightarrow$  k-particle cumulant does not contain contributions from lower order particle correlations



- Hybrid: highest v<sub>2</sub> in central collisions, decreasing in less central events
- SMASH and Angantyr: Opposite trend (Non-flow)
- Sub-event method can significantly reduce non-flow effects

## **4**-particle cumulant $c_2$ {4}



- Angantyr and SMASH transport close to zero (mostly)
- Incorrect sign prohibits calculation of v<sub>2</sub>{4}
- In general, Correct/Incorrect sign does not mean the flow is collective or not

$$v_n{4} = \sqrt[4]{-c_n{4}}$$

#### **Flow Fluctuations**

- Cumulants give estimates for the squared/quadrupled flow coefficients
- $\rightarrow\,$  Biased by event-by-event fluctuations
- $\rightarrow$  Sources: Multiplicity fluctuations, fluctuating initial geometry

Impact of flow fluctuations:

$$\langle v_n^2 \{2\} = \langle v_n \rangle^2 + \sigma_{v_n}^2$$
  $v_n^2 \{4\} \approx \langle v_n \rangle^2 - \sigma_{v_n}^2$ 

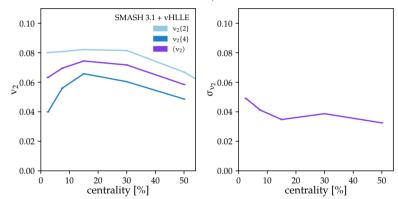
Unbiased flow estimate:

$$\langle v_n \rangle = \sqrt{\frac{v_n^2 \{2\} + v_n^2 \{4\}}{2}}$$

Flow fluctuations:

$$\sigma_{v_n} = \sqrt{\frac{v_n^2\{2\} - v_n^2\{4\}}{2}}.$$

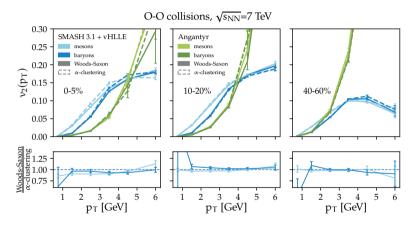
#### **Unbiased Flow Results**



O-O collisions,  $\sqrt{s_{NN}}$ =7 TeV

- Finally, the expected  $v_2$  picture in heavy-ion collisions emerges!
- $\bullet\,$  Flow fluctuations highest in central collisions  $\rightarrow\,$  Multiplicity fluctuations

#### **Differential Flow**



- Hybrid: Mass ordering at low p<sub>T</sub> visible, no baryon meson splitting at high p<sub>T</sub> → No individual parton description
- Angantyr: Non-flow goes through the roof once the number of particles decreases 22

## Conclusions

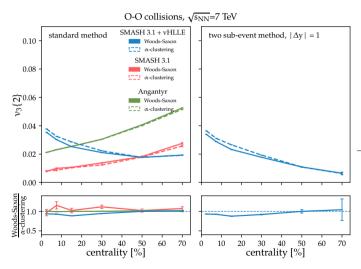
#### **Conclusion & Outlook**

- Hydrodynamic evolution causes big enhancement of  $v_2$  in central to mid-central collisions
- $\rightarrow\,$  Non-flow and flow fluctuations need to be considered
- $\rightarrow$  The  $\alpha\text{-clustered}$  configuration creates a denser medium, making a hydrodynamic description especially sensitve to the nuclear structure

#### **Outlook:**

- NLEFT configurations are available now
- Compare to Ne-Ne collisions
- $\rightarrow\,$  Light ions run at LHC in July 2025

# Backup



- Hybrid: similar picture
- SMASH and Angantyr: dominated by non-flow again
   → Flow from sub-event method

is imaginary!

