

Strangeness production in heavy-ion collisions at SIS energies

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FIAS Frankfurt Institute
for Advanced Studies



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SMASH transport approach

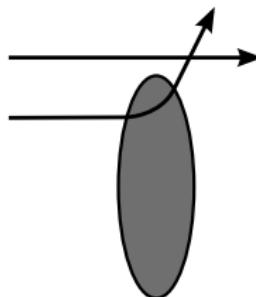
- ▶ **S**imulating **M**any **A**ccelerated **S**trongly-interacting **H**adrons
- ▶ 56 meson and 60 baryon species (+ anti particles)
= most of established hadrons from PDG (2017) made of *uds*
- ▶ Use cases: Nuclear collisions, infinite matter, afterburner for hydro
- ▶ Dileptons and photons
- ▶ Open source code will be published this year
- ▶ Test physics at SIS energies ($E_{\text{kin}} = 0.5 - 3.5A \text{ GeV}$), baseline for future NICA/FAIR predictions ($\approx 10A \text{ GeV}$)

J. Weil et al. In: *Phys. Rev.* C94.5 (2016), p. 054905. arXiv: [1606.06642 \[nucl-th\]](https://arxiv.org/abs/1606.06642)

Collision finding

- ▶ Geometric collision criterion (as used by UrQMD):

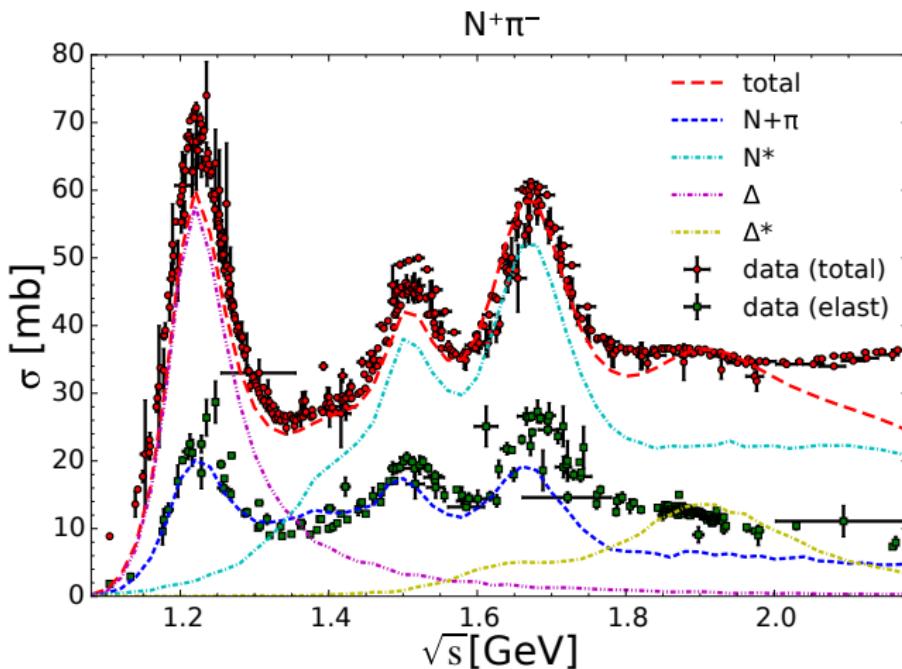
$$d_{\text{trans}} < d_{\text{int}} = \sqrt{\frac{\sigma_{\text{tot}}}{\pi}} \quad (1)$$



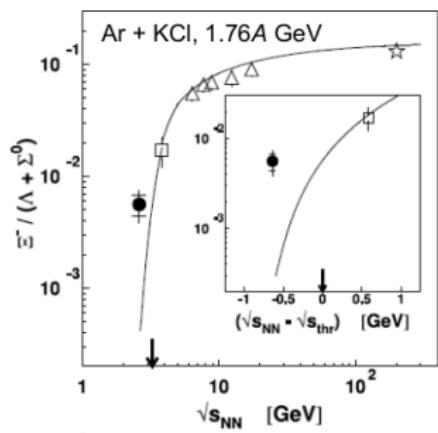
- ▶ Products of same reaction are forbidden to collide again
- ▶ Grid with cell size $\sqrt{\sigma_{\text{max}} / (\pi N_{\text{test}})}$ for collision finding

Cross section in SMASH

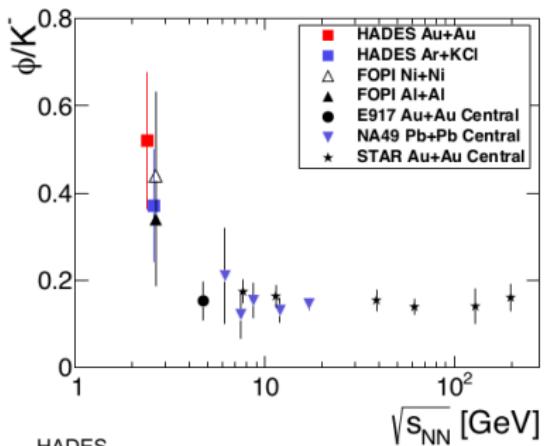
- ▶ Calculated from resonance masses, decay widths and branching ratios
- ▶ Parametrization of experimental data for non-resonant cross sections



Strangeness in heavy-ion collisions



HADES
PRL 103 (2009) 132301



HADES
arXiv:1703.08418

- ▶ Interesting probe for studying evolution of HIC
- ▶ High ϕ, Ξ measured by HADES \rightarrow sub-threshold strangeness enhancement
- ▶ $KN/\bar{K}N$ potentials? In-medium cross sections?
- ▶ Production mechanism in equilibrium (thermal model) and non-equilibrium (resonances)?

Strangeness production via resonances in SMASH

- ▶ Kaons and 11 kaonic resonances (+ anti particles)
- ▶ $\Lambda, \Sigma, \Xi, \Omega$ and 27 resonances (+ anti particles)
- ▶ Most important channels in heavy-ion collisions:
 - ▶ K^+ production ($Y \in \{\Lambda, \Sigma\}, B^* \in \{N^*, \Delta^*\}$):



- ▶ K^- production:



- ▶ No $BY \rightarrow NY^* \rightarrow NN\bar{K}$ channel

G. Graef et al. In: *Phys. Rev.* C90 (2014), p. 064909. arXiv: 1409.7954 [nucl-th]

Experimental branching ratios (2016-2017)

resonance	branching ratio $N^* \rightarrow \Lambda K$		
	PDG ¹	HADES ²	SMASH
$N(1650)$	5 – 15%	$7 \pm 4\%$	4%
$N(1710)$	5 – 25%	$15 \pm 10\%$	13%
$N(1720)$	4 – 5%	$8 \pm 7\%$	5%
$N(1875)$	> 0	$4 \pm 2\%$	2%
$N(1880)$		$2 \pm 1\%$	
$N(1895)$		$18 \pm 5\%$	
$N(1900)$	2 – 20%	$5 \pm 5\%$	2%
$N(1990)$			2%
$N(2080)$			0.5%
$N(2190)$	0.2 – 0.8%		0.8%
$N(2250)$	0		0.5%

¹C. Patrignani et al. *Chin. Phys. C* 40.10 (2016), p. 100001.

²R. Münzer et al. (2017). arXiv: 1703.01978 [nucl-ex].

Experimental branching ratios (2018)

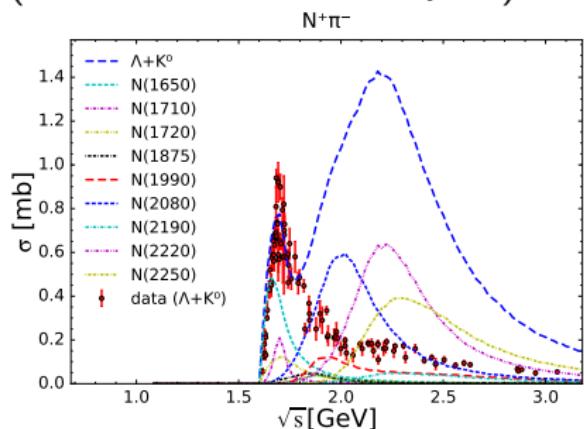
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$N(1875)$	> 0	$4 \pm 2\%$	2%
$N(1880)$	12 – 28%	$2 \pm 1\%$	
$N(1895)$	13 – 23%	$18 \pm 5\%$	
$N(1900)$	2 – 20%	$5 \pm 5\%$	2%
$N(1990)$			2%
$N(2060)$	> 0		
$N(2080)$			0.5%
$N(2100)$	> 0		
$N(2120)$	0		
$N(2190)$	$\Lambda K^* : 0.2 – 0.8\%$		0.8%
$N(2250)$	0		0.5%

¹ M. Tanabashi et al. (Particle Data Group), Phys. Rev. D 98, 030001 (2018)

Tuning branching ratios

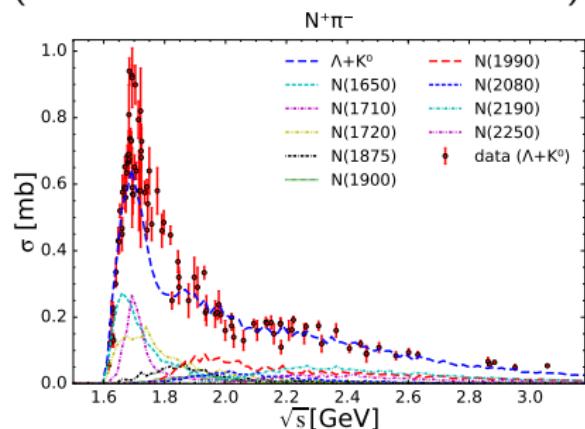
before

(based on PDG and UrQMD)



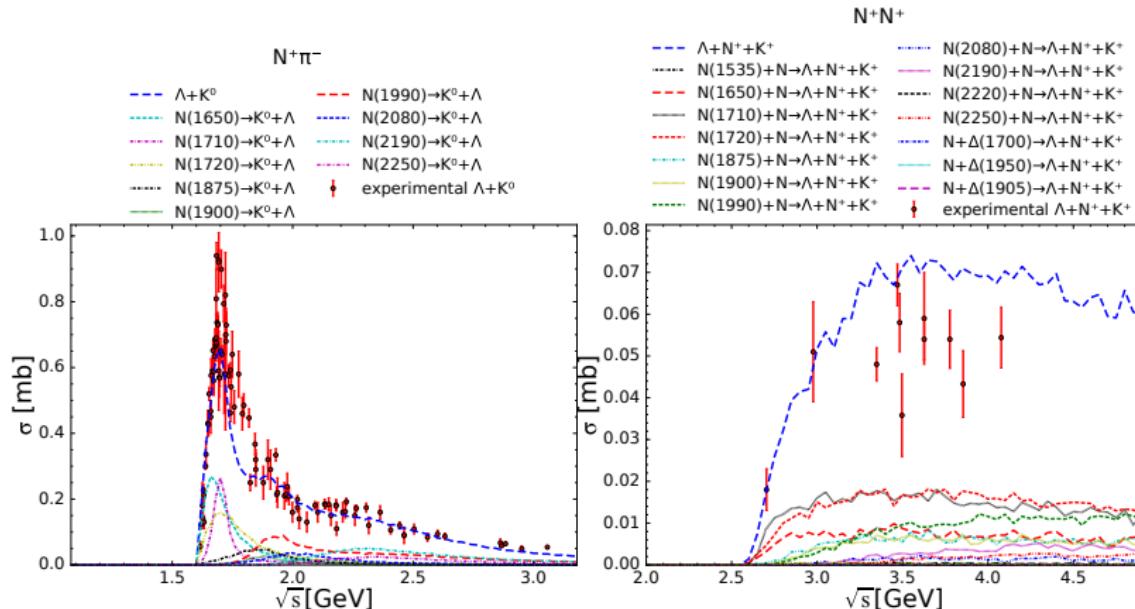
after

(tuned to exclusive cross section)



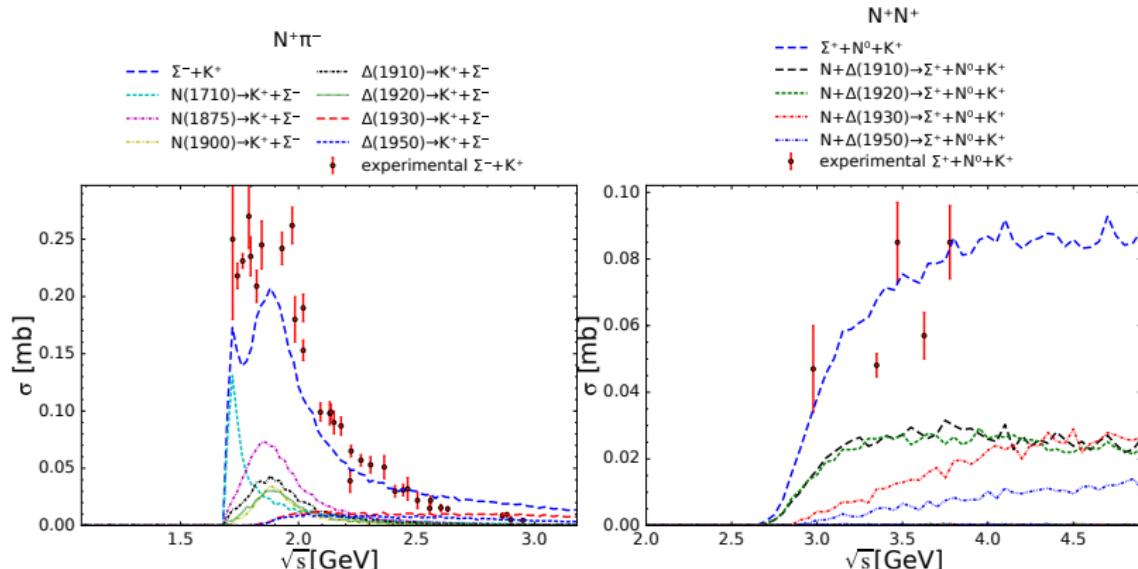
- ▶ Improve cross section by reducing $N(2080), N(2220), N(2250) \rightarrow \Lambda K, N\pi$

Λ production



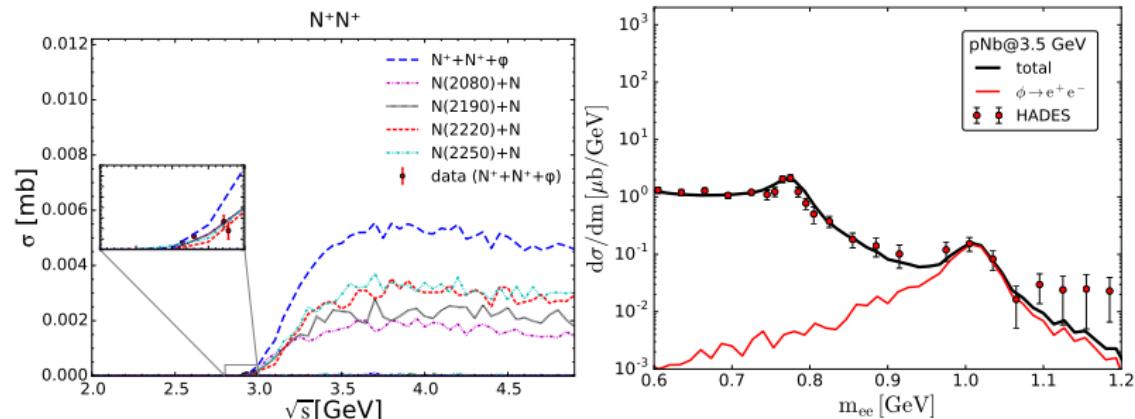
- ▶ Relevant branching ratios: $N^* \rightarrow \Lambda K, \pi N$

Σ production



- Relevant branching ratios: $N^*, \Delta^* \rightarrow \Sigma K, \pi N$

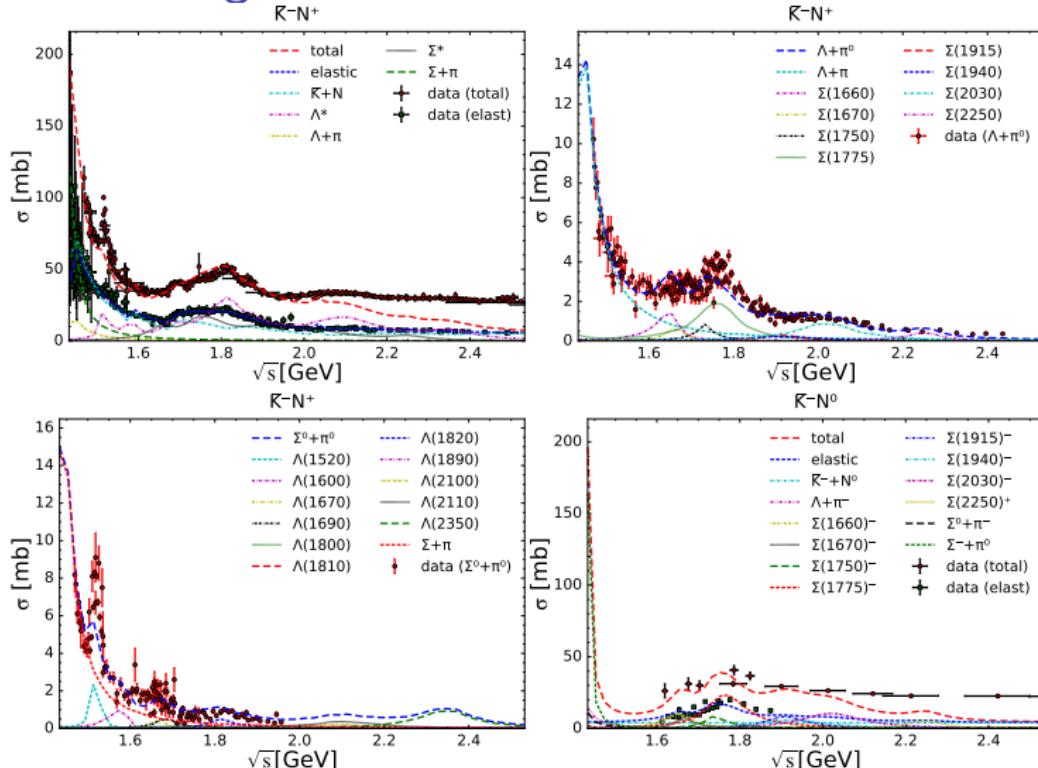
ϕ production



- ▶ $pp \rightarrow pp\phi \rightarrow ppK^+\bar{K}^-$ only measured at threshold
- ▶ ϕ production not well constrained by cross section
- ▶ Significant ϕ peak in p Nb dileptons
- ▶ Model ϕ production via $N^*(> 2000) \rightarrow N\phi$

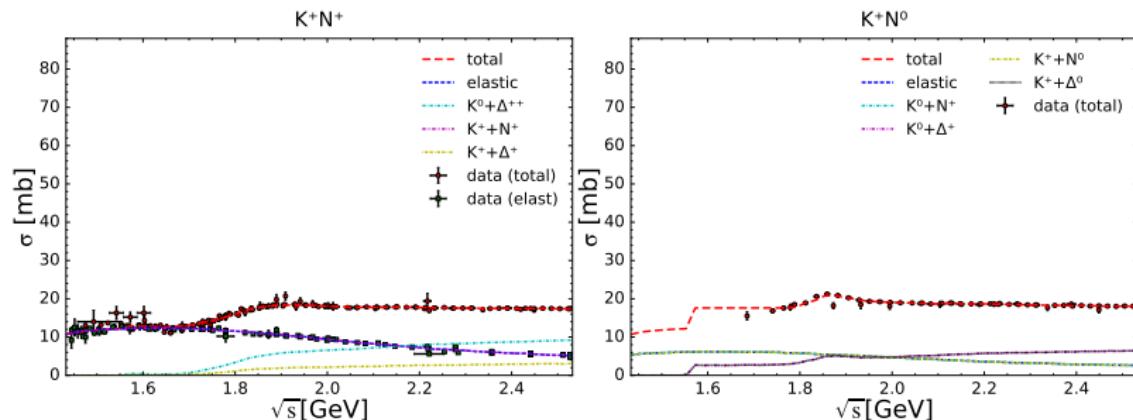
J. Steinheimer et al. In: *J. Phys. G* 43.1 (2016), p. 015104. arXiv:
1503.07305 [nucl-th]

$\bar{K}N$ scattering



- ▶ γ^* contribution + elastic background + charge exchange + strangeness exchange (like UrQMD)

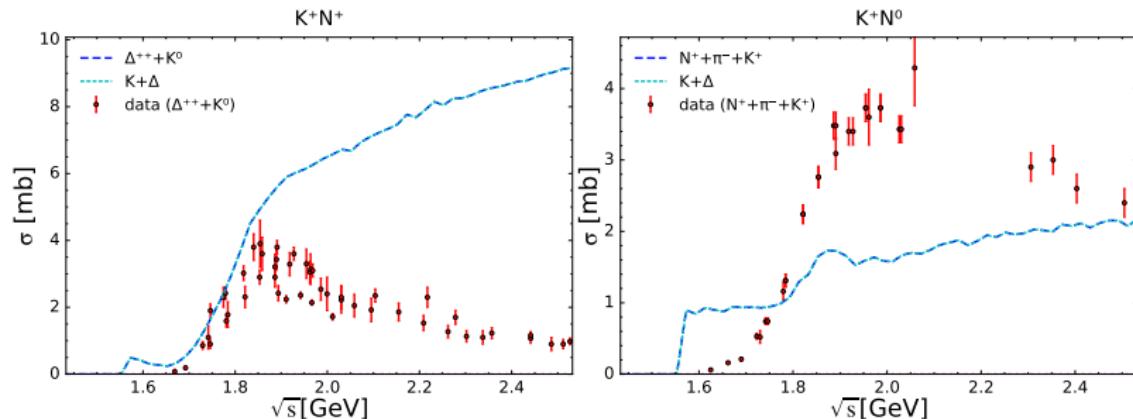
KN scattering



- ▶ K rescattering significantly affects kinematics³
- ▶ Elastic background + charge exchange + single-pion production (like GiBUU)
- ▶ Not designed to describe exclusive cross sections

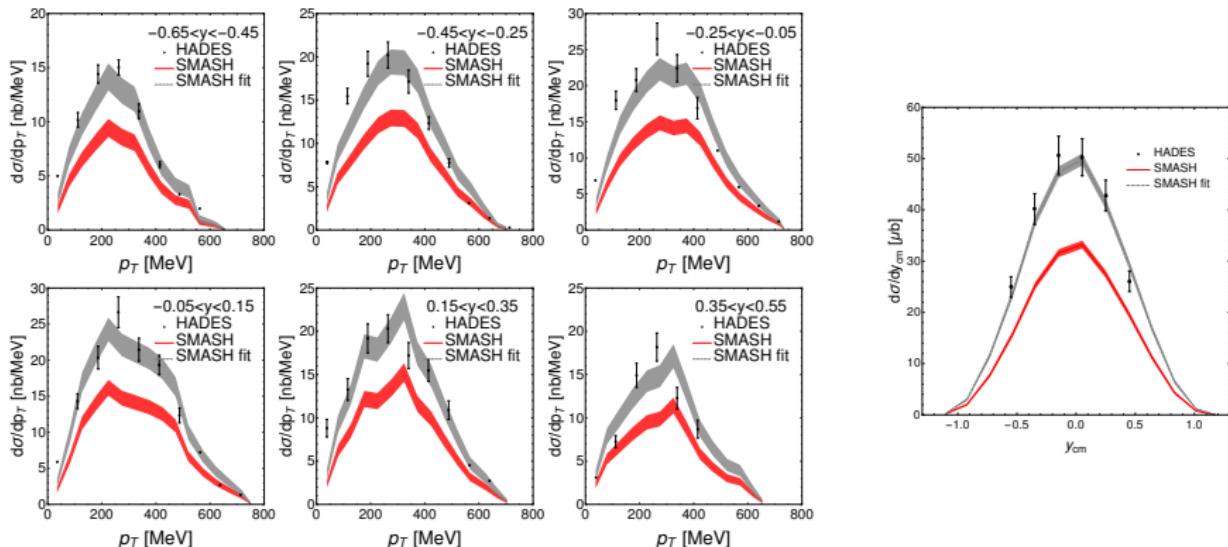
³C. Hartnack et al. *Phys. Rept.* 510 (2012), pp. 119–200. arXiv: 1106.2083 [nucl-th].

KN exclusive cross section



- ▶ $K^+ p \rightarrow \Delta^{++} K^0$: Not enough pions in the final state?
- ▶ $K^+ n \rightarrow p \pi^- K^+$: Not well reproduced
- ▶ Single-pion production too simplistic
- ▶ Probably not important for kaon rescattering

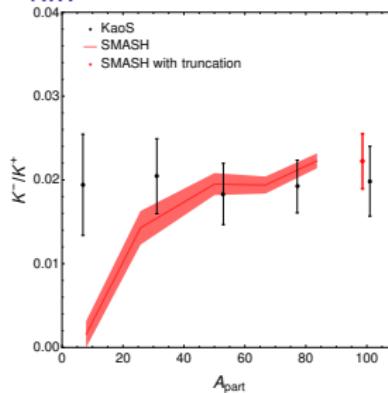
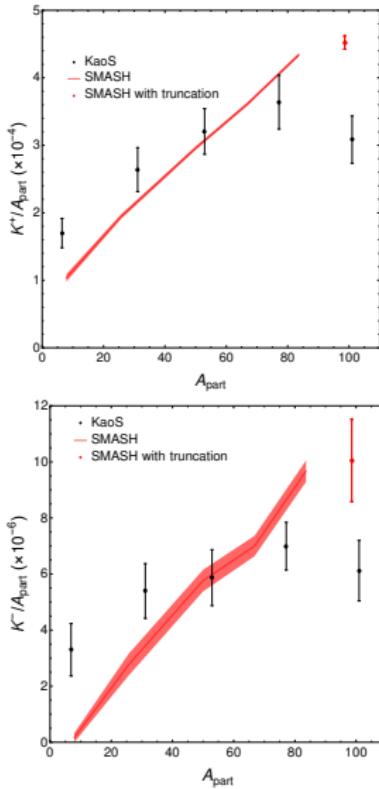
K_S^0 production in proton-proton at $E_{\text{kin}} = 3.5 \text{ GeV}$



G. Agakishiev et al. In: *Phys. Rev. C90* (2014), p. 054906. arXiv: 1404.7011 [nucl-ex]

- ▶ Assuming long-short symmetry: $\sigma_{K_S^0} = 0.5\sigma_{K^0} + 0.5\sigma_{\bar{K}^0}$
- ▶ Cross section too low by factor 1.5 (red), grey lines scaled accordingly (similar to GiBUU)

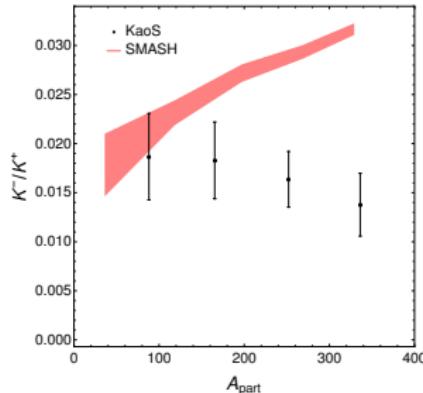
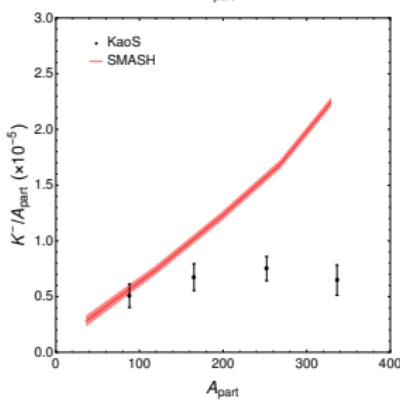
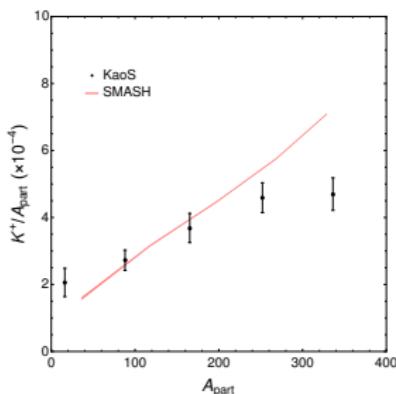
K^\pm production in ^{58}Ni - ^{58}Ni at $E_{\text{kin}} = 1.5A \text{ GeV}$



- Decent agreement with KaoS for ratio
- Absolute yields reasonable, but trend different (similar to IQMD)

A. Forster et al. In: *J. Phys.* G31.6 (2005), S693–S700. arXiv: nucl-ex/0411045
[nucl-ex]

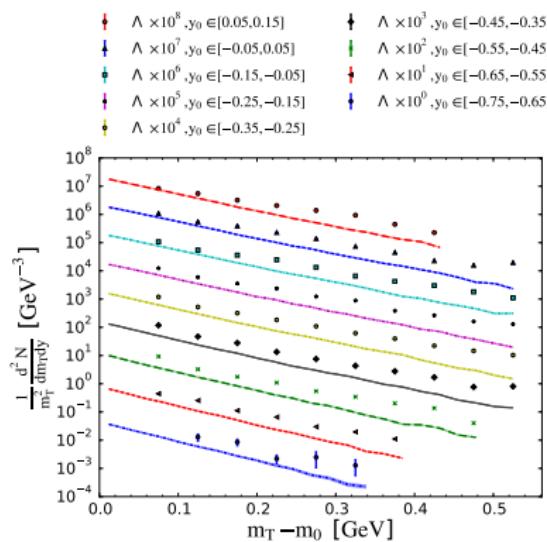
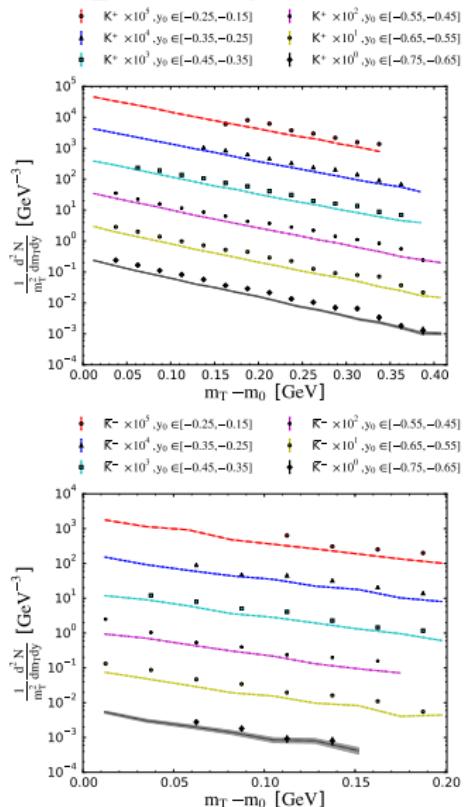
K^\pm production in gold-gold at $E_{\text{kin}} = 1.5A \text{ GeV}$



- ▶ Overestimation for high number of participants, more so for K^- than K^+
- ▶ Similar to IQMD

A. Forster et al. In: *J. Phys.* G31.6 (2005), S693–S700. arXiv: nucl-ex/0411045
[nucl-ex]

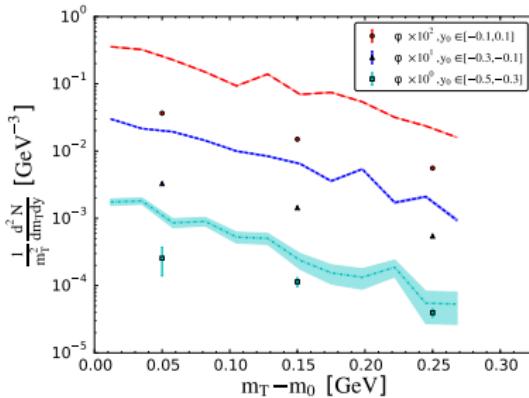
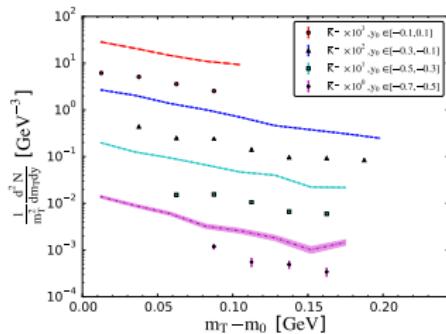
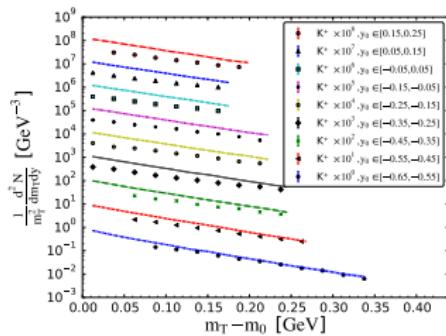
Strangeness production in ArKCl at $E_{\text{kin}} = 1.76A$ GeV



- ▶ HADES ArKCL compared to SMASH CaCa
- ▶ Kaons slightly and Λ significantly underestimated with wrong slope

G. Agakishiev et al. In: *Eur. Phys. J. A* 47 (2011), p. 21. arXiv: 1010.1675
[nucl-ex]

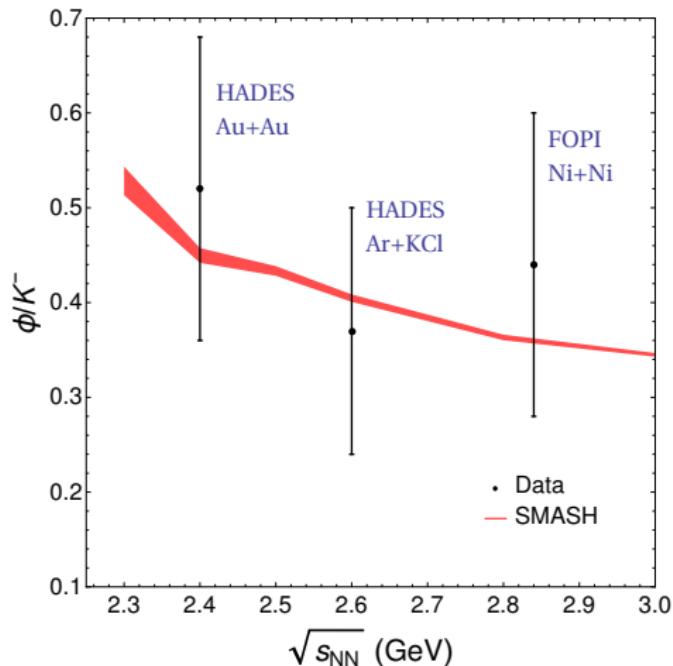
Strangeness production in AuAu at $E_{\text{kin}} = 1.23 \text{ GeV}$



- ▶ HADES AuAu compared to SMASH
- ▶ K^+ overestimated at midrapidities, K^- and ϕ at all rapidities

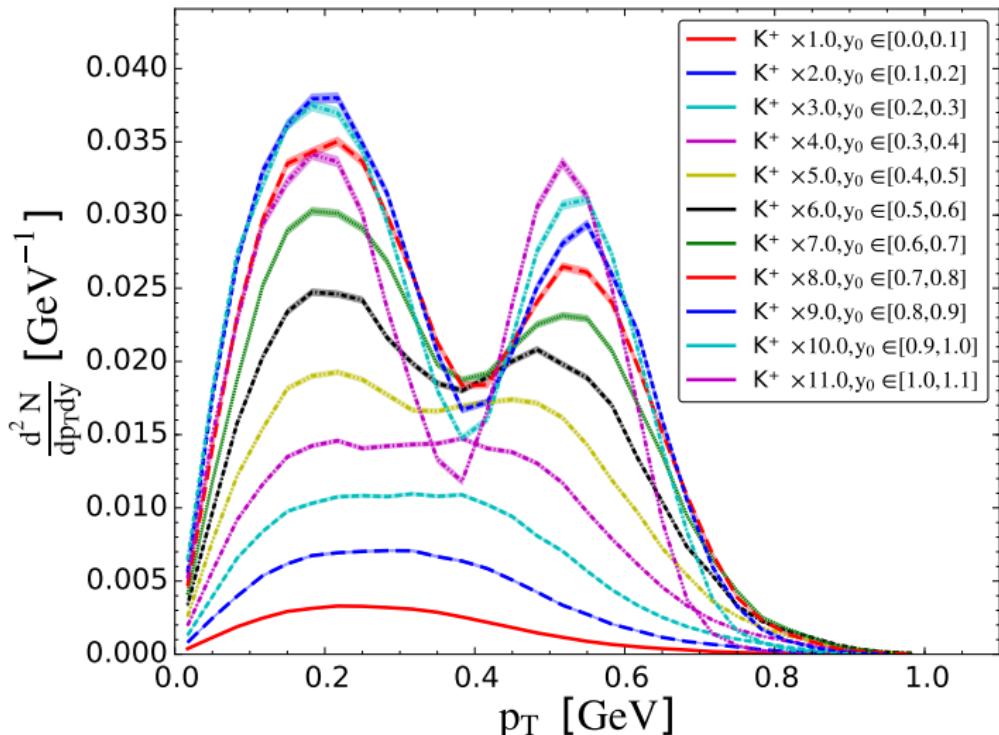
J. Adamczewski-Musch et al. In: *Phys. Lett. B* 778 (2018), pp. 403–407.
arXiv: 1703.08418 [nucl-ex]

ϕ/K^- ratio in AuAu



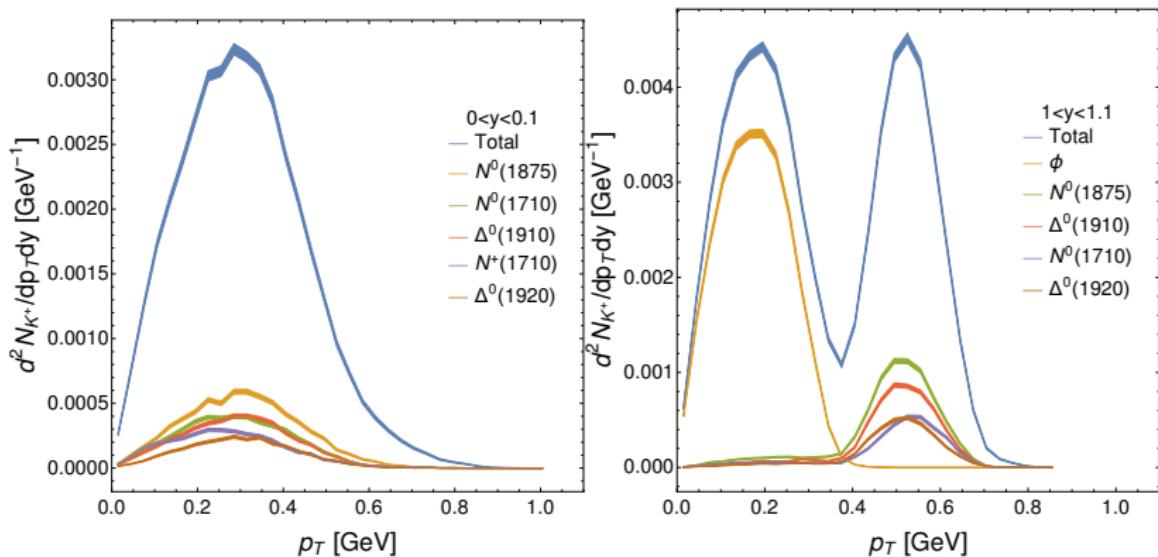
- ▶ SMASH AuAu compared to several systems
- ▶ Ratio well reproduced

K^+ production in $\pi^- C$ at $E_{\text{kin}} = 1.7 \text{ GeV}$



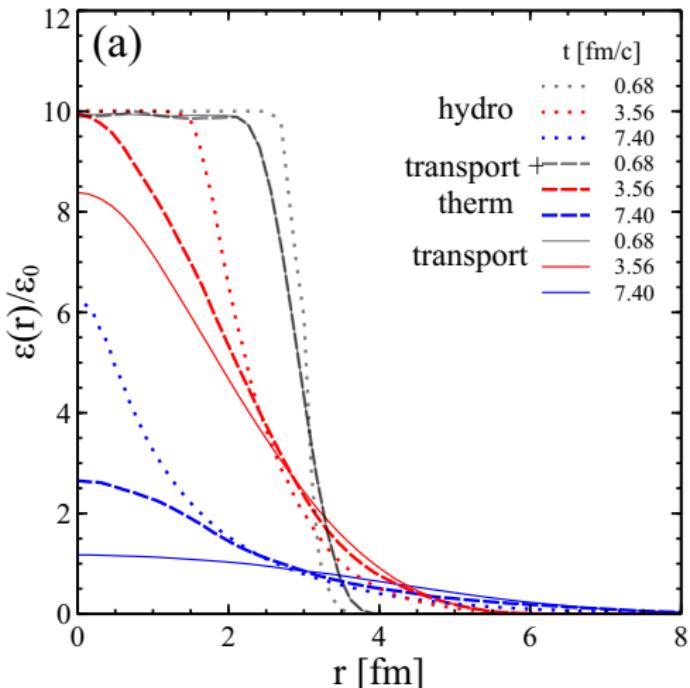
- ▶ Measured by HADES (not yet public)
- ▶ SMASH: Unphysical second peak at high rapidity

Contribution of resonances to pion beam



- ▶ Second peak from heavy resonances
- ▶ Not enough particles in final state of B^* decay?

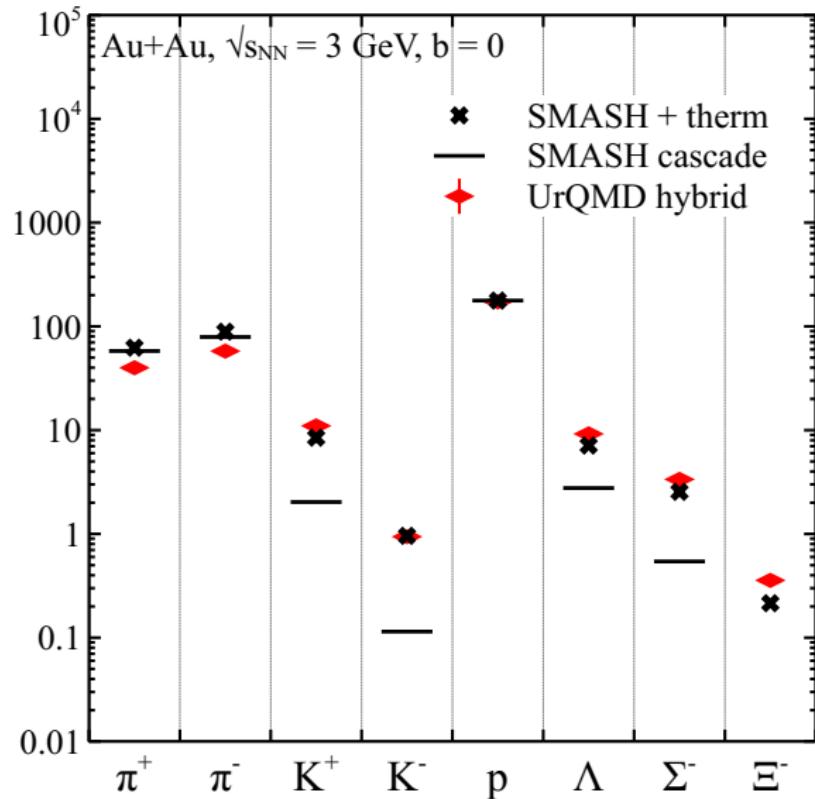
Particle production with forced thermalization



- ▶ Force thermalization in regions of high density by resampling particles
- ▶ Local, not global
- ▶ Effective many-particle scattering
- ▶ Similar to hydro-hybrid model, but more dynamic

D. Oliinychenko et al. In: *J. Phys. G* 44.3 (2017), p. 034001. arXiv: 1609.01087 [nucl-th]

Forced canonical thermalization vs. cascade + hydro

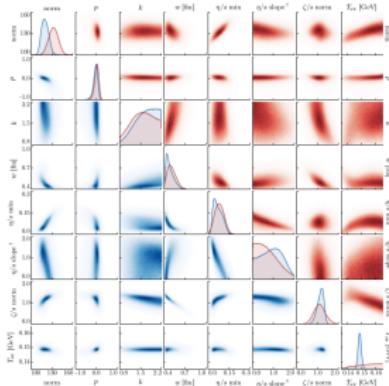


- ▶ Strangeness enhancement comparable to hybrid approach

Conclusion

- ▶ Elementary K, Λ, Σ, ϕ production at low energies can be reasonably modeled via resonances
- ▶ Dilepton data for p Nb constrains ϕ production
- ▶ Strangeness production at midrapidity in large systems overestimated
⇒ missing in medium-effects? (KN potentials, in-medium cross section/width)
- ▶ Pion beam very sensitive to resonance properties
- ▶ Not enough pions in final state of B^* decay?
- ▶ Effective many-particle interactions by forced thermalization enhance strangeness production

Outlook

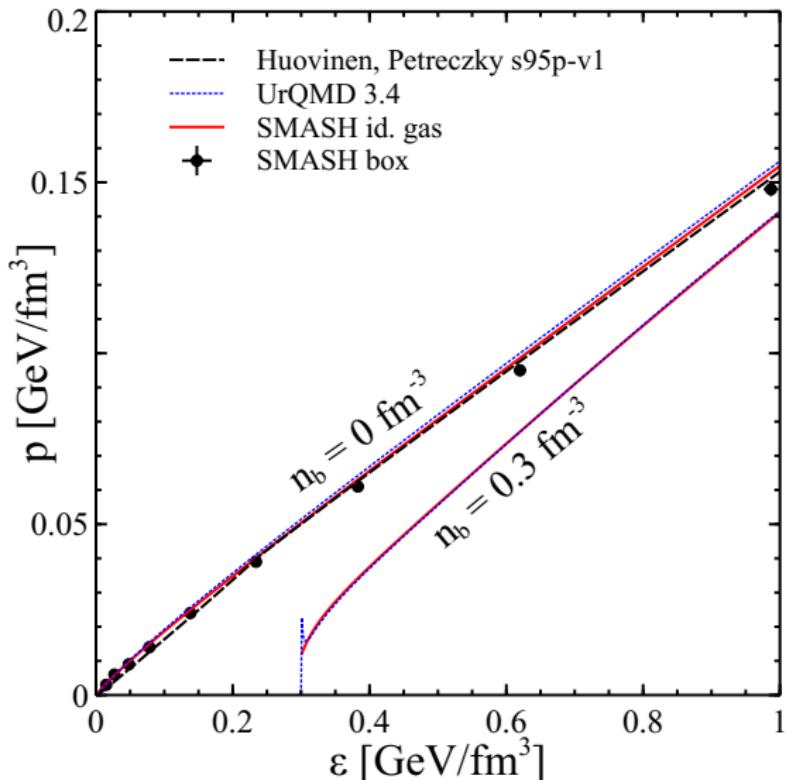


J. E. Bernhard et al. In: *Phys. Rev.* C94.2 (2016), p. 024907. arXiv:
[1605.03954 \[nucl-th\]](https://arxiv.org/abs/1605.03954)

- ▶ Future work: use Bayesian modeling for tuning branching ratios
- ▶ Higher energies require string fragmentation and additive quark model
- ▶ More sophisticated hyperon-nucleon potential⁴ planned

⁴S. Petschauer et al. *Eur. Phys. J.* A52.1 (2016), p. 15. arXiv:
[1507.08808 \[nucl-th\]](https://arxiv.org/abs/1507.08808).

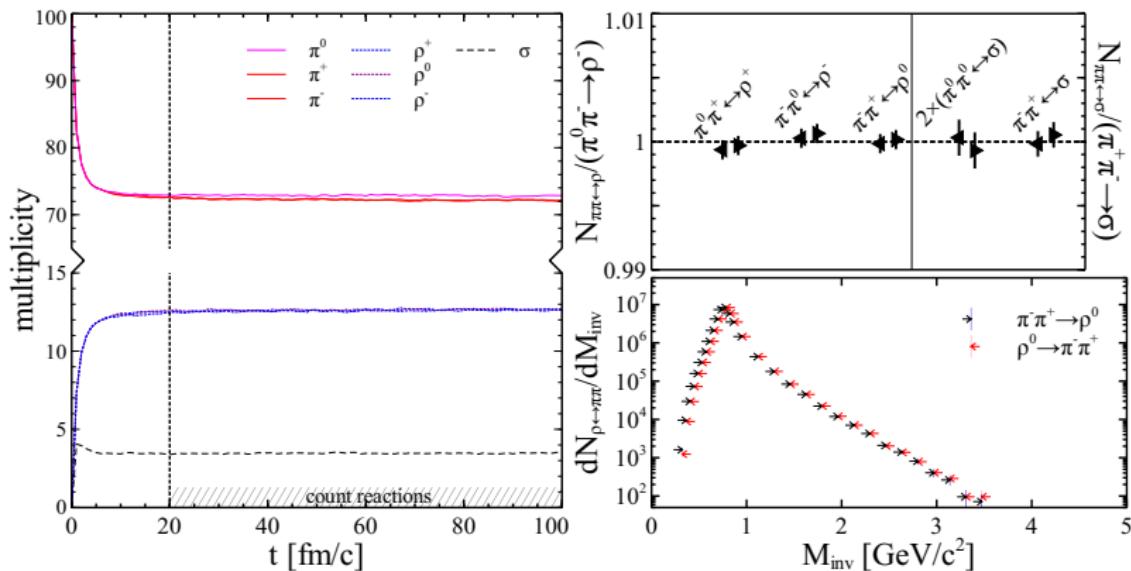
Equation of state



- ▶ SMASH hadron gas vs. UrQMD vs. lattice QCD

Test detailed balance in a $\pi\rho\sigma$ box

- ▶ Initialize periodic box with pions
- ▶ Wait until it equilibrates
- ▶ Count and compare number of forward and backward reactions



Comparison to exact solution of Boltzmann equation

- ▶ Boltzmann equation in curved spacetime

$$p^\mu \frac{\partial f(x, p)}{\partial x^\mu} + p_\lambda p^\mu \Gamma_{\mu i}^\lambda \frac{\partial f(x, p)}{\partial p_i} = C(f) \quad (6)$$

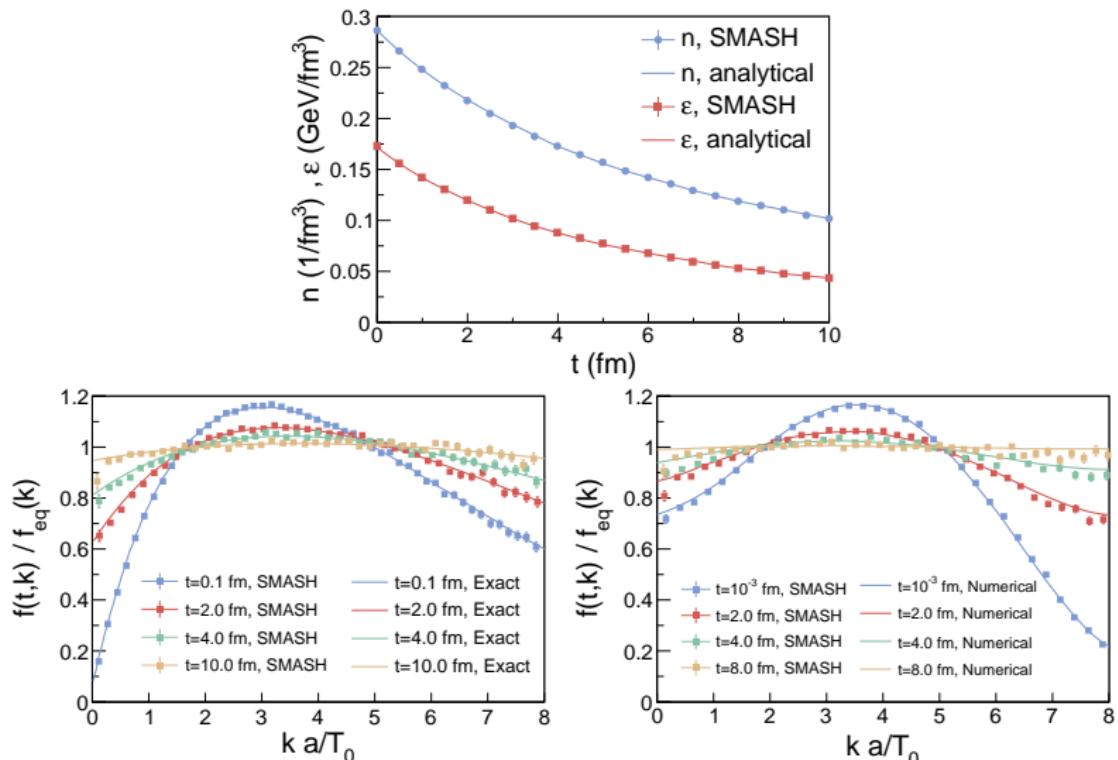
- ▶ Expanding universe with Friedmann-Lemaître-Robertson-Walker metric

$$ds^2 = dt^2 - a(t)^2(dx^2 + dy^2 + dz^2) \quad (7)$$

- ▶ Infinite gas of massless particles with constant elastic cross section
- ▶ An analytic solution exists

D. Bazow et al. In: *Phys. Rev.* D94.12 (2016), p. 125006. arXiv: 1607.05245 [hep-ph]

Comparison to exact solution of Boltzmann equation



J. Tindall et al. In: *Phys. Lett. B* 770 (2017), pp. 532–538. arXiv: 1612.06436 [hep-ph]

Resonances in SMASH

- ▶ Breit-Wigner spectral function

$$\mathcal{A}(m) = \frac{2N}{\pi} \frac{m^2 \Gamma(m)}{(m^2 - m_0^2)^2 + m^2 \Gamma(m)^2} \quad (8)$$

- ▶ Manley-Saleski ansatz⁵ for off-shell decay branching ratio

$$\Gamma_{R \rightarrow ab} = \Gamma_{R \rightarrow ab}^0 \frac{\rho_{ab}(m)}{\rho_{ab}(m_0)} \quad (9)$$

$$\rho_{ab}(m) = \int dm_a dm_b \mathcal{A}_a(m_a) \mathcal{A}_b(m_b) \frac{p_f}{m} B_L^2(p_f R) \mathcal{F}_{ab}^2(m) \quad (10)$$

- ▶ Post form factor⁶ for unstable decay products

$$\mathcal{F}_{ab}(m) = \frac{\lambda^4 + (s_0 - m_0^2)^2/4}{\lambda^4 + (m^2 - (s_0 + m_0^2)/2)^2} \quad (11)$$

⁵D. M. Manley et al. *Phys. Rev.* D45 (1992), pp. 4002–4033.

⁶M. Post et al. *Nucl. Phys.* A741 (2004), pp. 81–148. arXiv: nucl-th/0309085.

Cross sections in SMASH

- ▶ $2 \rightarrow 1$ resonance production

$$\sigma_{ab \rightarrow R}(s) = \frac{2J_R + 1}{(2J_a + 1)(2J_b + 1)} S_{ab} \frac{2\pi^2}{p_i^2} \Gamma_{ab \rightarrow R}(s) \mathcal{A}(\sqrt{s}) \quad (12)$$

- ▶ $2 \rightarrow 2$ resonance production

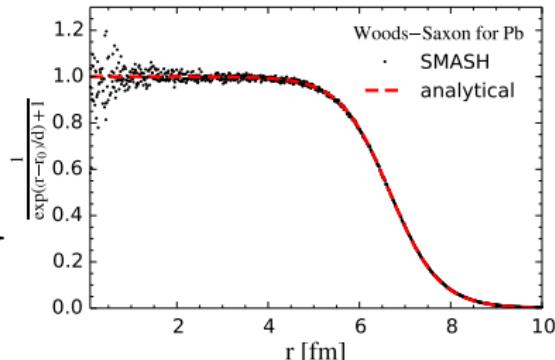
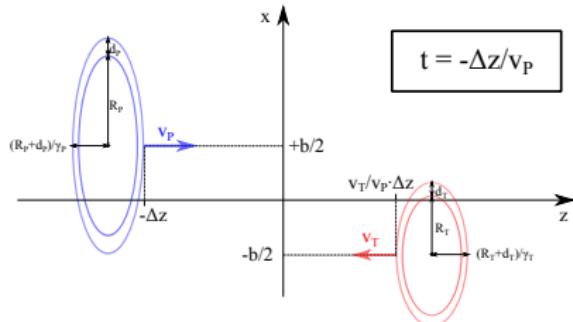
$$\begin{aligned} \sigma_{ab \rightarrow Rc}(s) &= \sum_I (C_{ab}(I) C_{Rc}(I))^2 \frac{|M|_{ab \rightarrow Rc}^2(s, I)}{16\pi} \\ &\times \frac{(2J_R + 1)(2J_c + 1)}{s p_i} \frac{4\pi}{p_{cm}^i} \int dm \mathcal{A}(m) p_f \end{aligned} \quad (13)$$

- ▶ Can model most cross sections like this, some have to be parametrized instead

Modifying particle species and decay modes in SMASH

					N(1440)	
	0.60	1				N π
	0.24	1				Δ π
	0.16	0				N σ
# NAME MASS[GEV] WIDTH[GEV] PDG					N(1520)	
##### unflavored mesons #####						
π	0.138	7.7e-9	111	211	0.65	2 N π
η	0.548	1.31e-6	221		0.10	0 Δ π
σ	0.800	0.400	9000221		0.15	0 N ρ
ρ	0.776	0.149	113	213	0.10	2 Δ π
ω	0.783	8.49e-3	223		0.50	0 N π
η'	0.958	1.98e-4	331		0.40	0 N η
f ₀ (980)	0.990	0.070	9010221		0.06	0 N(1440) π
...					0.02	0 N ρ
...					0.02	0 N σ
##### N baryons #####					N(1535)	
N	0.938	0	2112	2212	0.50	0 N π
N(1440)	1.462	0.350	12112	12212	0.10	0 N η
N(1520)	1.515	0.115	1214	2124	0.08	0 Λ K
N(1535)	1.535	0.150	22112	22212	0.01	0 N ρ
N(1650)	1.655	0.140	32112	32212	0.12	2 N ρ
N(1675)	1.675	0.150	2116	2216		

Nucleus collision



- ▶ Woods-Saxon distribution

$$\frac{dN}{dr} = \frac{\rho_0}{\exp\left(\frac{r-r_0}{d}\right) + 1} \quad (14)$$

- ▶ Deformed nuclei

Fermi motion

- ▶ Local density approximation

$$p_F(\vec{r}) = \hbar c \sqrt[3]{3\pi^2 \rho(\vec{r})} \quad (15)$$

- ▶ Sample momenta p_i from Fermi sphere in nucleus rest frame
- ▶ Boost Fermi momenta to calculation frame

$$p'_{iz} = \gamma(p_{iz} + \beta E_i) = \gamma p_{iz} + \frac{p_A}{A} \quad (16)$$

- ▶ Without potentials:
Ignore Fermi motion for propagation until first interaction

Skyrme and symmetry potential

$$U = a \frac{\rho}{\rho_0} + b \left(\frac{\rho}{\rho_0} \right)^\tau + 2S_{\text{pot}} \frac{\rho_p - \rho_n}{\rho_0} \frac{I_3}{I} \quad (17)$$

$$H_i = \sqrt{\vec{p}_i^2 + m_i^2} + U(\vec{r}_i) \quad (18)$$

where

$$a = -209.2 \text{ MeV} \quad b = 156.4 \text{ MeV} \quad c = 1.35 \quad S_{\text{pot}} = 18 \text{ MeV} \quad (19)$$

- ▶ Nucleus-nucleus only
- ▶ Soft potential with incompressibility $K_0 = 240 \text{ MeV}$
- ▶ Makes nucleus stable despite Fermi motion

Pauli blocking

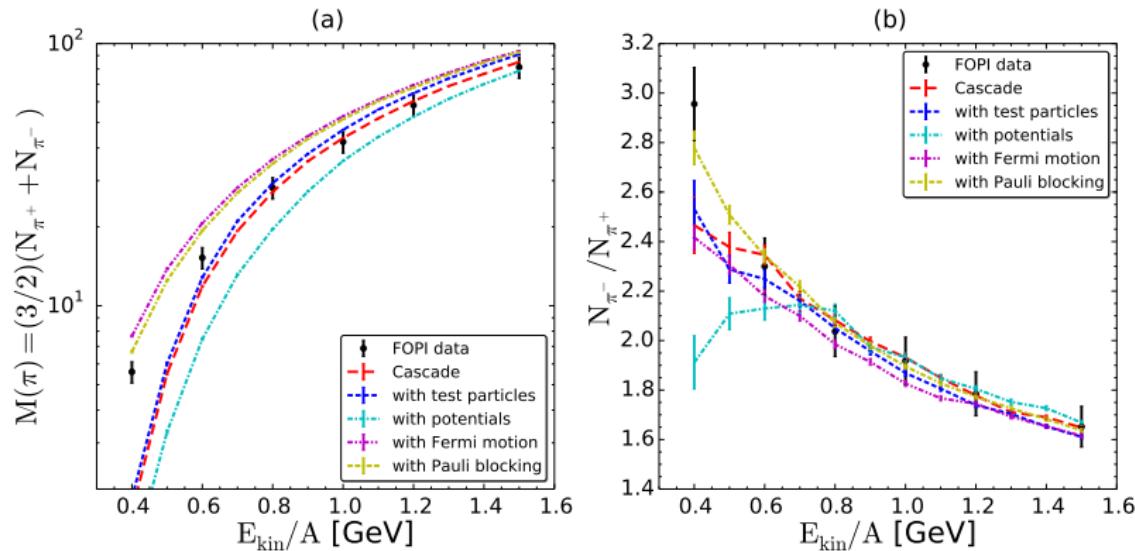
- ▶ Collision integral in Boltzmann-Uehling-Uhlenbeck equation

$$C(f) = \frac{1}{2} \int \frac{d^3 p_2}{E_2} \frac{d^3 p'_1}{E_1} \frac{d^3 p'_2}{E'_2} W(p_1, p_2, p'_1, p'_2) \times (f'_1 f'_2 (1 \pm f)(1 \pm f_2) - f f_2 (1 \pm f'_1)(1 \pm f'_2)) \quad (20)$$

- ▶ Pauli blocking and Bose enhancement
- ▶ Reject reactions with probability

$$P = 1 - \prod_{\text{final state fermion } i} (1 - f_i) \quad (21)$$

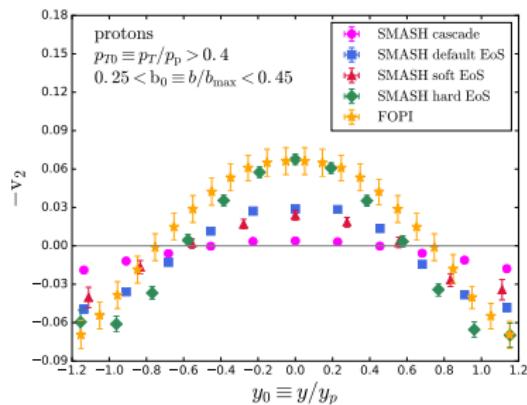
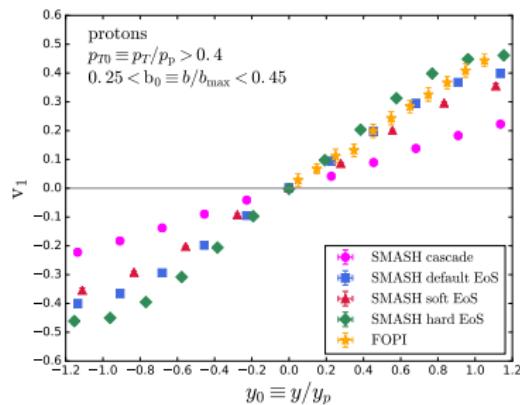
Pion production in central gold-gold collisions



W. Reisdorf et al. In: *Nucl. Phys.* A781 (2007), pp. 459–508. arXiv: [nucl-ex/0610025](https://arxiv.org/abs/nucl-ex/0610025)

- ▶ Yield overestimated, but ratio reproduced
- ▶ FOPI pion multiplicities sensitive to nucleonic potentials and Pauli blocking

Flow in gold-gold collisions at $E_{\text{kin}} = 1A \text{ GeV}$



- ▶ Sensitive to parameters of nucleonic potentials
- ▶ Hard equation of state reproduces data best

W. Reisdorf et al. In: *Nucl. Phys.* A876 (2012), pp. 1–60. arXiv:
1112.3180

Analysis suite

- ▶ Extensive collection of tests for the model
- ▶ Fully automated, checked for every SMASH release
- ▶ Consistency checks:
 - ▶ Detailed balance: Check equilibrium in thermalized box
 - ▶ Elastic box: Comparison to ideal gas expectations
- ▶ Comparison to experimental data:
 - ▶ Angular distributions: pp , np at $\sqrt{s} \approx 2.5$ GeV
 - ▶ Elementary cross sections: NN , πN , $\pi\pi$, KN
 - ▶ FOPI pions: π multiplicities for $E_{\text{kin}} = 0.4 - 1.5A$ GeV
 - ▶ Spectra: dN/dy and dN/dm_T for π and p in AuAu at $E_{\text{kin}} = 1.5A$ GeV and CC at $E_{\text{kin}} \in \{1, 2\}A$ GeV
- ▶ Of interest to other models targeting NICA/FAIR energies?
- ▶ Systematic comparison of models?