

My adventures with thermal model

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multiplicities of identified particles in high energy collisions

$$\frac{dN}{dy} \quad \text{of} \quad \pi, K, p, \bar{p}, \Lambda, \dots$$

well reproduced by statistical *i.e.* thermal models

Thermal model

Particles emitted from

- single source
- in chemical equilibrium
- characterised by V and T, μ_B, μ_S, μ_I

Yields given by

$$N_i = V n_i = -T \frac{\partial Z_i(T, \mu_i)}{\partial \mu}$$

Thermal model

Dashen-Ma-Berstein theorem: If interactions mediated by *narrow* resonances, properties of interacting hadron gas are those of noninteracting hadron-resonance gas \Rightarrow **Hadron resonance gas model**

Partition function of noninteracting hadron resonance gas:

$$\ln Z_i = \frac{V g_i}{2\pi^2} \int_0^\infty \pm p^2 dp \ln(1 \pm \exp(-(E_i - \mu_i)/T))$$

Stable particle yields must include contribution from resonance decays:

$$N_\pi = V(n_\pi + 2n_\rho + n_\Delta + \dots)$$

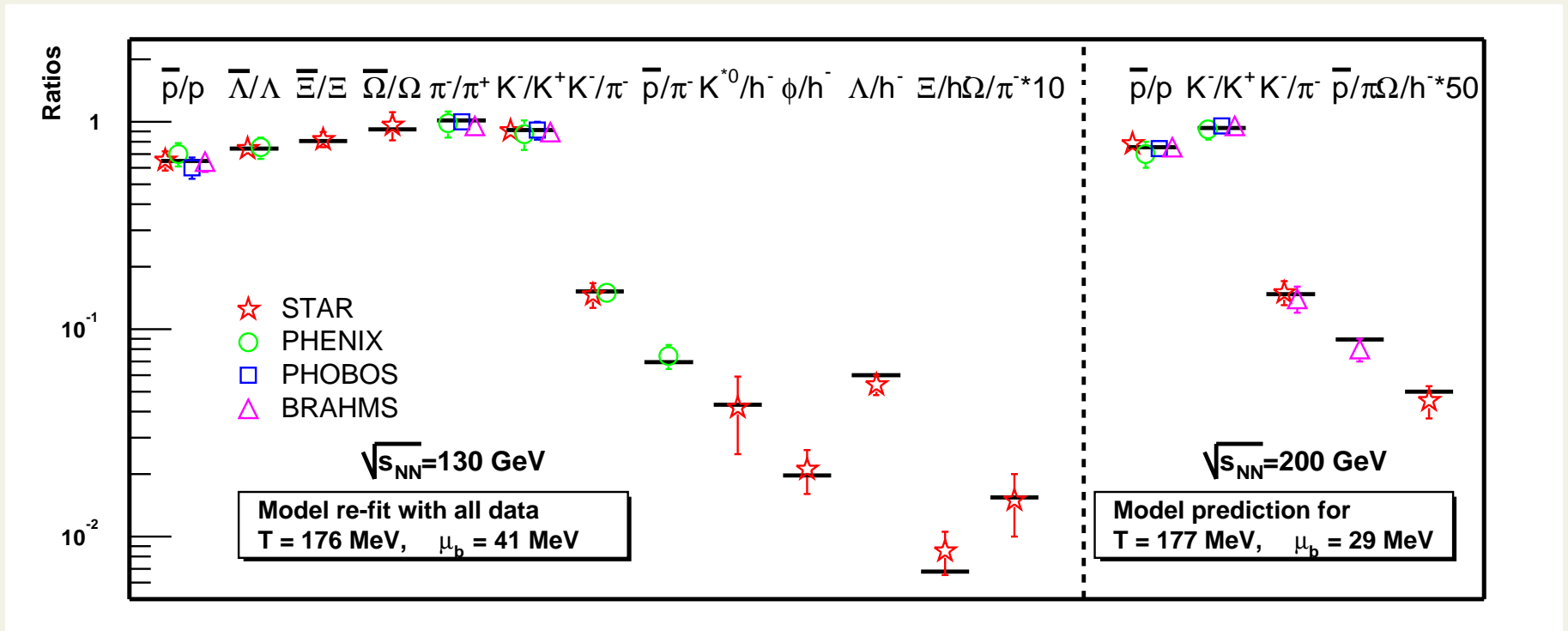
Conservation laws constrain μ_S, μ_I

• **at LHC $\mu_B = \mu_S = \mu_I = 0$**

Fit at each collision energy provides values for $V, T,$ and μ_B

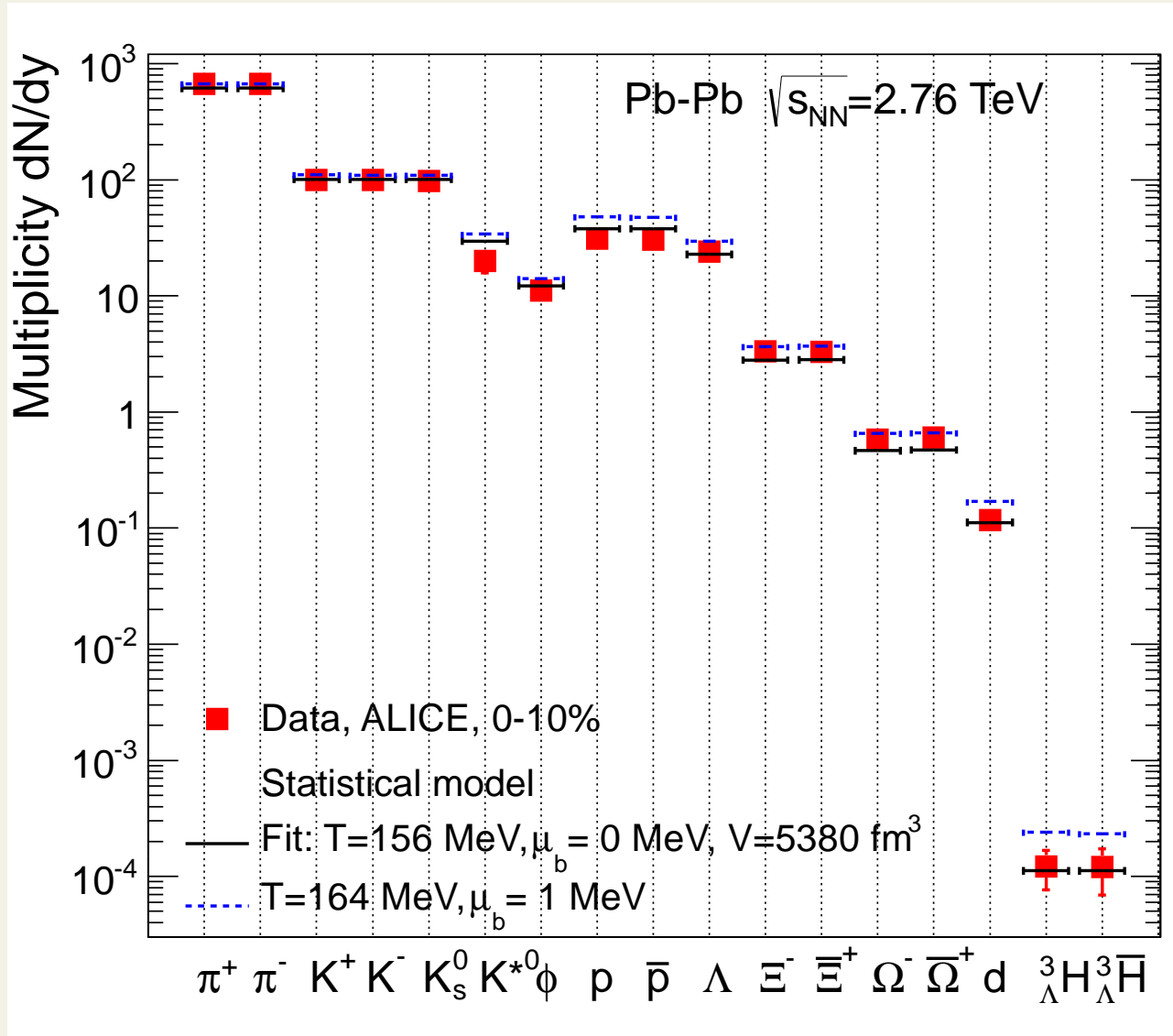
Fit to RHIC data

Braun-Munzinger *et al.*, PLB518, 41 (2001)



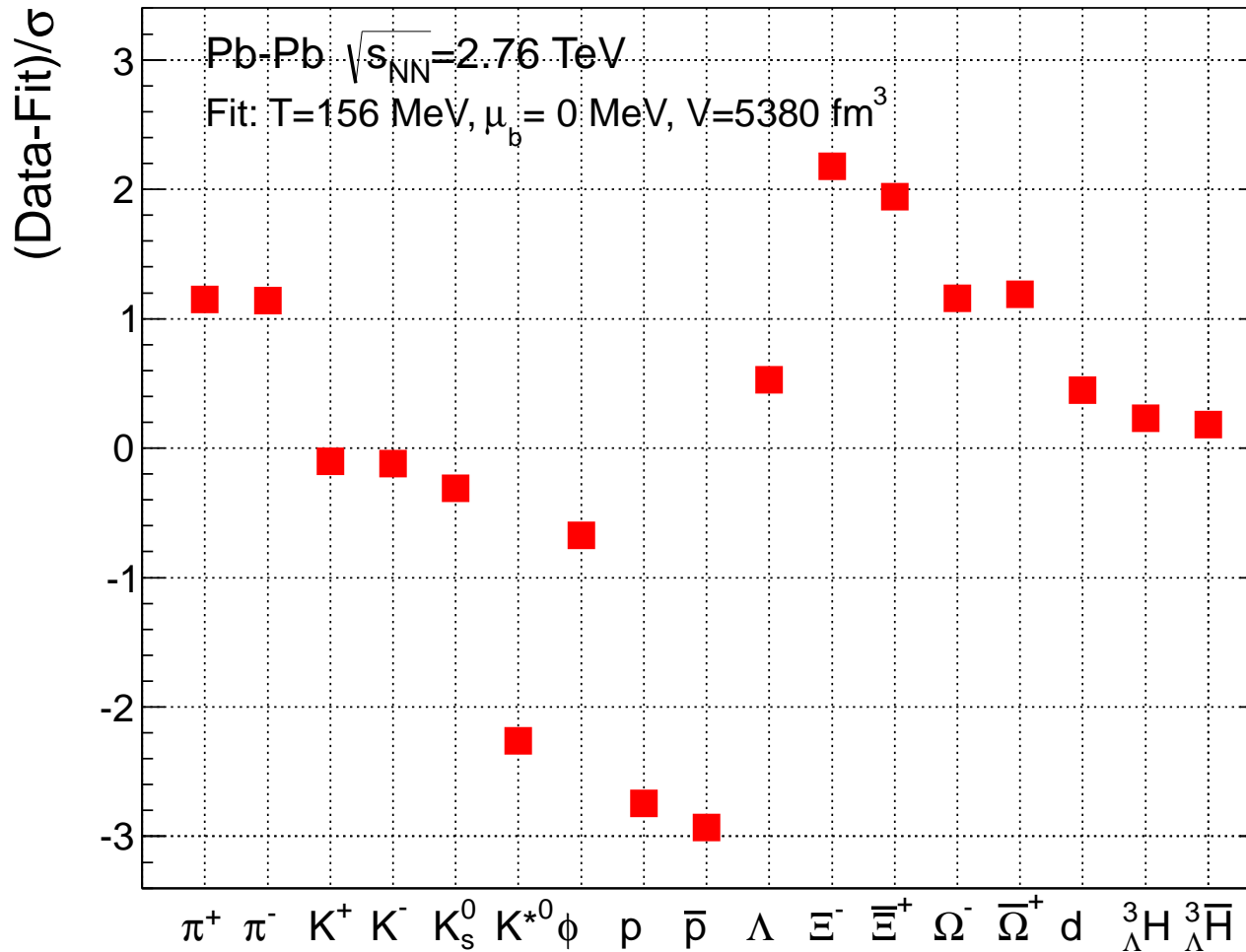
Fit to ALICE data

J. Stachel @ sQM13, arXiv:1311.4662



Fit to ALICE data

J. Stachel @ sQM13, arXiv:1311.4662



- **more protons/antiprotons** than in the data, $\sim 3\sigma$ difference
- **slightly too few cascades**, $\sim 2\sigma$ difference

usual explanation:

- $N\bar{N}$ annihilations continue after chemical equilibrium is lost
- ⇒ hybrid models (UrQMD)

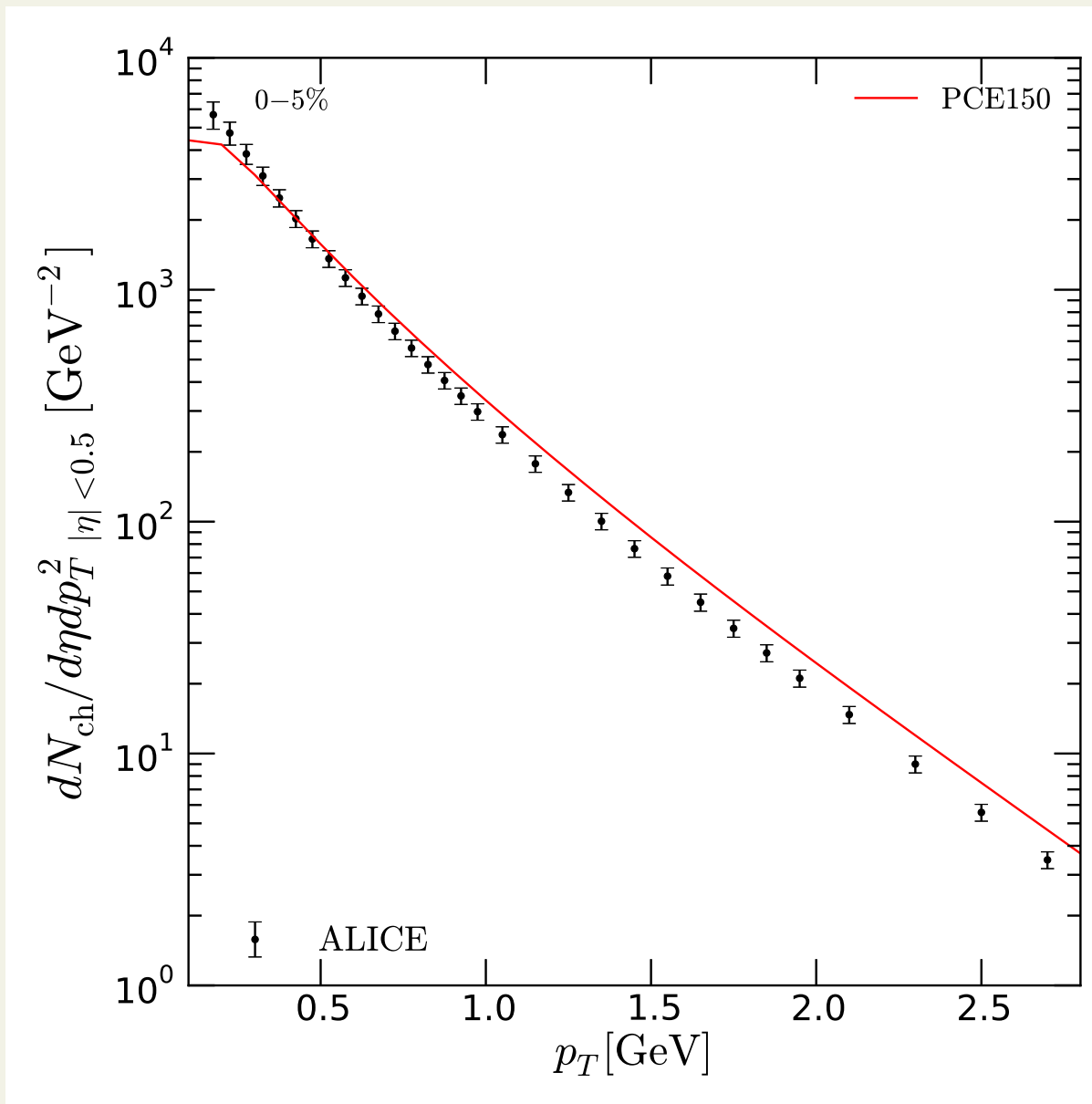
BUT

- hyperons?
- detailed balance?



- how to do fits?
- pure hydro?

Charged hadron p_T spectrum at LHC

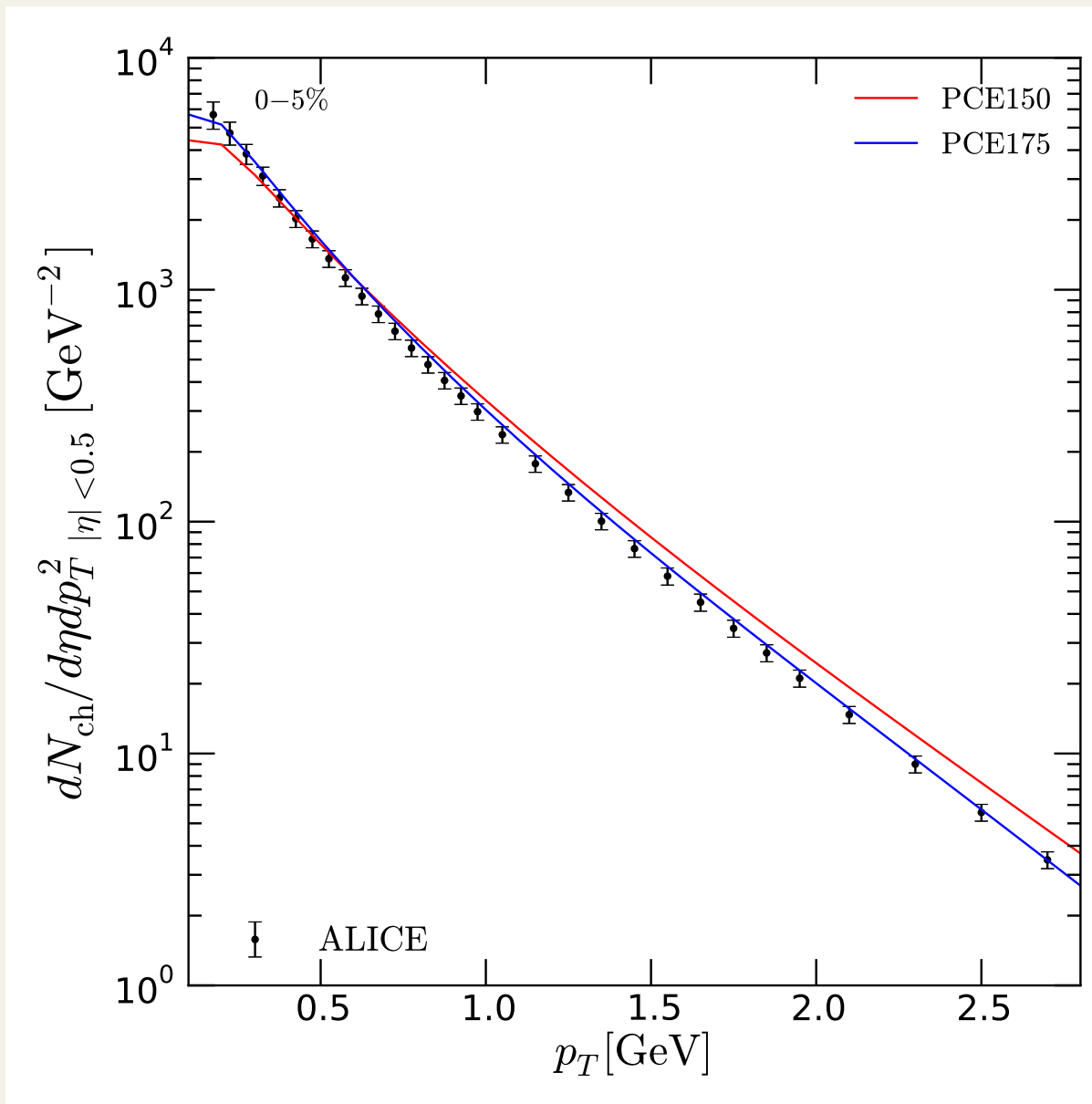


- viscous hydro
- initial state:
pQCD+saturation
- $\tau_0 \approx 0.2\text{fm}/c$

PCE150:
fit to π , K , p yields
no fit to spectrum

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Charged hadron p_T spectrum at LHC



- viscous hydro
- initial state:
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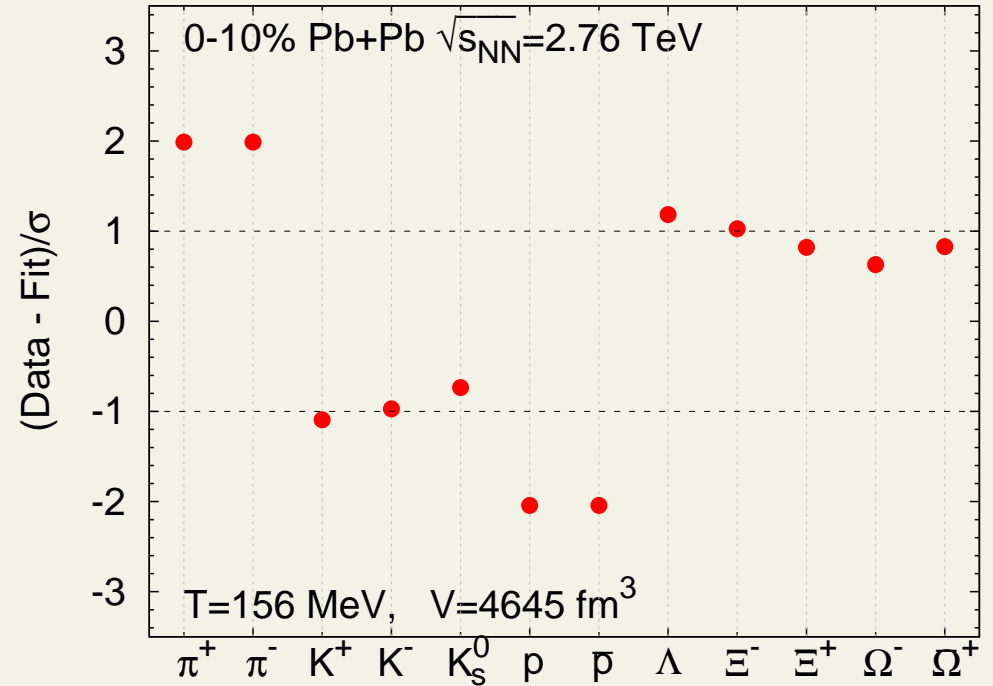
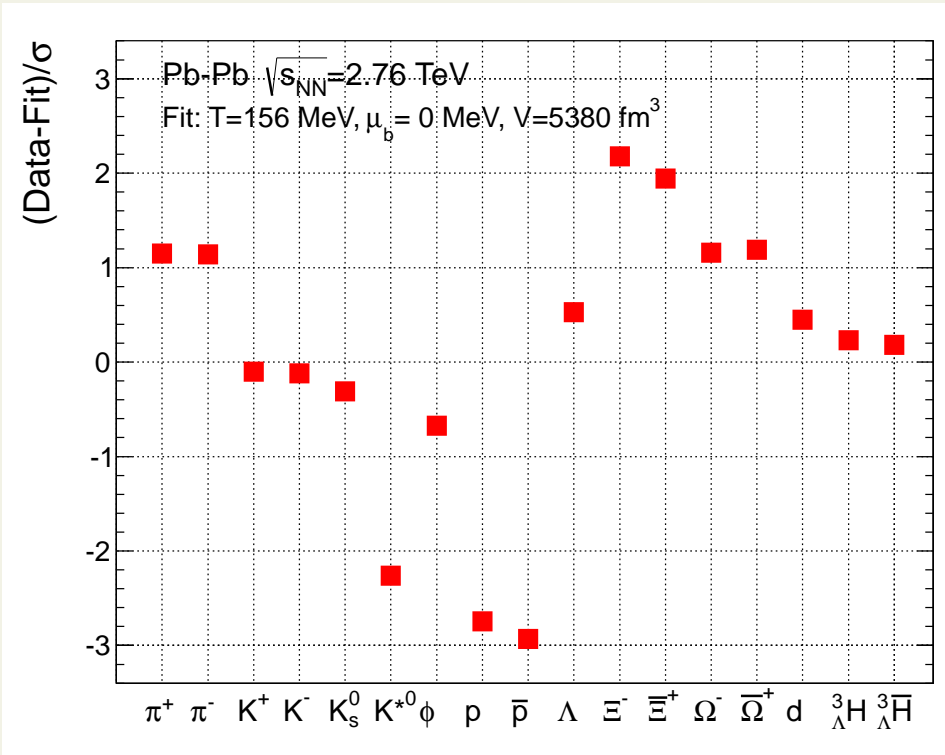
PCE150:
fit to π , K , p yields
no fit to spectrum

PCE175:
no fit to yields
fits the spectrum

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Fit to ALICE data

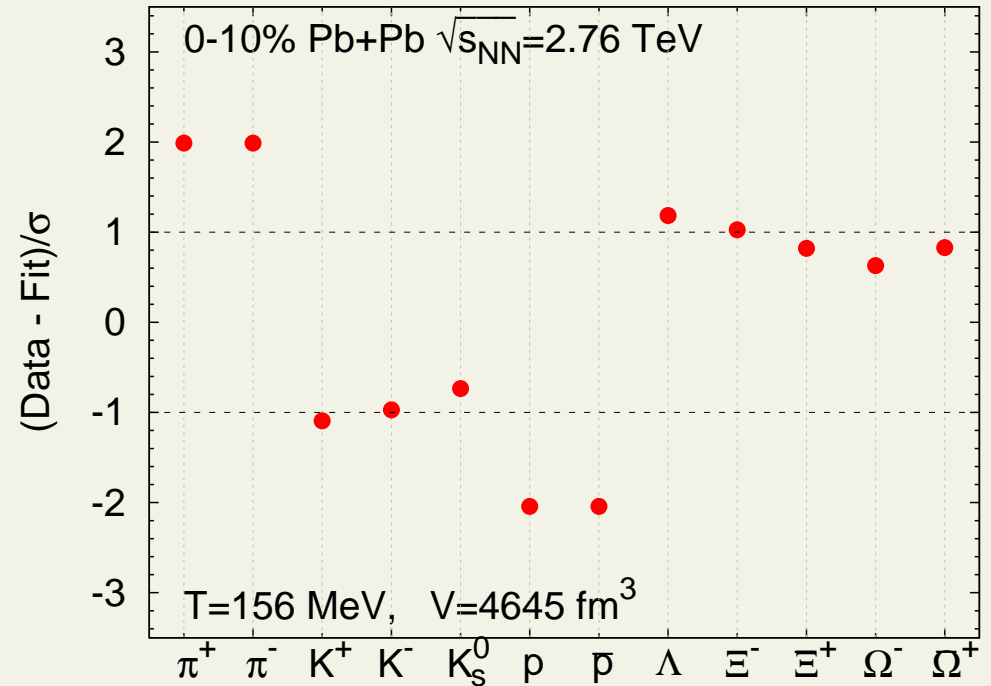
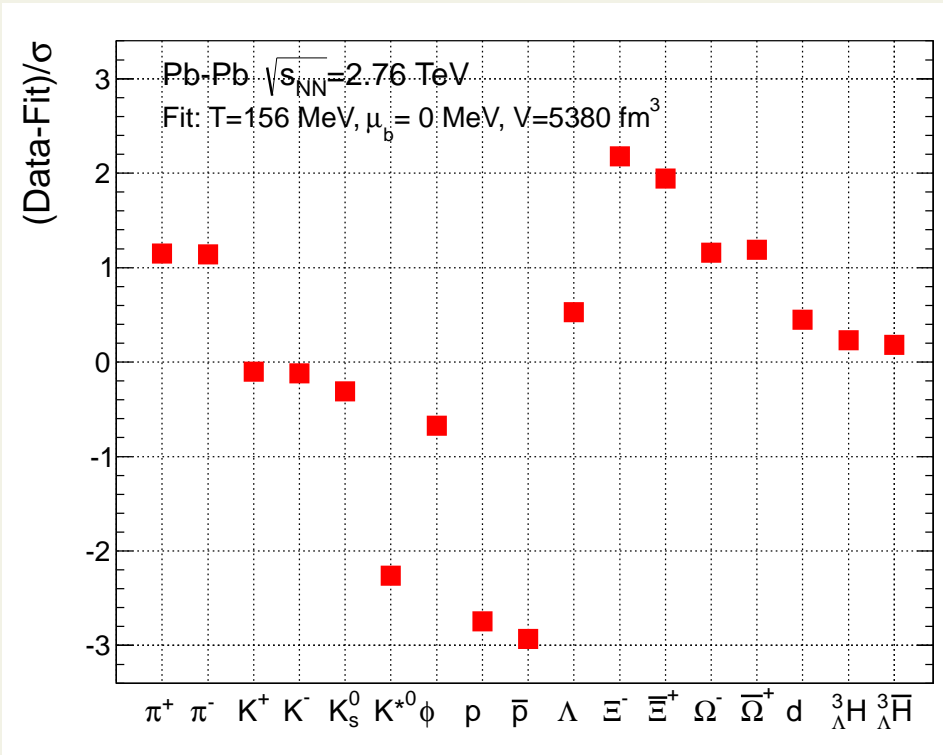
J. Stachel @ sQM13, arXiv:1311.4662



● Why the difference?

Fit to ALICE data

J. Stachel @ sQM13, arXiv:1311.4662

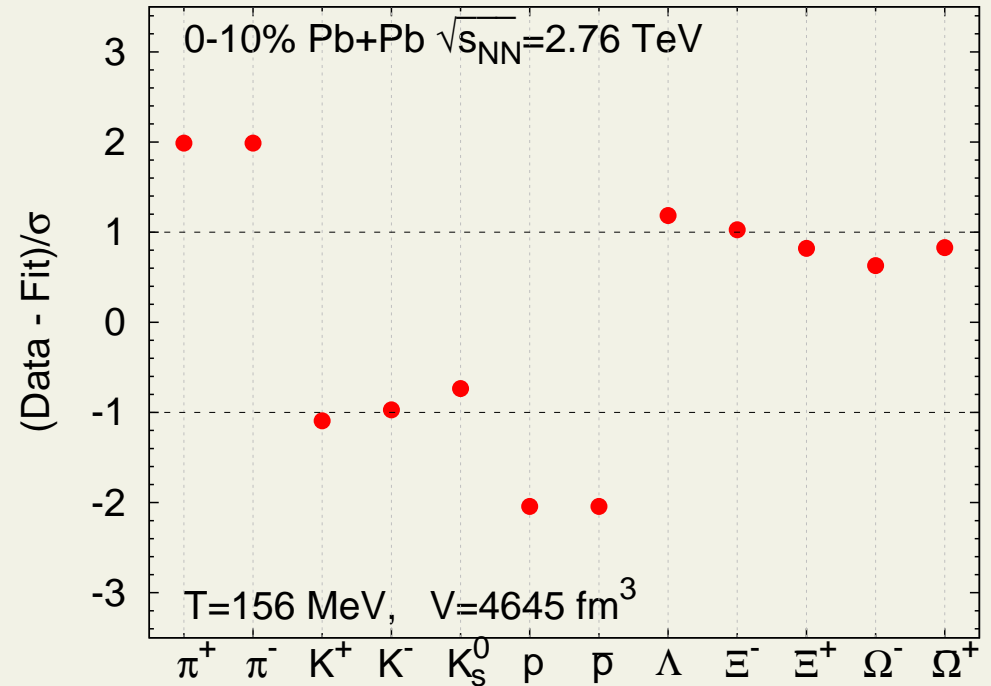
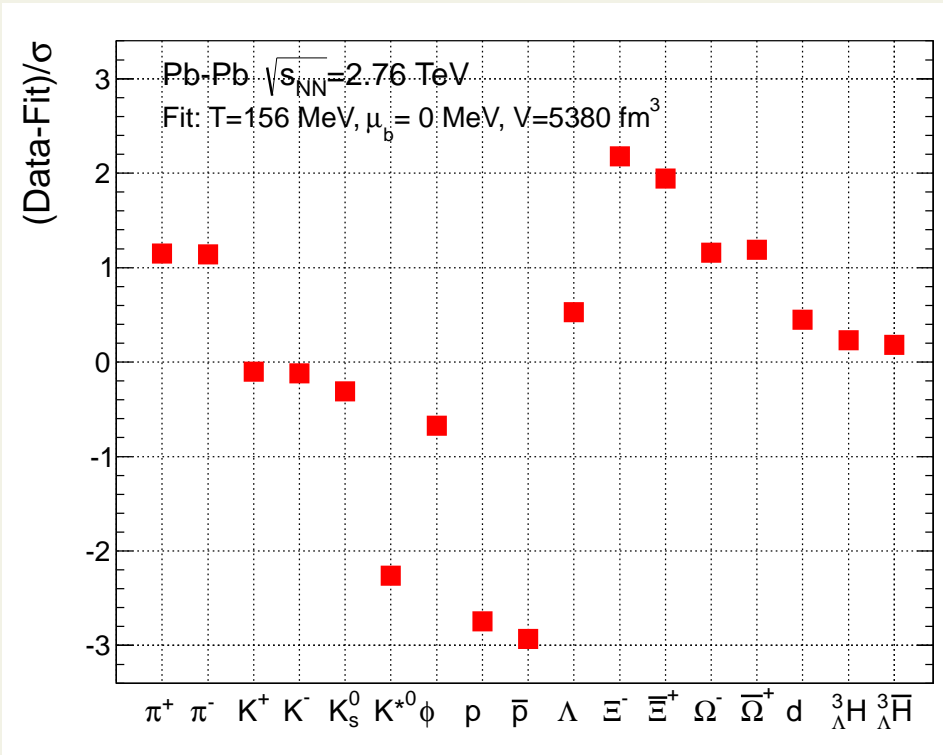


● Why the difference?

- Excluded volume correction does not affect ratios

Fit to ALICE data

J. Stachel @ sQM13, arXiv:1311.4662

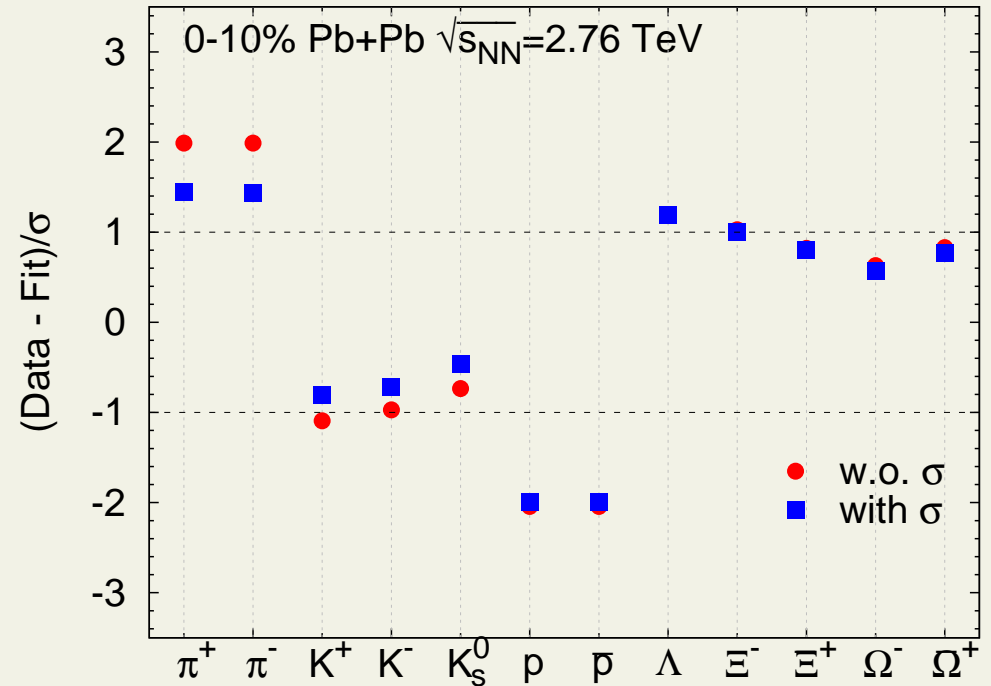
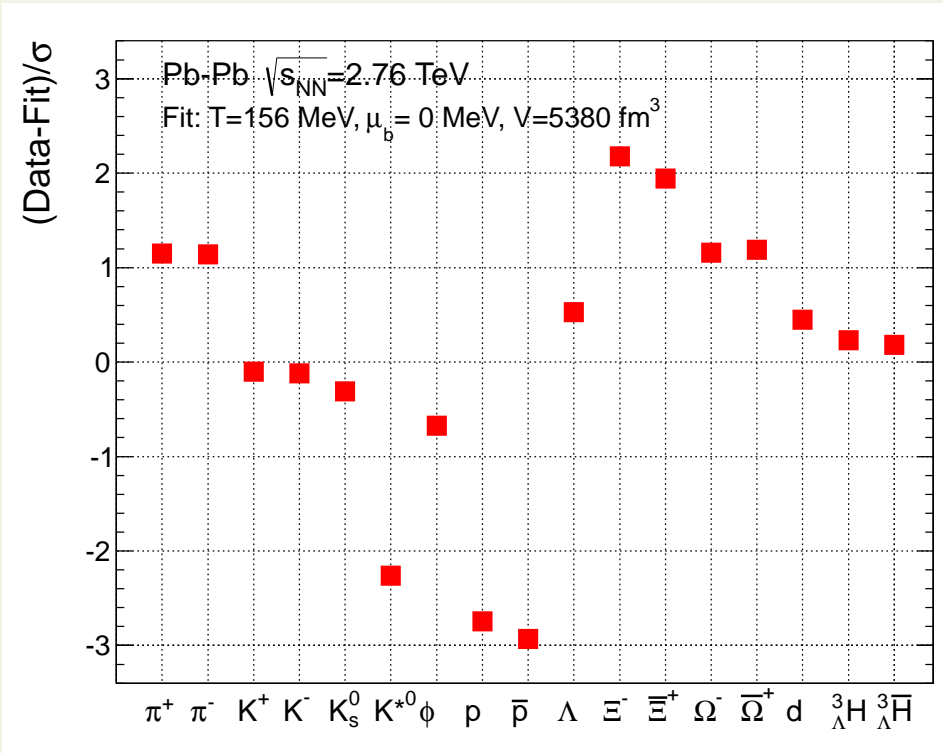


Why the difference?

- Excluded volume correction does not affect ratios
- **sigma-meson!**

Fit to ALICE data

J. Stachel @ sQM13, arXiv:1311.4662



● Inclusion of σ -meson **improves the fit**

● **but. . .**

$f_0(500)$ a.k.a. σ -meson

- mass 400-550 MeV
- width 400-700 MeV

Dashen-Ma-Bernstein: **resonances with narrow widths. . .**

Prakash & Venugopalan, NPA546, 718 (1992): experimental phase shifts

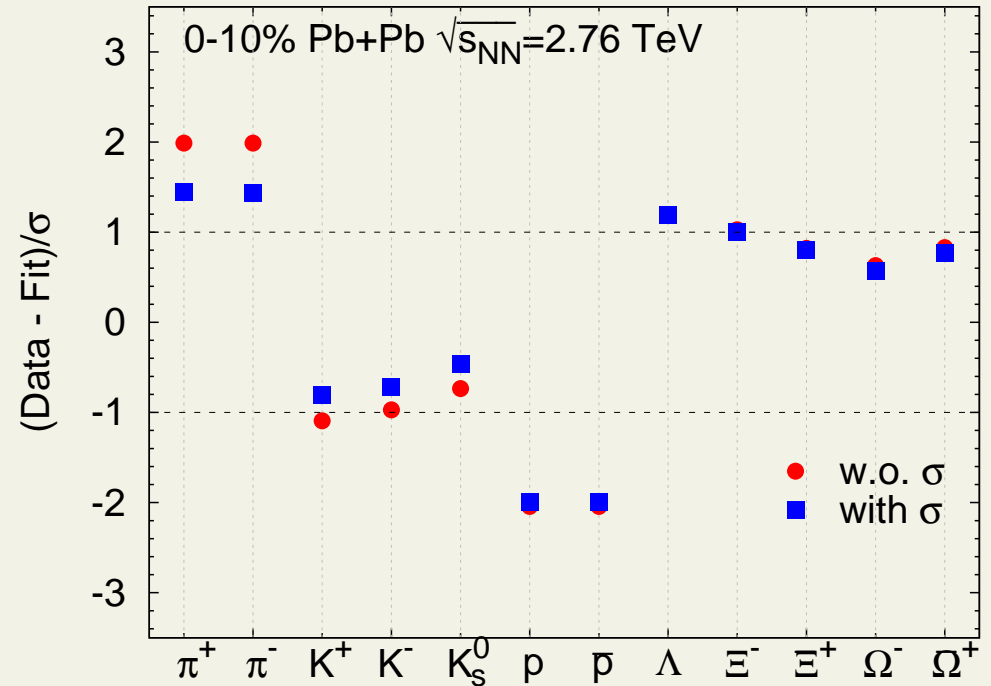
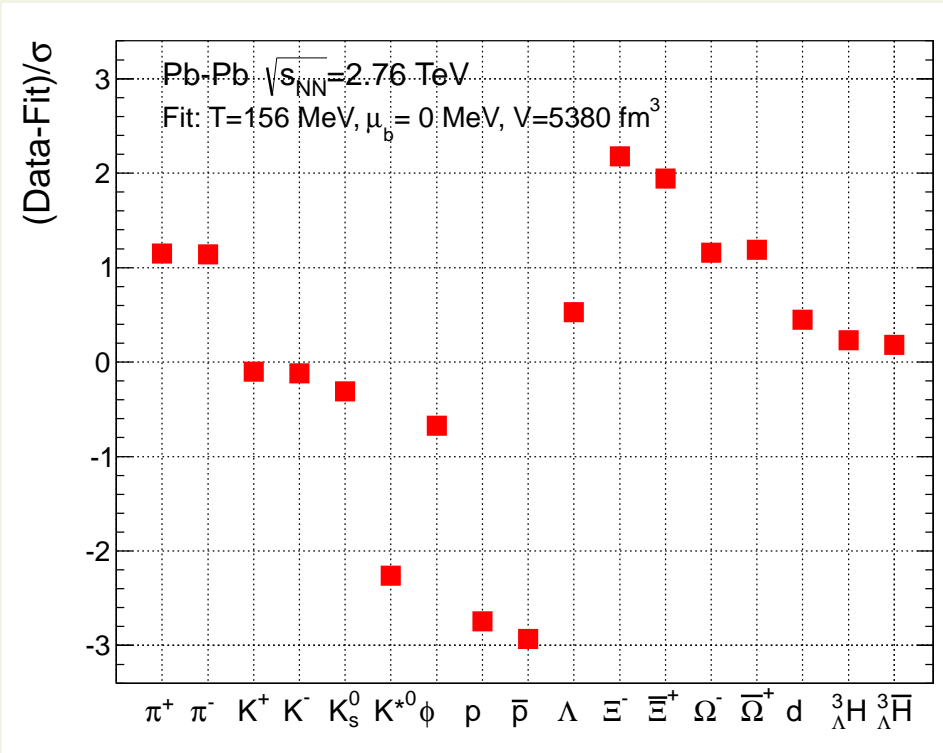
⇒ **HRG good approximation at low temperatures**

- phase shifts **contain** σ
- their HRG **did not**

⇒ **Does σ -meson belong to HRG??**

Fit to ALICE data

J. Stachel @ sQM13, arXiv:1311.4662



● Still fairly large difference. . .

● Finite resonance widths!

Particle densities:

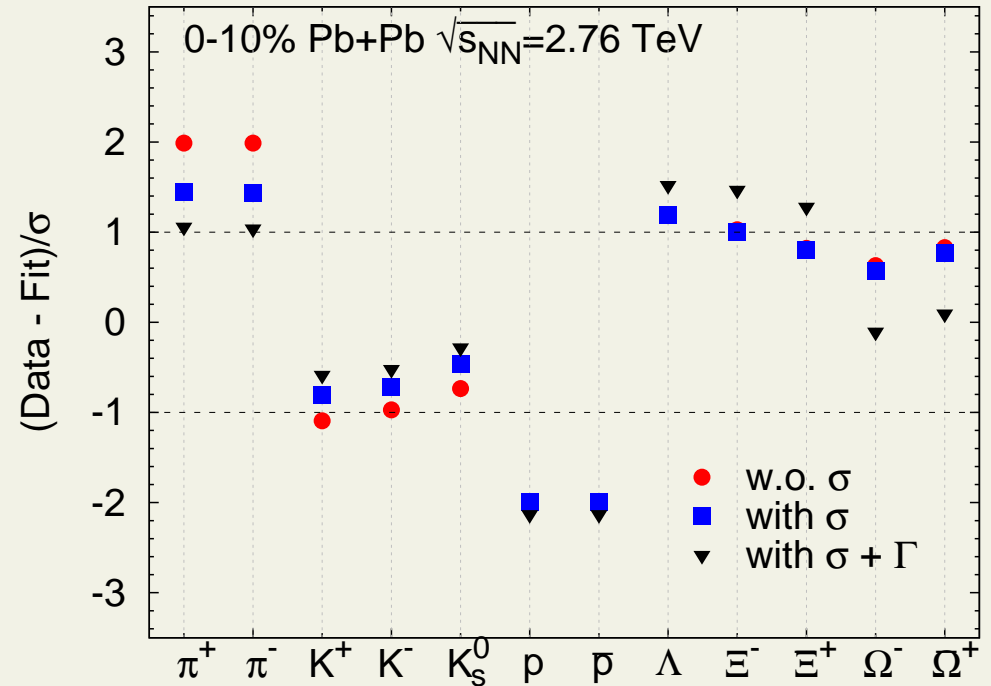
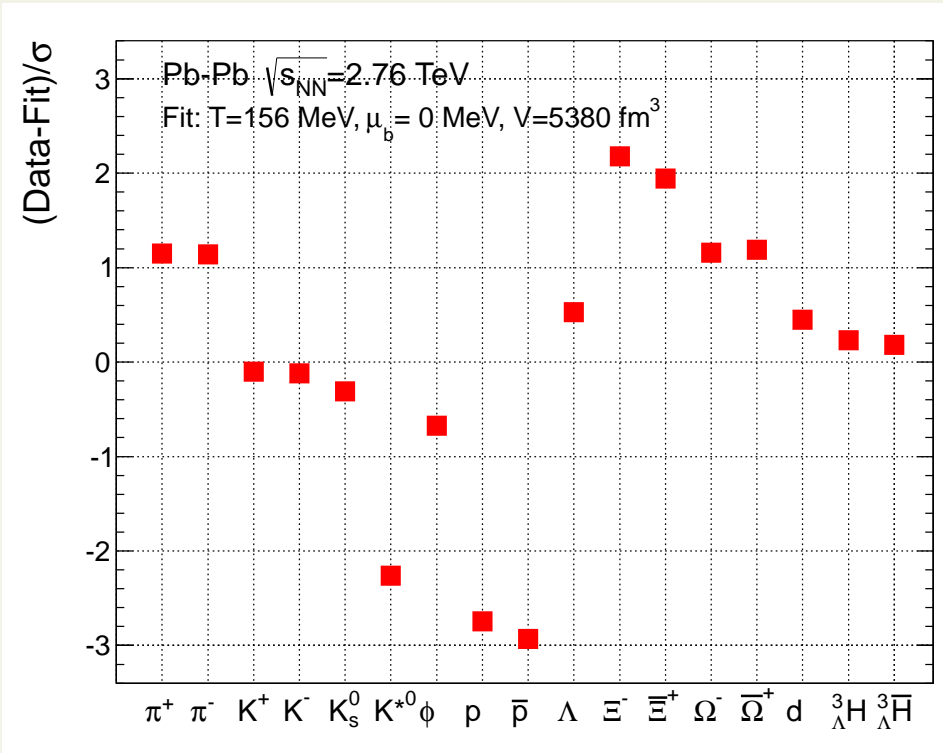
$$n_i = \frac{g_i}{2\pi^2} \int_0^\infty dp \frac{p^2}{\exp\{(\sqrt{m_i^2 + p^2} - \mu_i)/T\} \pm 1}$$

Finite widths can be taken into account by **integrating over mass with Breit-Wigner distribution** as weight:

$$n_i = \frac{g_i}{2\pi^2} \frac{1}{N_{\text{BW}}} \int_{M_0}^\infty dm \frac{\Gamma_i^2}{(m - m_i)^2 + \Gamma_i/4} \int_0^\infty dp \frac{p^2}{\exp\{(\sqrt{m_i^2 + p^2} - \mu_i)/T\} \pm 1}$$

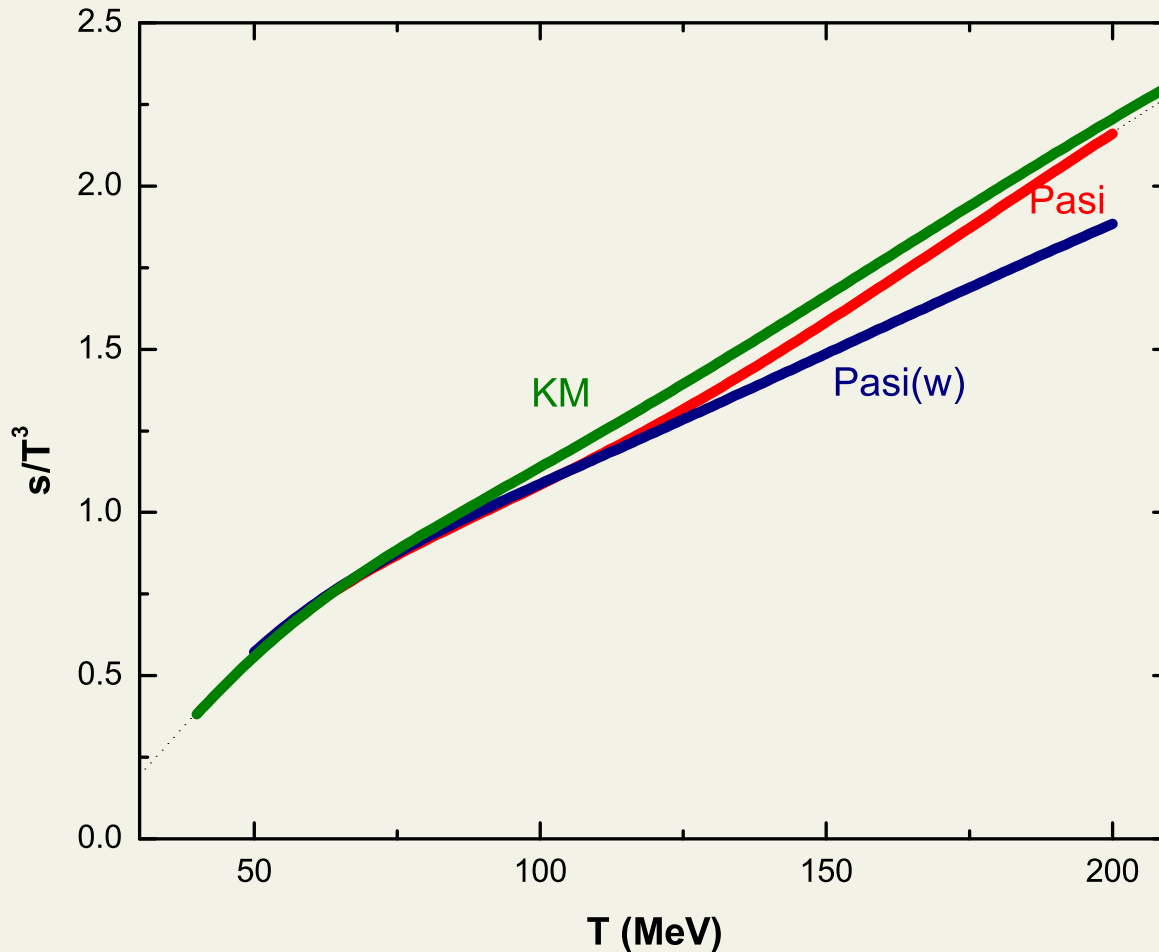
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● differences due to different resonance lists

Entropy in interacting pion gas



KM: Anton Wiranata's
K-matrix calculation

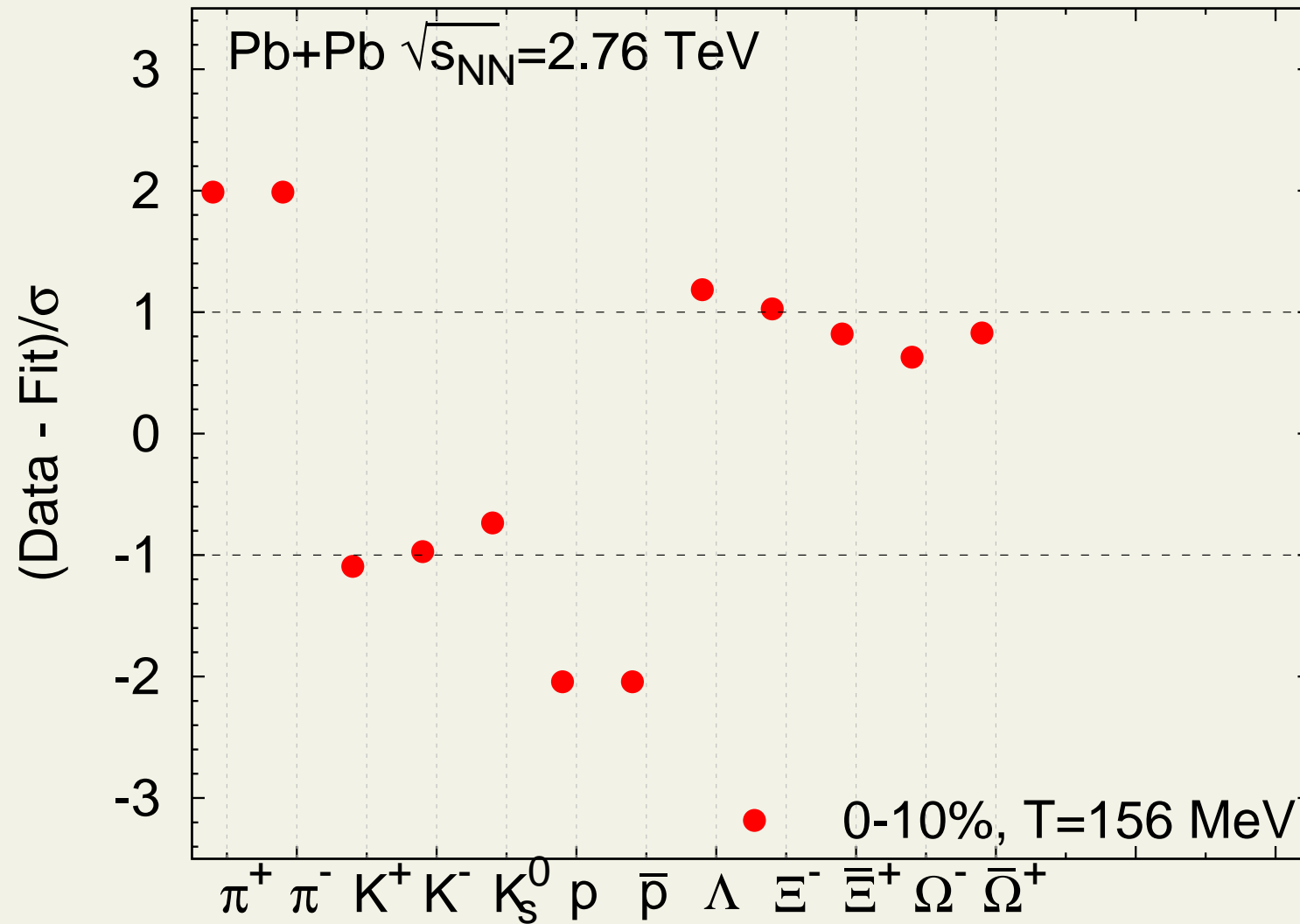
Pasi: $\pi + \rho$ -gas, no widths

Pasi(w): $\pi + \rho$ -gas,
finite width

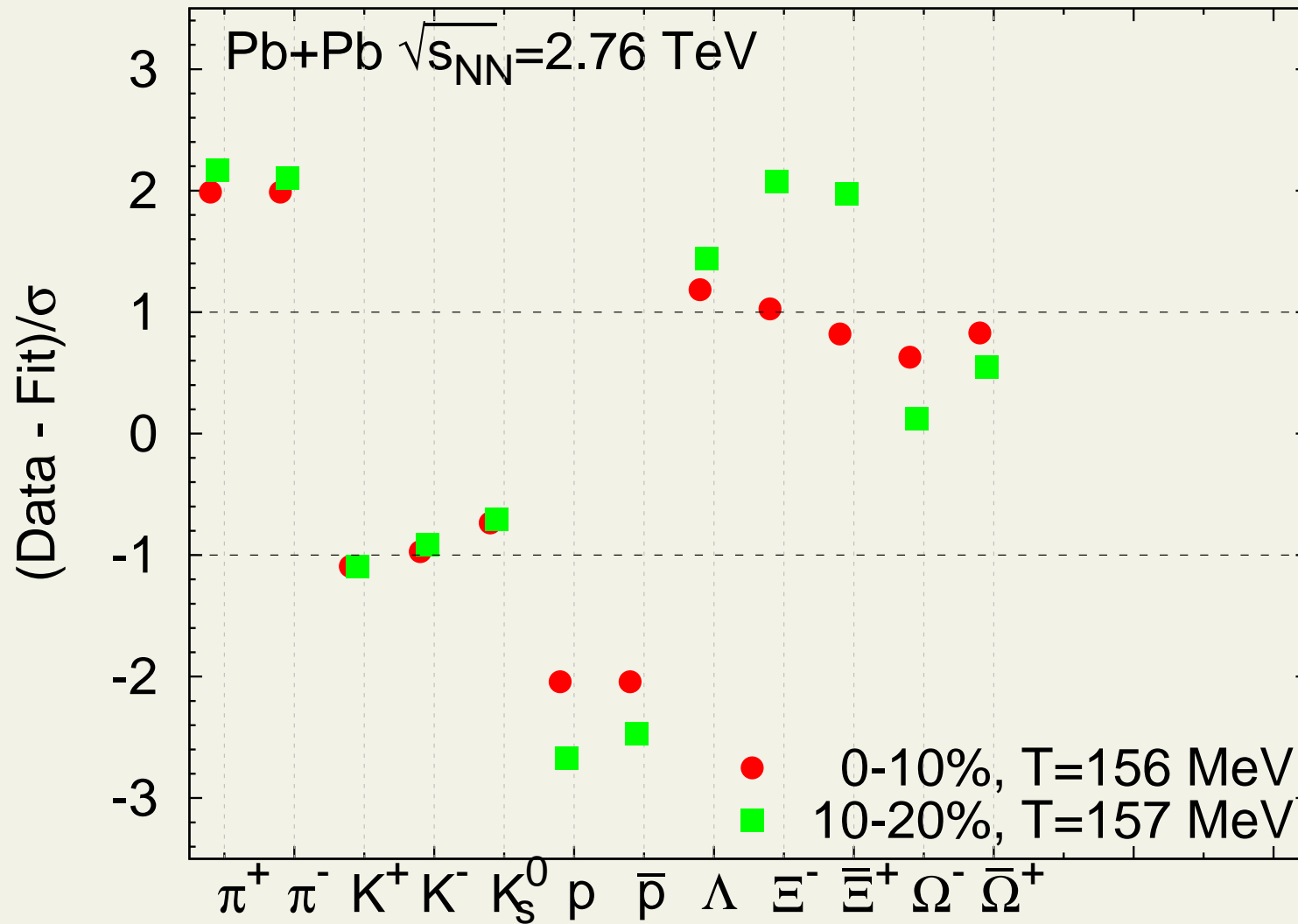
- zero width gives a better approximation!
- spectra calculation very tedious

⇒ no widths nor σ -meson in the following

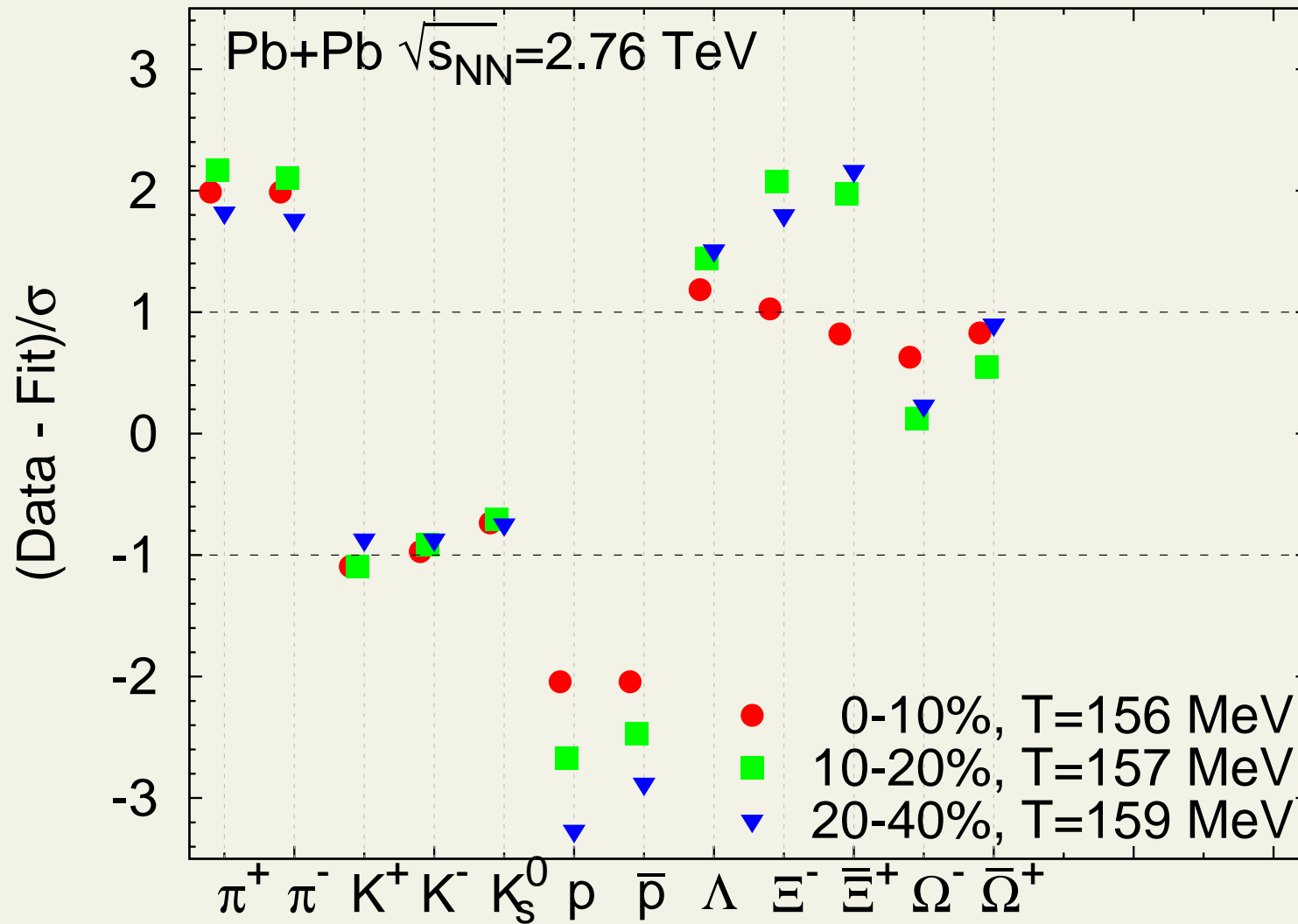
Centrality dependence



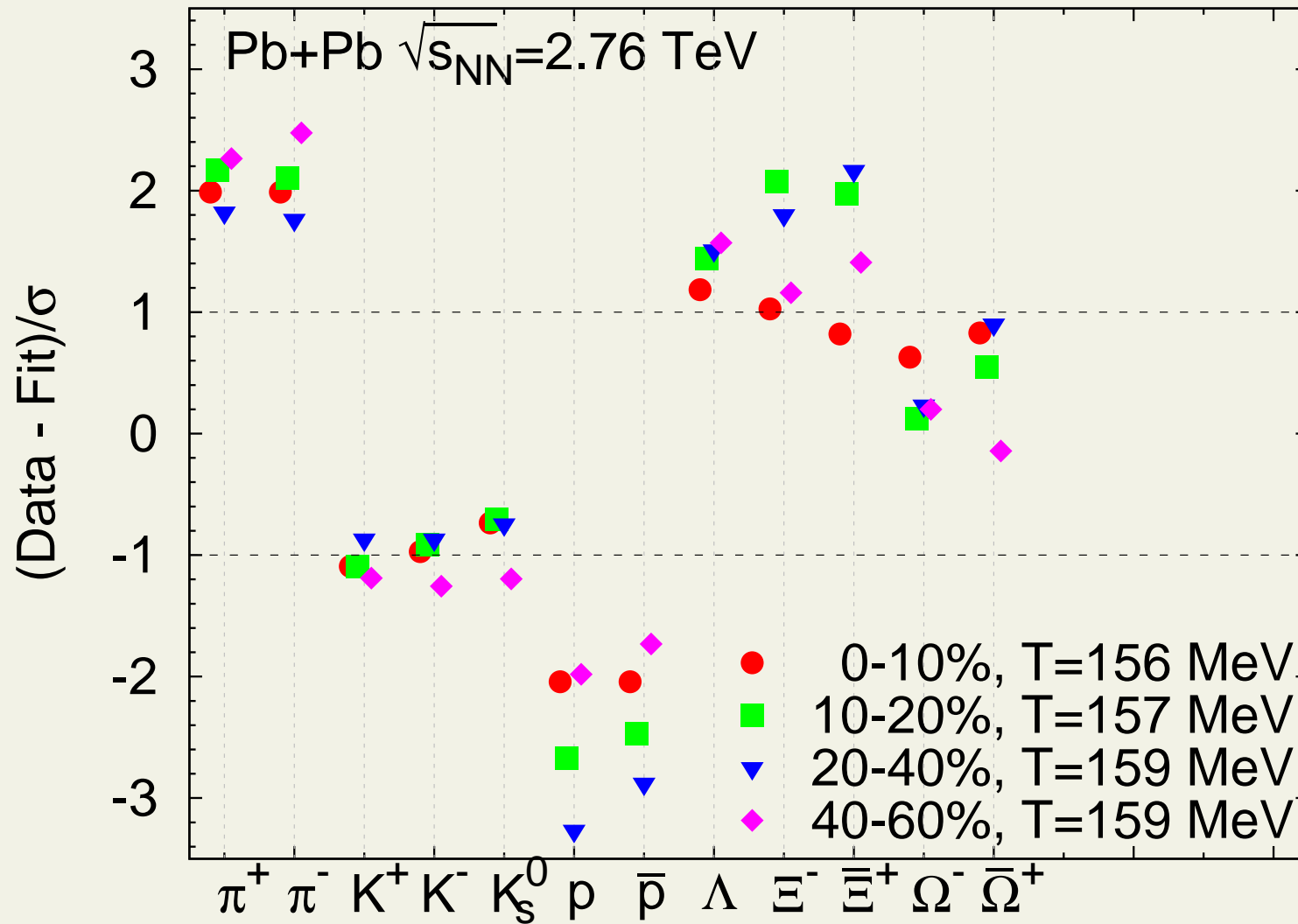
Centrality dependence



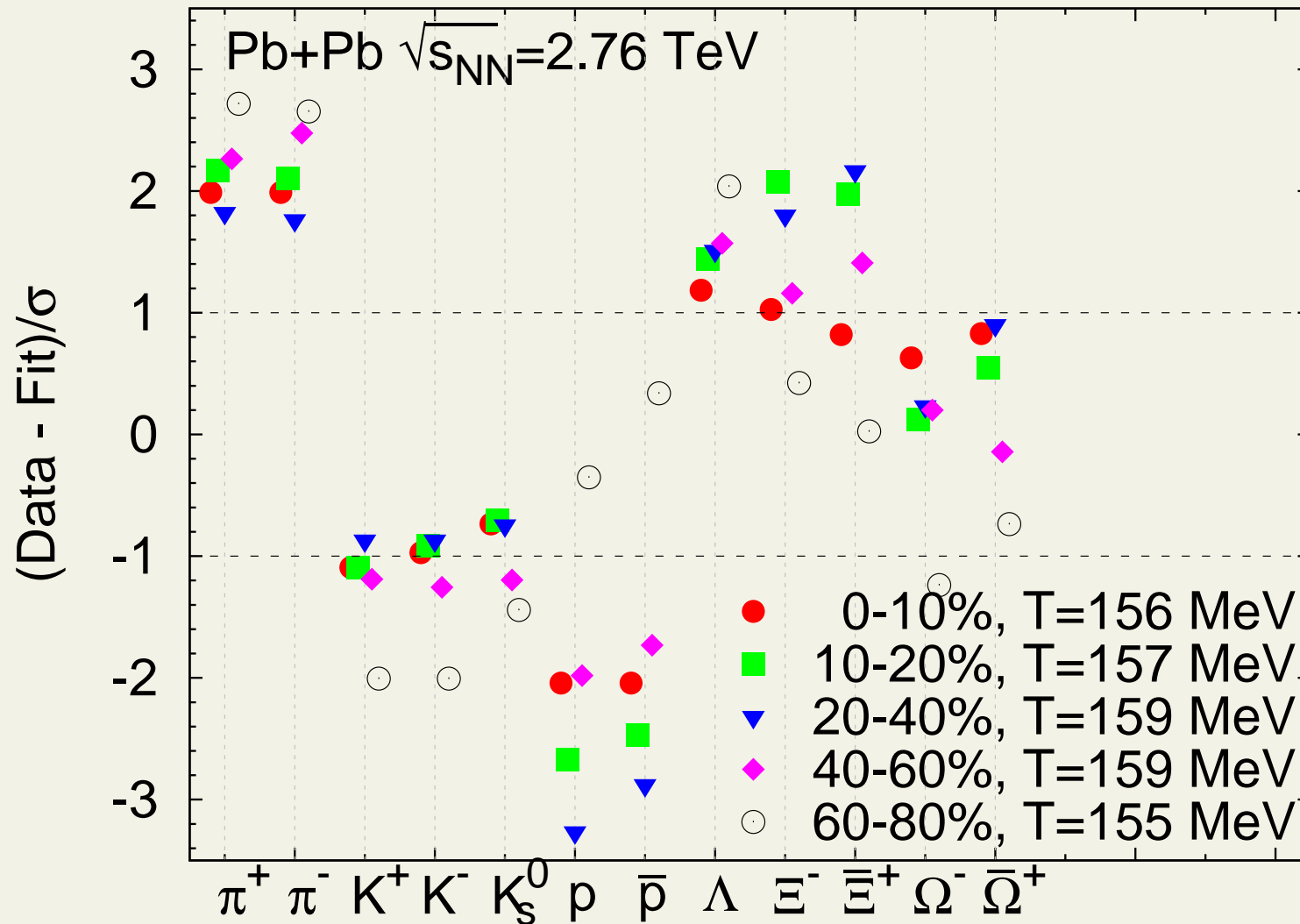
Centrality dependence



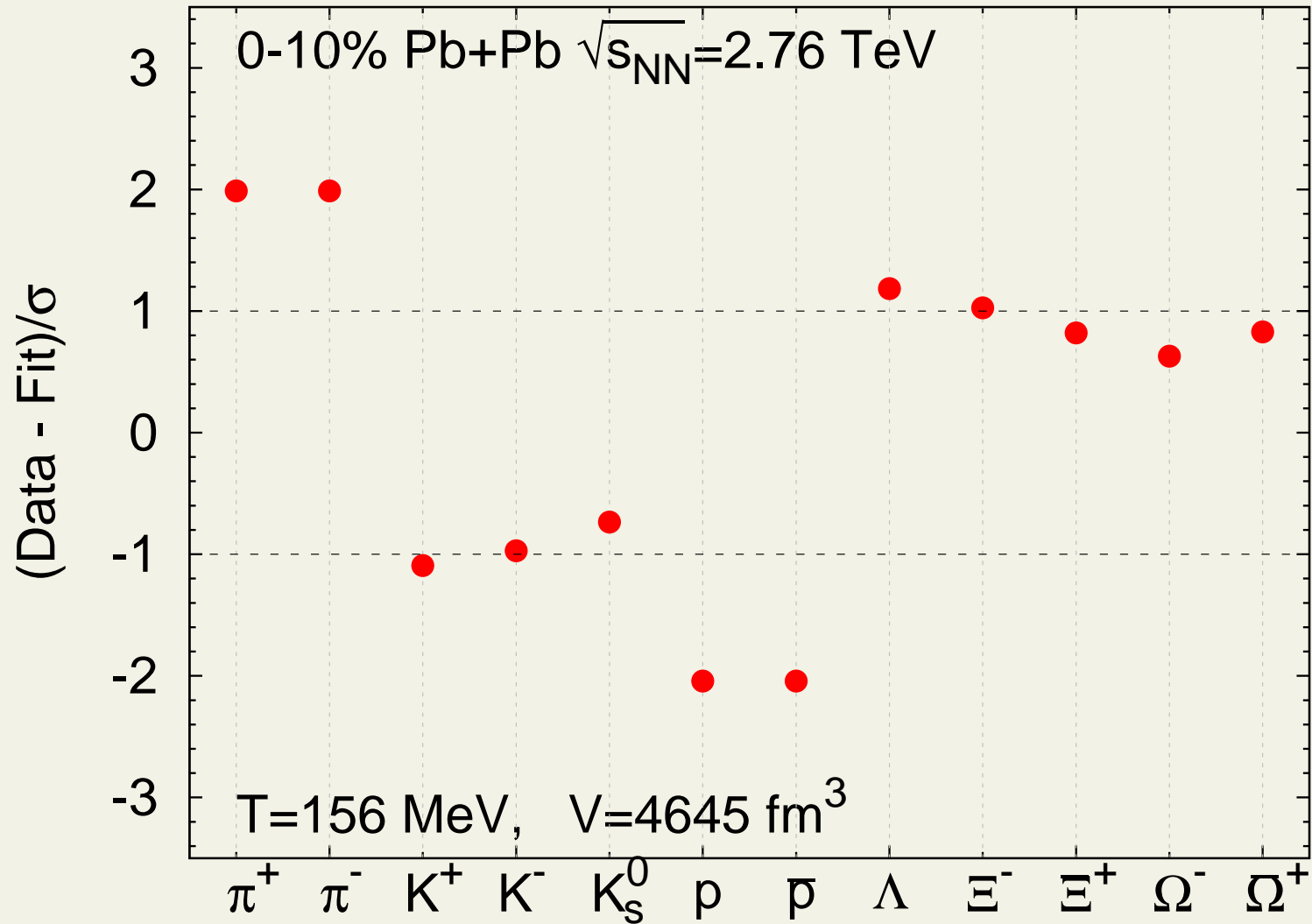
Centrality dependence



Centrality dependence



Fit to ALICE data



● how to include nonequilibrium effects in thermal model?

Partial chemical equilibrium

Bebie *et al.*, NPB378, 95 (1992)

After **loss of chemical equilibrium at T_1** , resonance formation/decay continues:



These processes in equilibrium if

$$\mu_\rho = 2\mu_\pi, \quad \mu_\Delta = \mu_\pi + \mu_N, \quad \mu_{K^*} = \mu_K + \mu_\pi, \dots$$

Conservation of particle # \Rightarrow buildup of chemical potentials

Entropy conservation \Rightarrow

$$\frac{\hat{n}_i(T, \mu_i)}{s(T, \{\mu_i\})} = \frac{\hat{n}_i(T_1)}{s(T_1)}$$

fixes the chemical potentials

Note that $\hat{n}_\pi = n_\pi + 2n_\rho + n_\Delta + n_{K^*} + \dots$

Partial chemical equilibrium

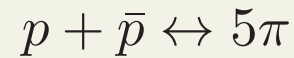
After loss of chemical equilibrium:

System characterised by

- T_1 , temperature where equilibrium is lost
- T_2 , temperature of final freeze-out
- V , volume at T_2

Partial chemical equilibrium

Assume $N\bar{N}$ -annihilations continue after loss of equilibrium at T_1 :



Equilibrium requires

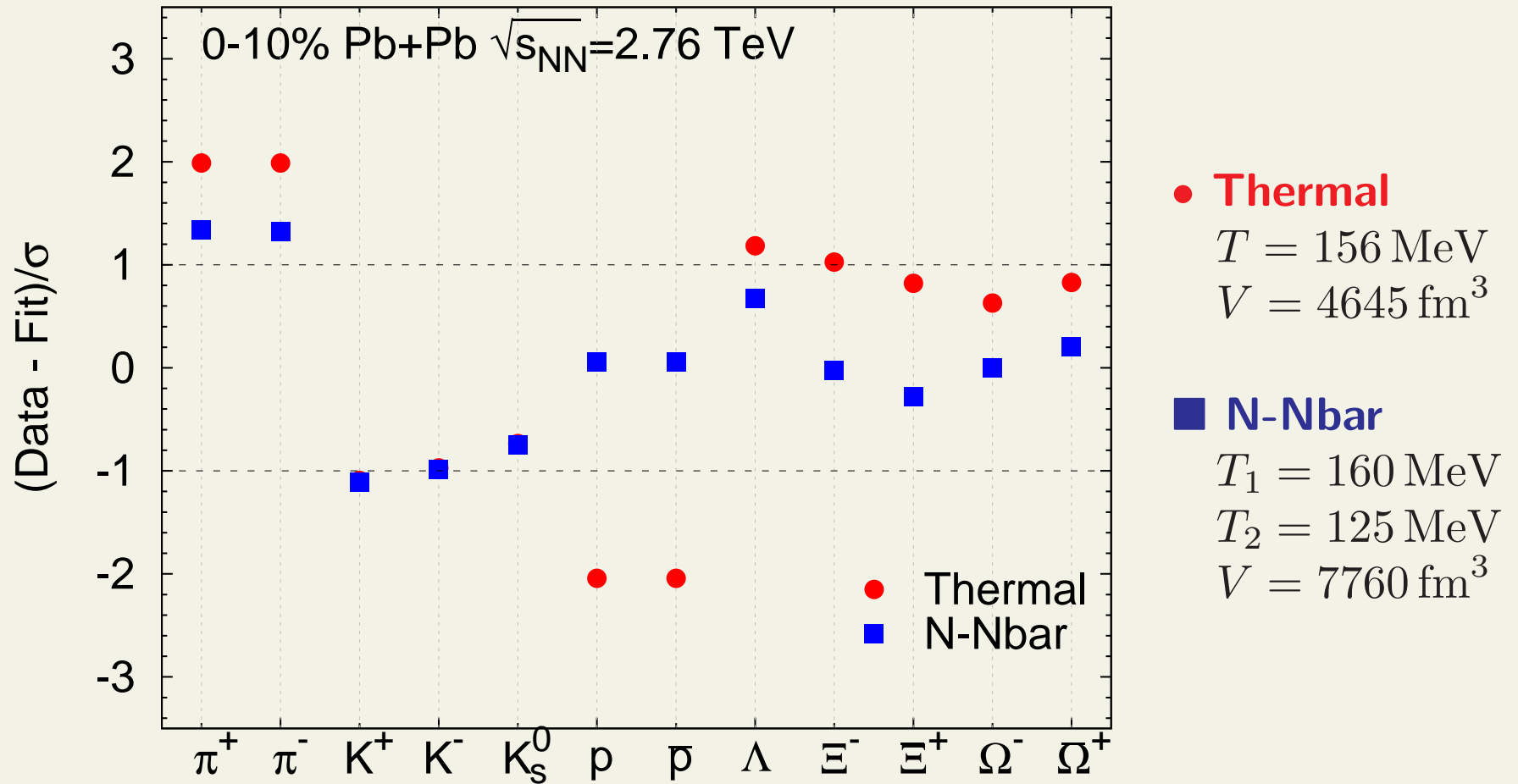
$$\mu_p + \mu_{\bar{p}} = 5\mu_\pi \quad \Leftrightarrow \quad \mu_p = \mu_{\bar{p}} = 2.5\mu_\pi$$

- \hat{N}_N , $\hat{N}_{\bar{N}}$ or \hat{N}_π **not conserved**
- $N_\pi + 2.5N_N + 2.5N_{\bar{N}}$ **is conserved**

$$\frac{\hat{n}_\pi(T, \mu_\pi) + 2.5\hat{n}_N(T, \mu_\pi) + 2.5\hat{n}_{\bar{N}}(T, \mu_\pi)}{s(T, \{\mu_i\})} = \text{const.}$$

gives $\mu_\pi(T)$

Fit allowing $N\bar{N}$ -annihilations

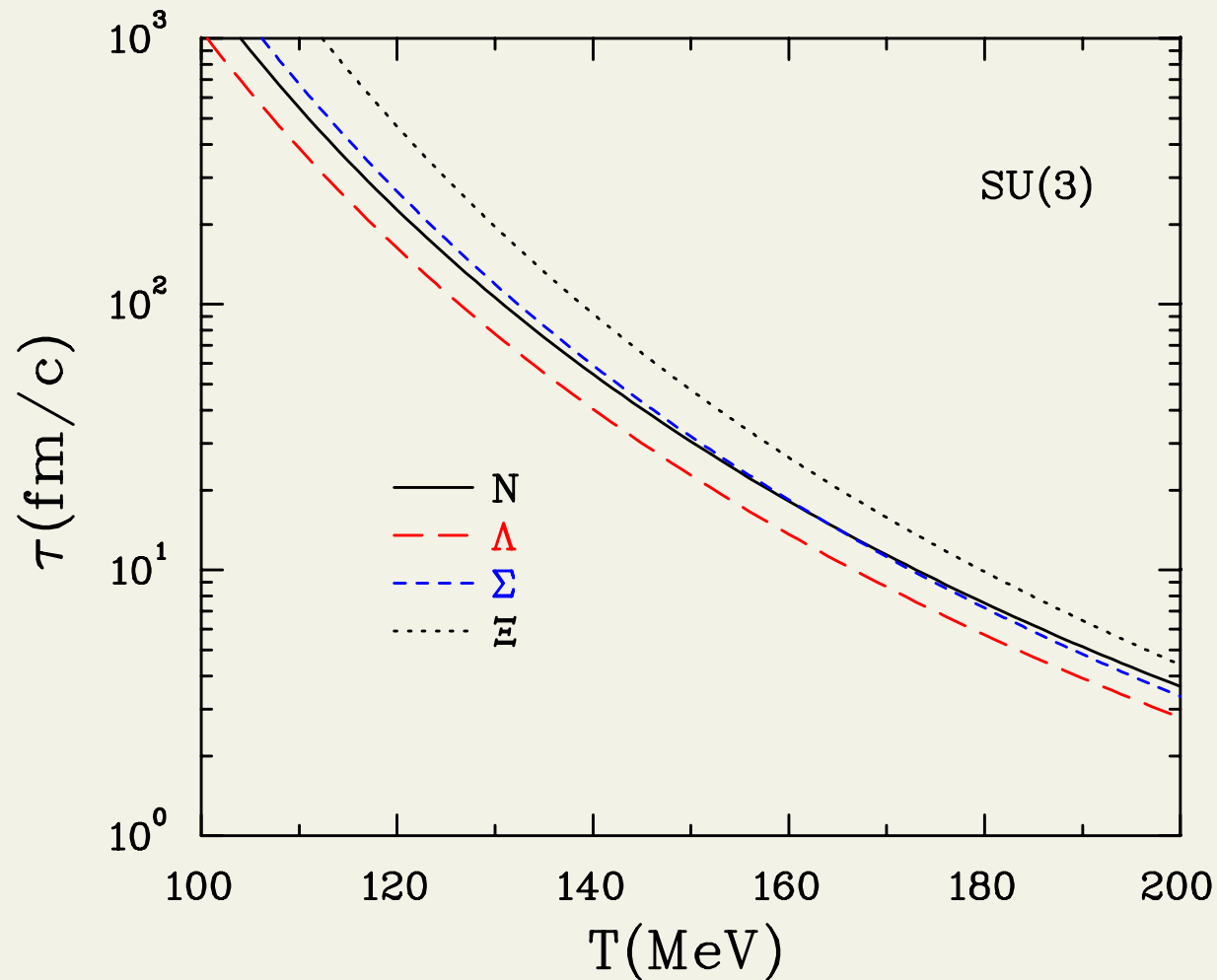


● very good fit!

● but. . .

Chemical equilibration time for nucleons

Kapusta & Shovkovy, Phys. Rev. C68, 014901 (2003)



- Λ 's equilibrate faster than nucleons
 - Σ 's equilibrate as fast as nucleons
- \Rightarrow include Λ and Σ annihilations

Λ - and Σ -annihilations

Processes



in equilibrium when

$$\mu_{\Lambda} = \mu_{\Sigma} = 2.5\mu_{\pi}$$

But



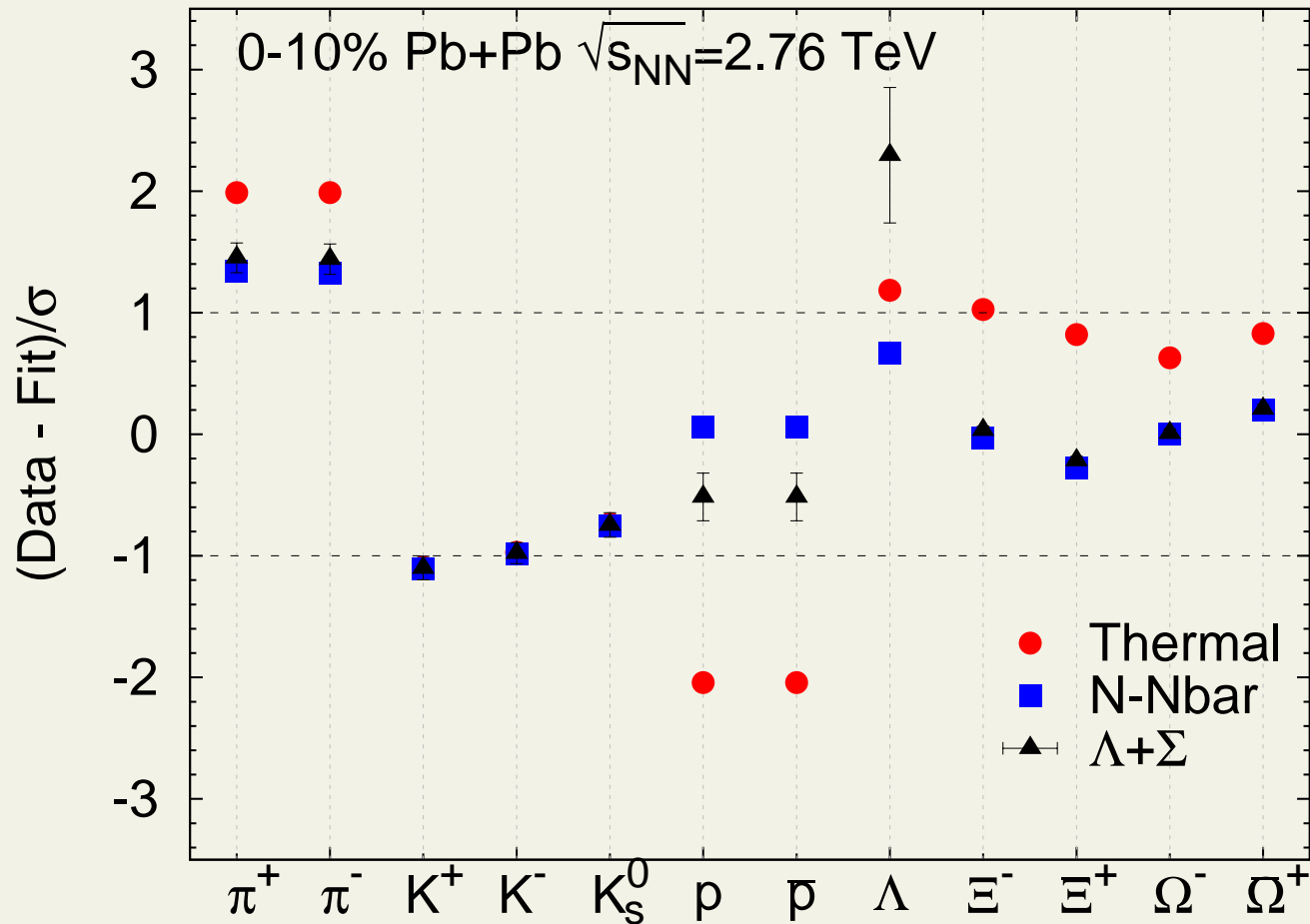
in equilibrium when

$$\mu_{\Lambda} = \mu_{\Sigma} = 1.5\mu_{\pi} + \mu_{\bar{K}}$$

\Rightarrow Weighted average with rates as weight

rates unknown \Rightarrow upper and lower limit

Fit including Λ - and Σ -annihilations



- **Thermal**

$$T = 156 \text{ MeV}$$

$$V = 4645 \text{ fm}^3$$

- **N-Nbar**

$$T_1 = 160 \text{ MeV}$$

$$T_2 = 125 \text{ MeV}$$

$$V = 7760 \text{ fm}^3$$

- ▲ **$\Lambda+\Sigma$**

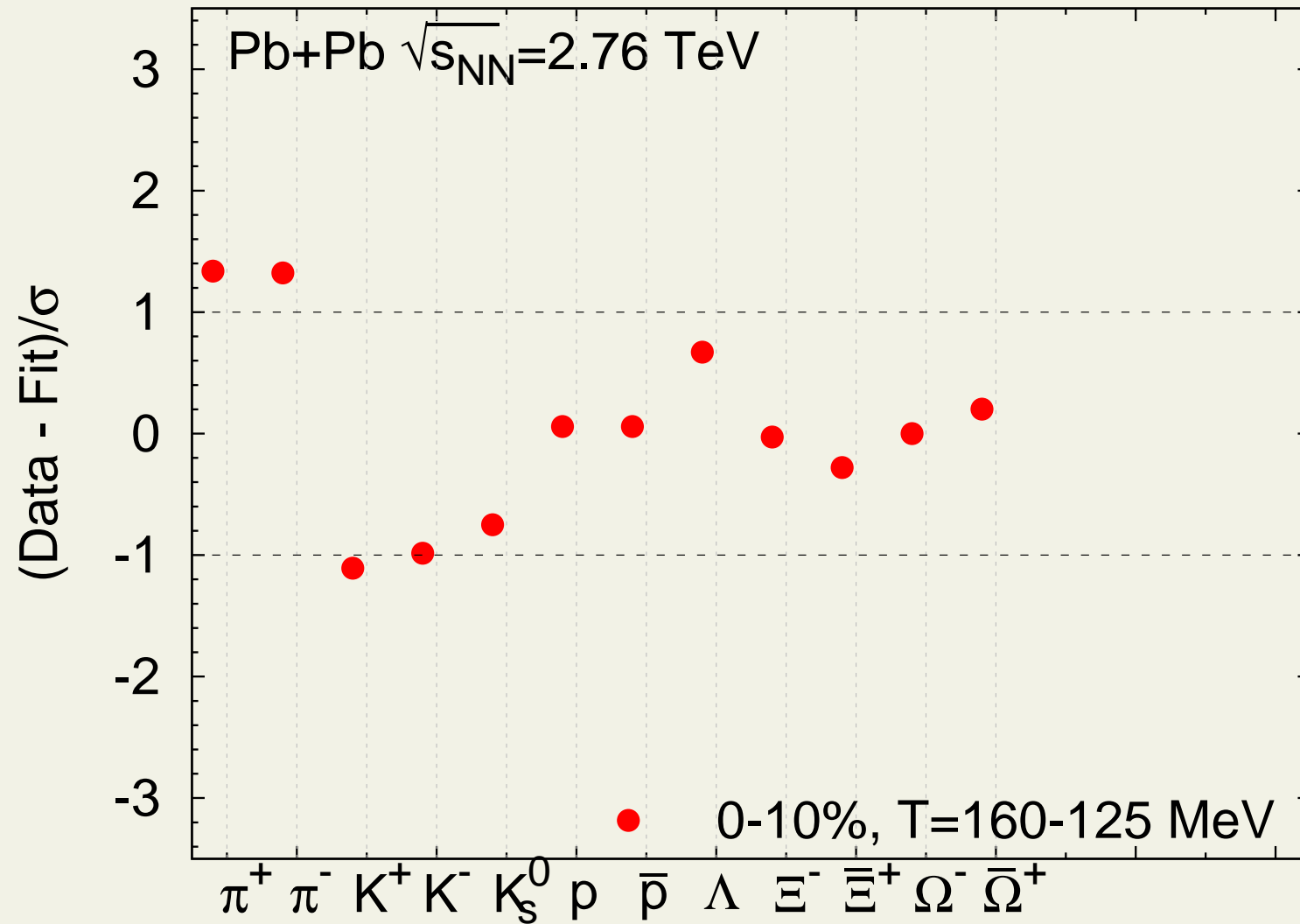
$$T_1 = 159 - 161 \text{ MeV}$$

$$T_2 = 133 - 130 \text{ MeV}$$

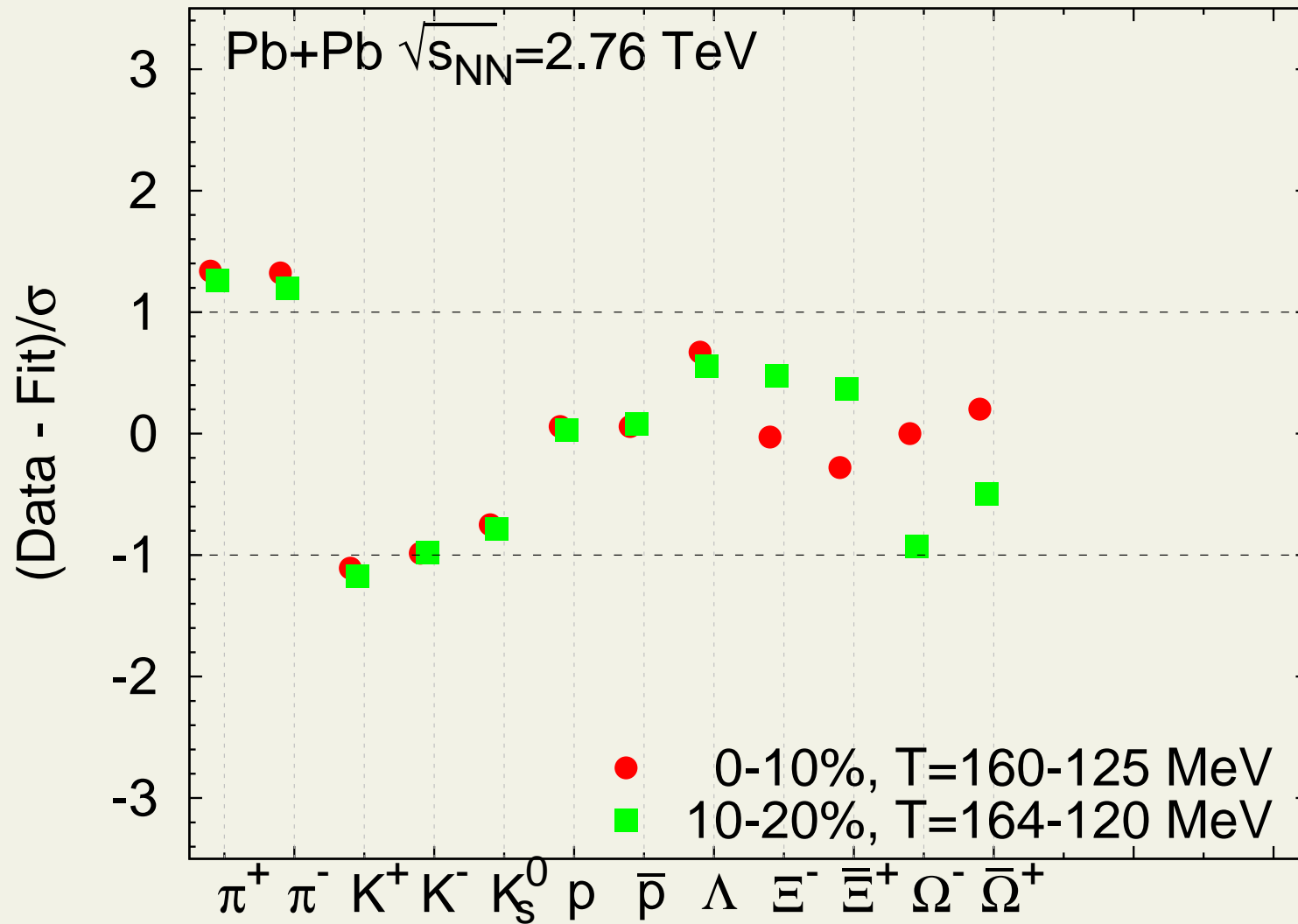
$$V = 6750 - 6650 \text{ fm}^3$$

- **serious lack of Λ 's!**

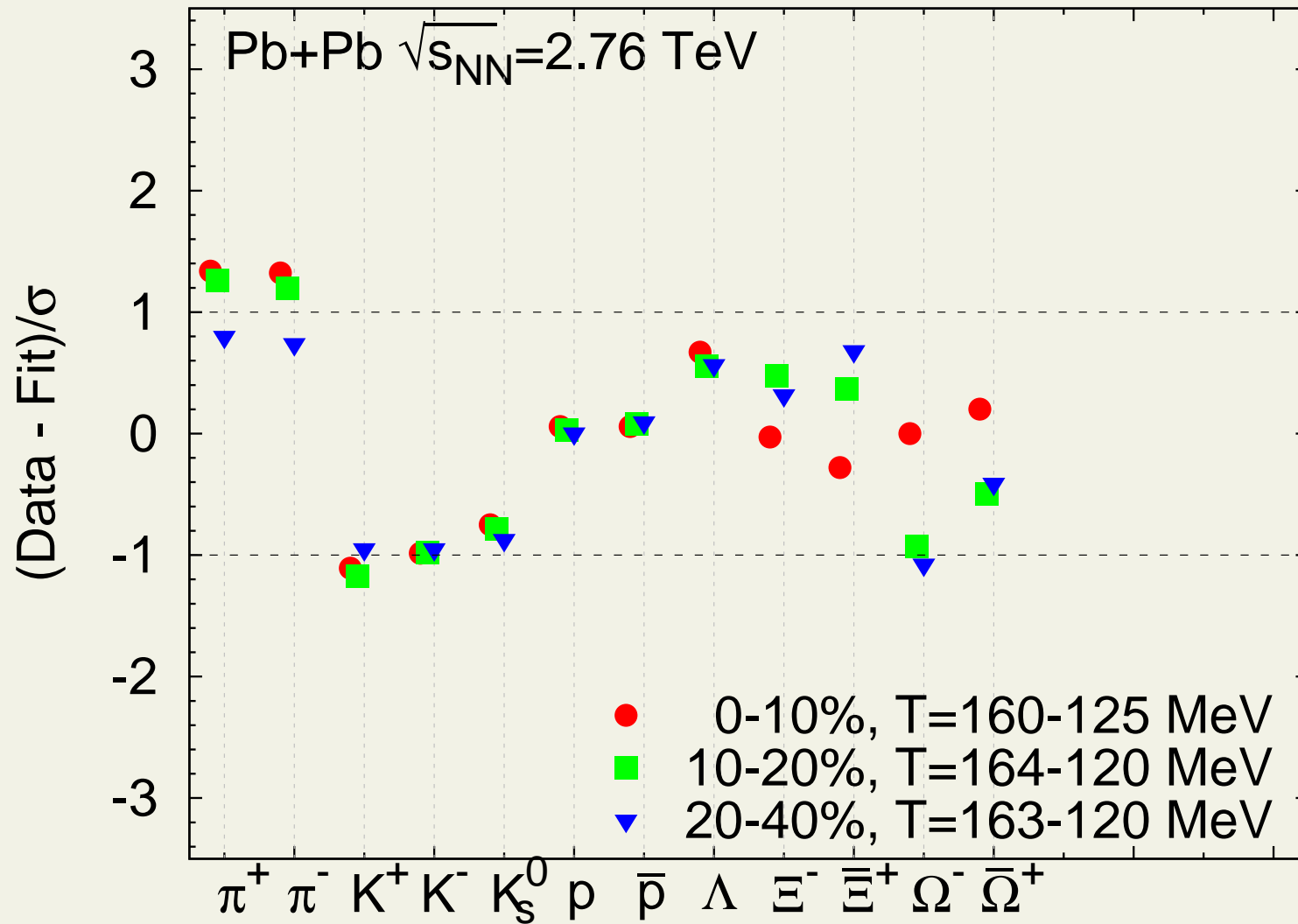
Centrality dependence



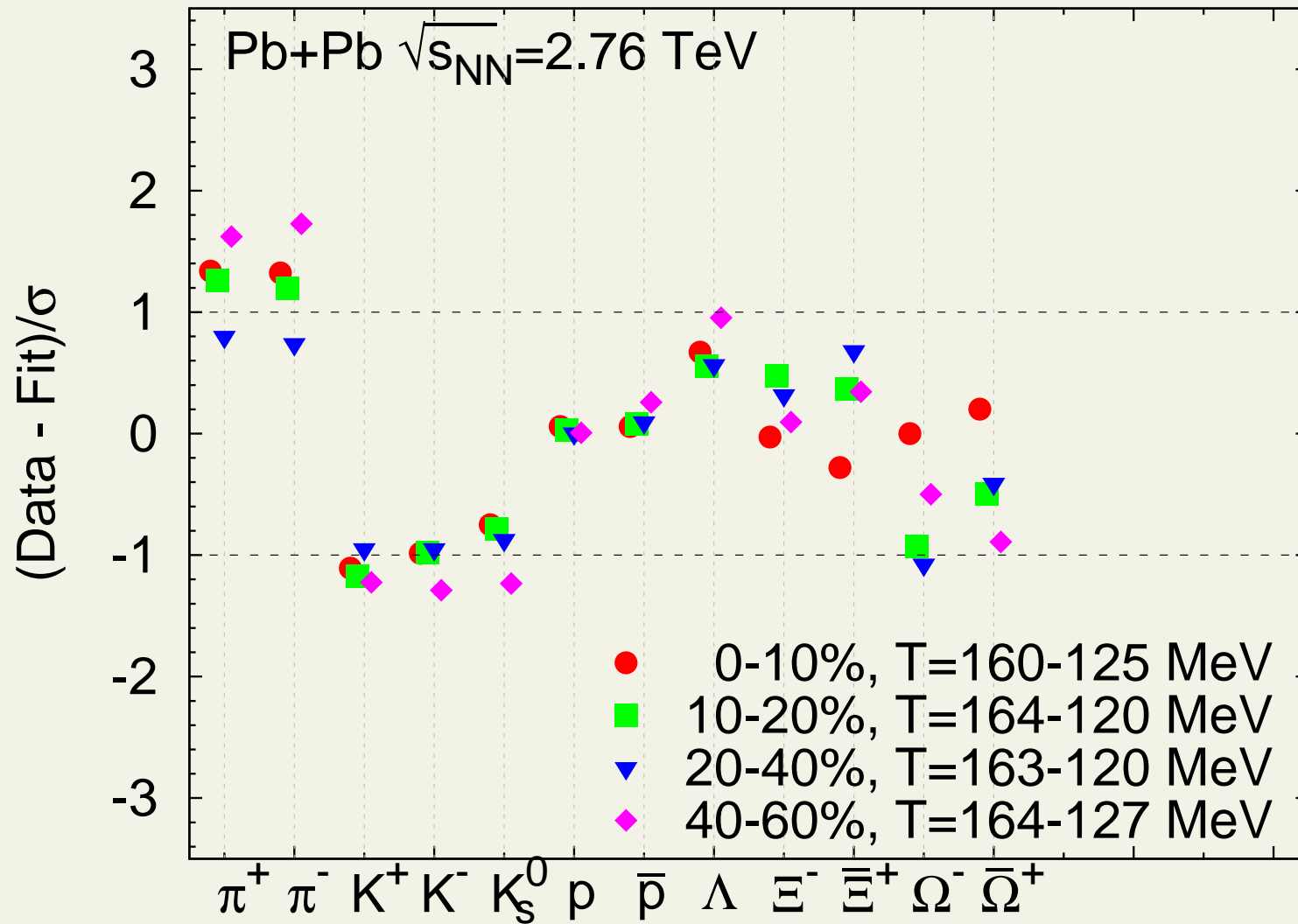
Centrality dependence



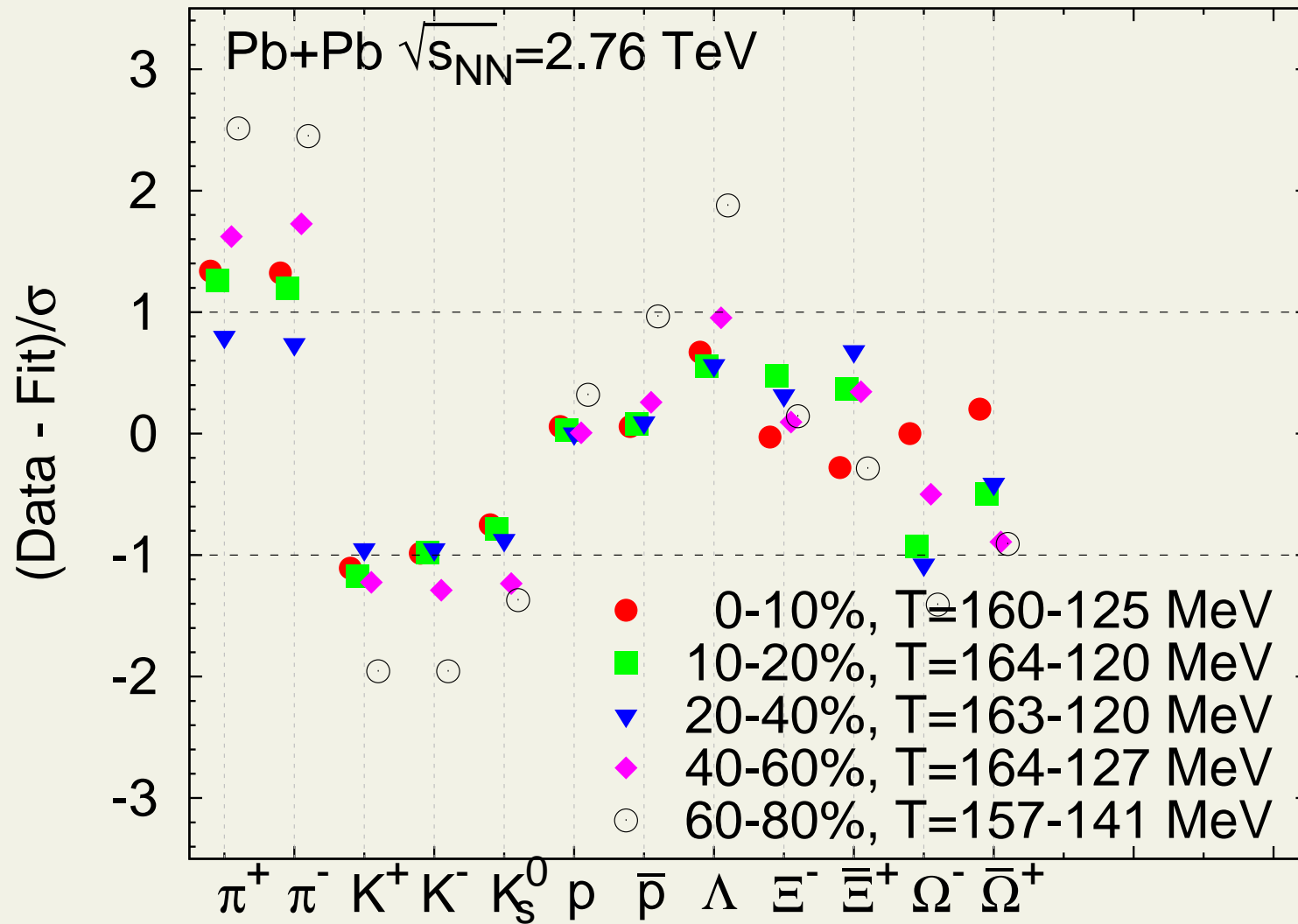
Centrality dependence



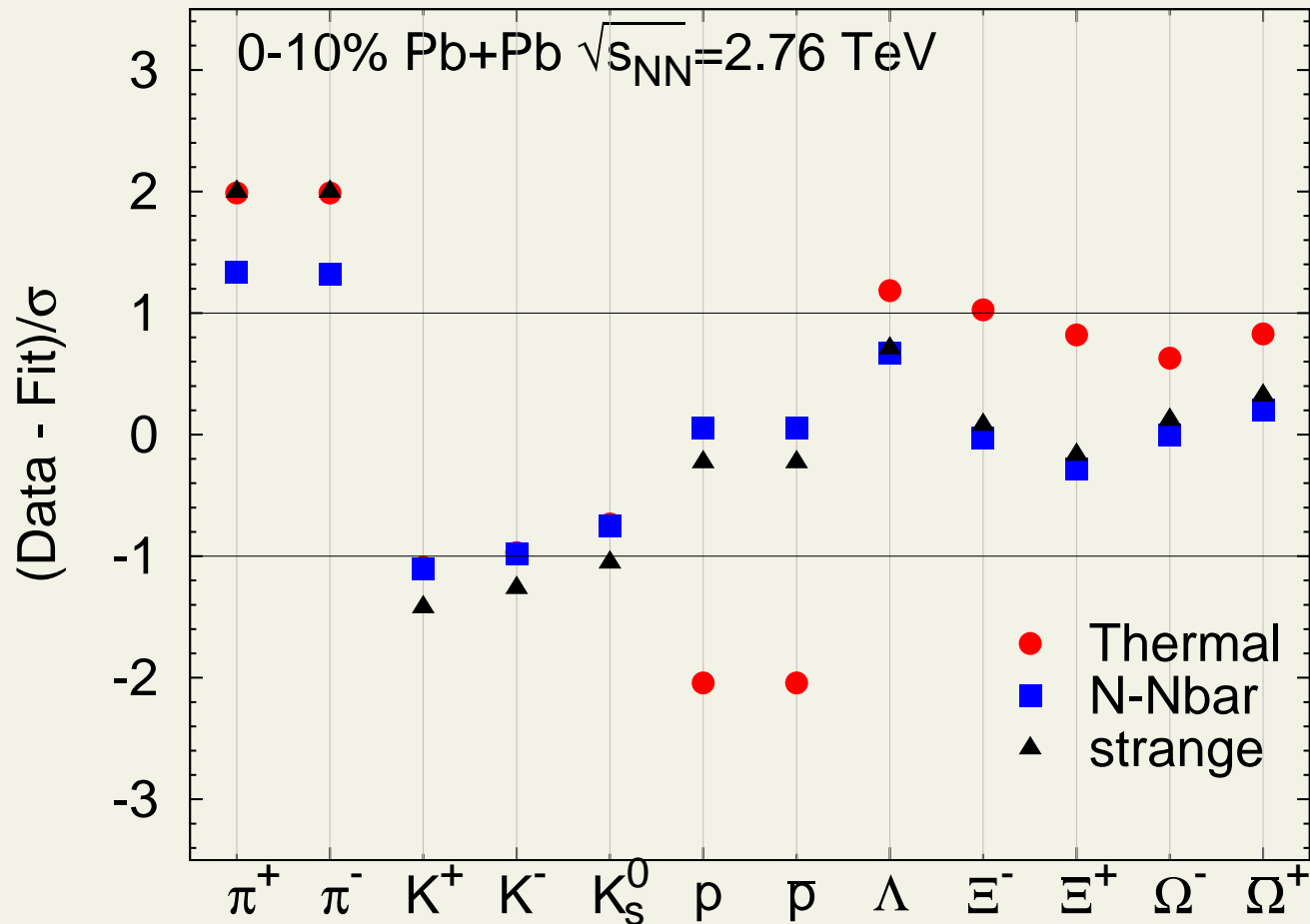
Centrality dependence



Centrality dependence



What if strange particles freeze-out before non-strange?



- **Thermal**

$$T = 156 \text{ MeV}$$

$$V = 4645 \text{ fm}^3$$

- **N-Nbar**

$$T_1 = 160 \text{ MeV}$$

$$T_2 = 125 \text{ MeV}$$

$$V = 7760 \text{ fm}^3$$

- ▲ **strange**

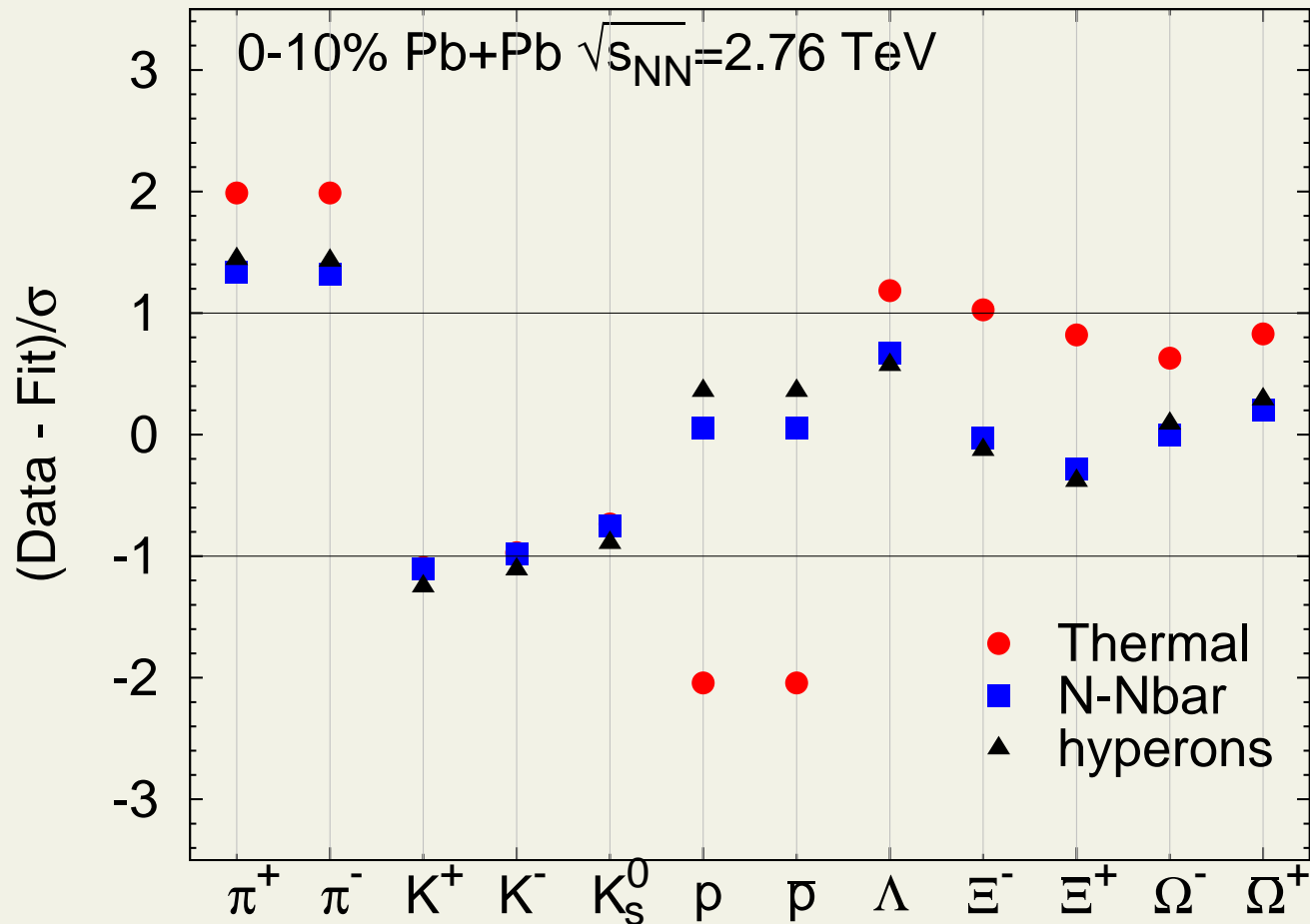
$$T_{\text{strange}} = 158 \text{ MeV}$$

$$T_{\text{rest}} = 149 \text{ MeV}$$

$$V = 6230 \text{ fm}^3$$

- **K/π -ratio ?**

What if hyperons freeze-out before the rest?



- **Thermal**

$$T = 156 \text{ MeV}$$

$$V = 4645 \text{ fm}^3$$

- **N-Nbar**

$$T_1 = 160 \text{ MeV}$$

$$T_2 = 125 \text{ MeV}$$

$$V = 7760 \text{ fm}^3$$

- ▲ **hyperons**

$$T_{\text{hyperon}} = 157 \text{ MeV}$$

$$T_{\text{rest}} = 144 \text{ MeV}$$

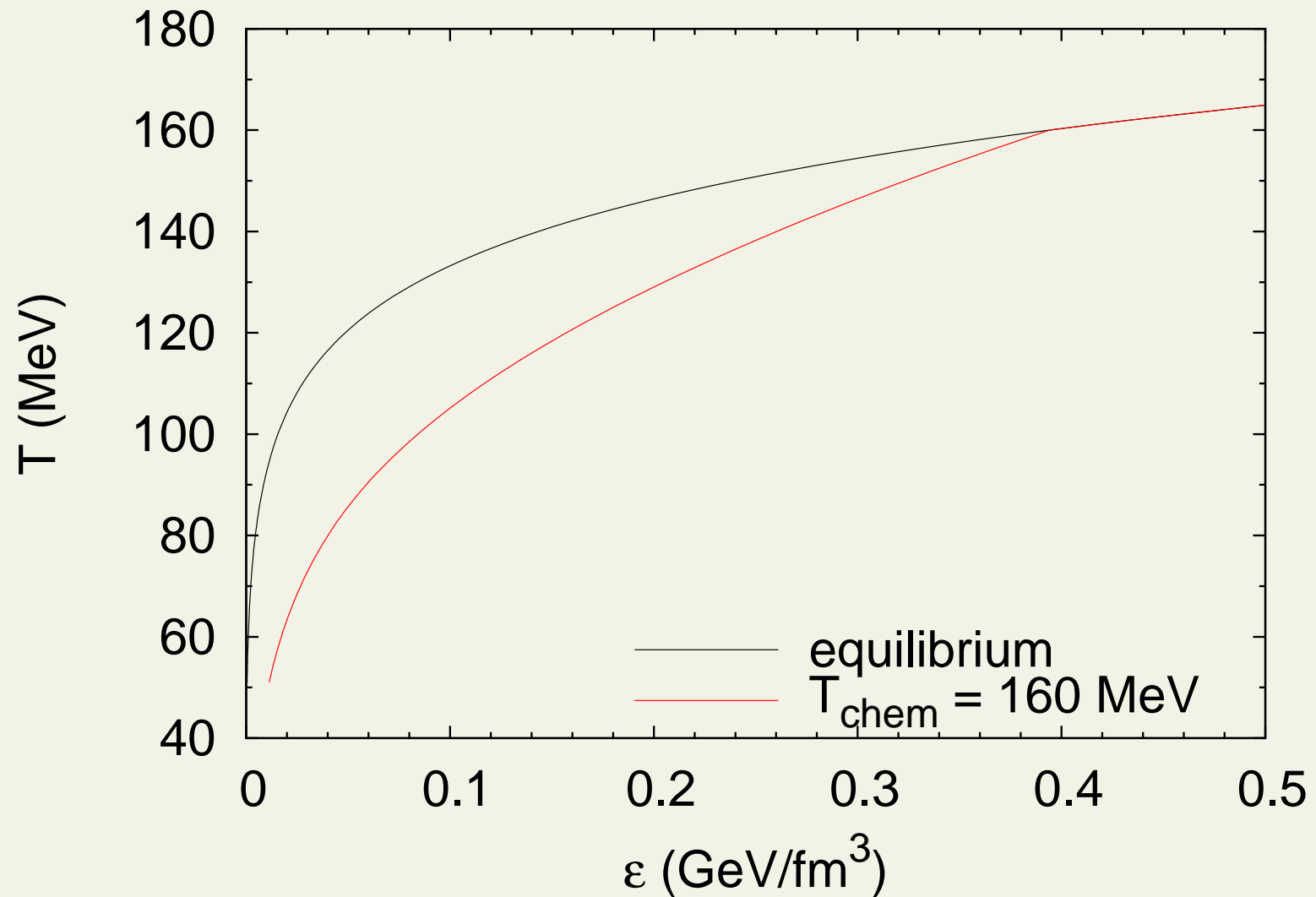
$$V = 7890 \text{ fm}^3$$

- equal to N-Nbar fit

Equation of state

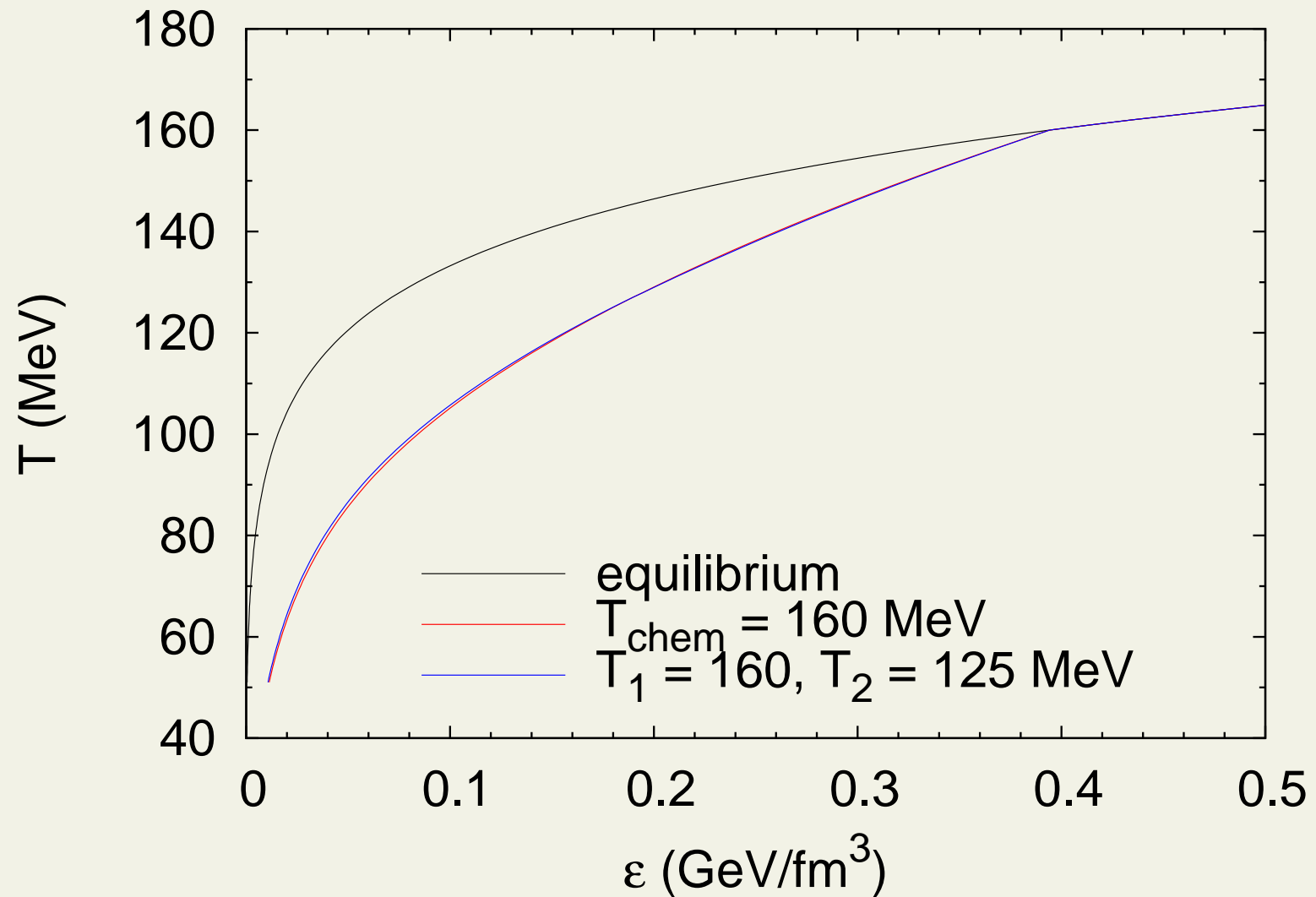
- built as conventional PCE EoS
- loss of chemical equilibrium at T_1
- conserved species change at T_2

Equation of state



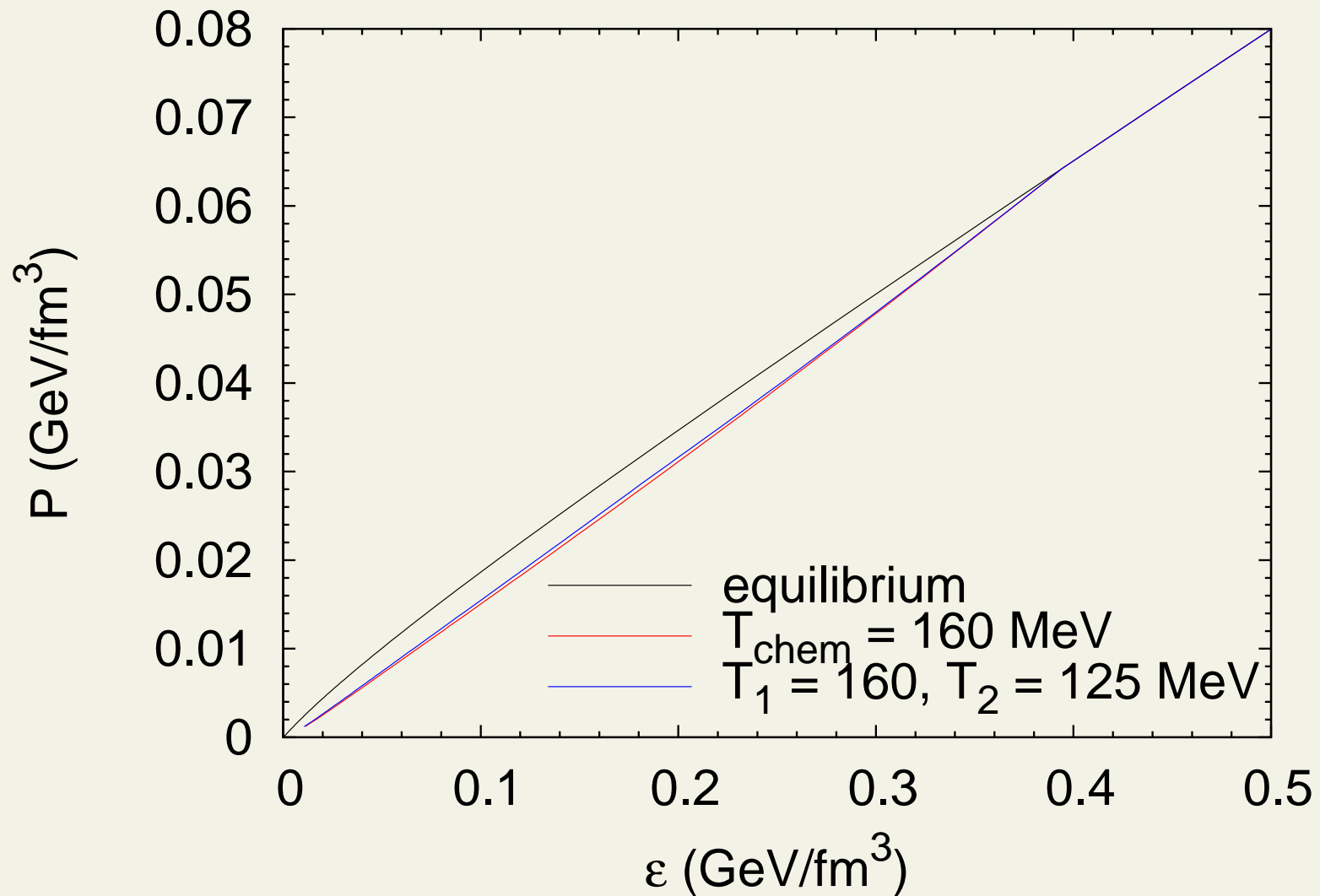
chemical freeze-out \Rightarrow main effect via temperature

Equation of state



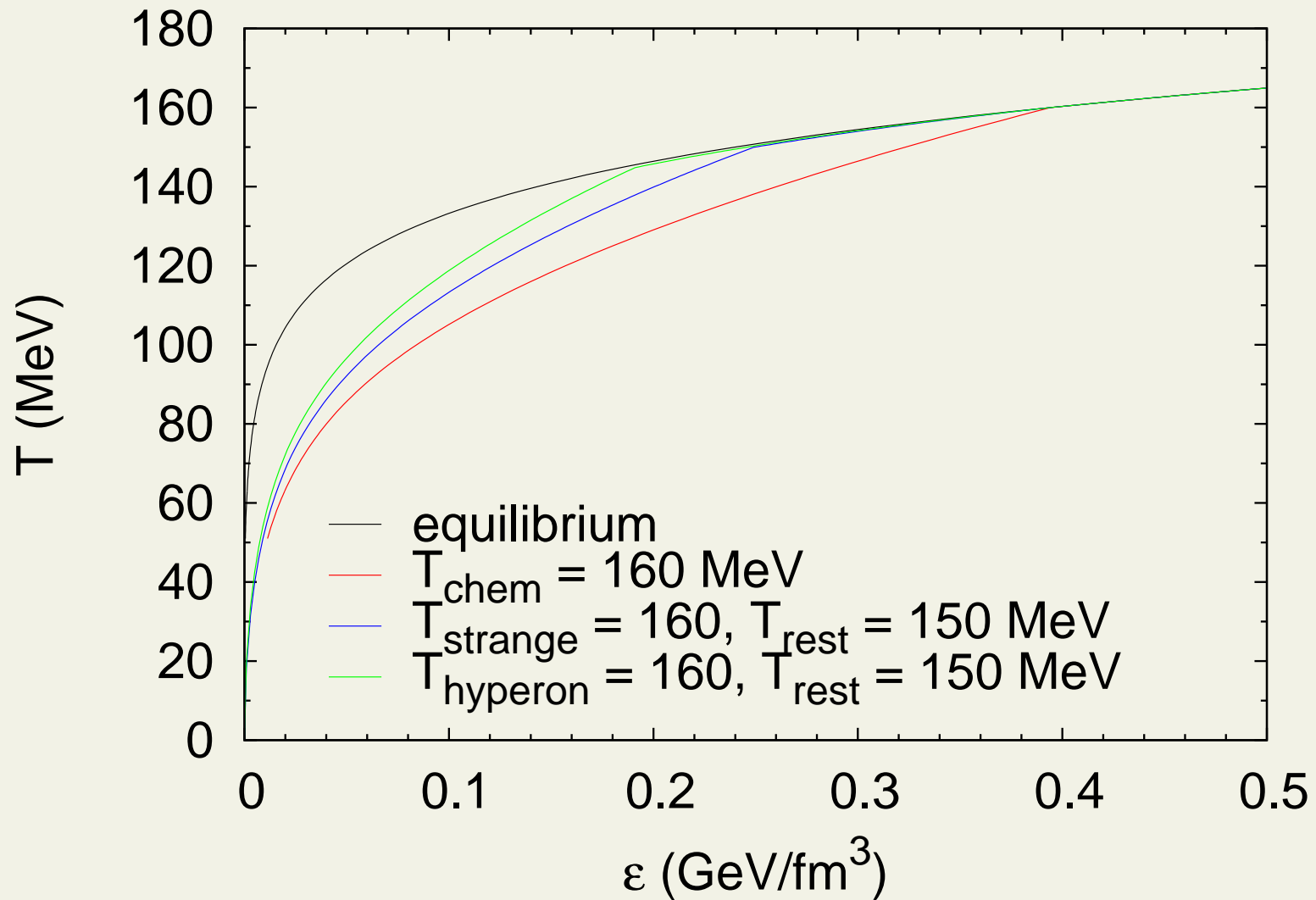
$N\bar{N}$ -annihilations hardly change the EoS

Equation of state



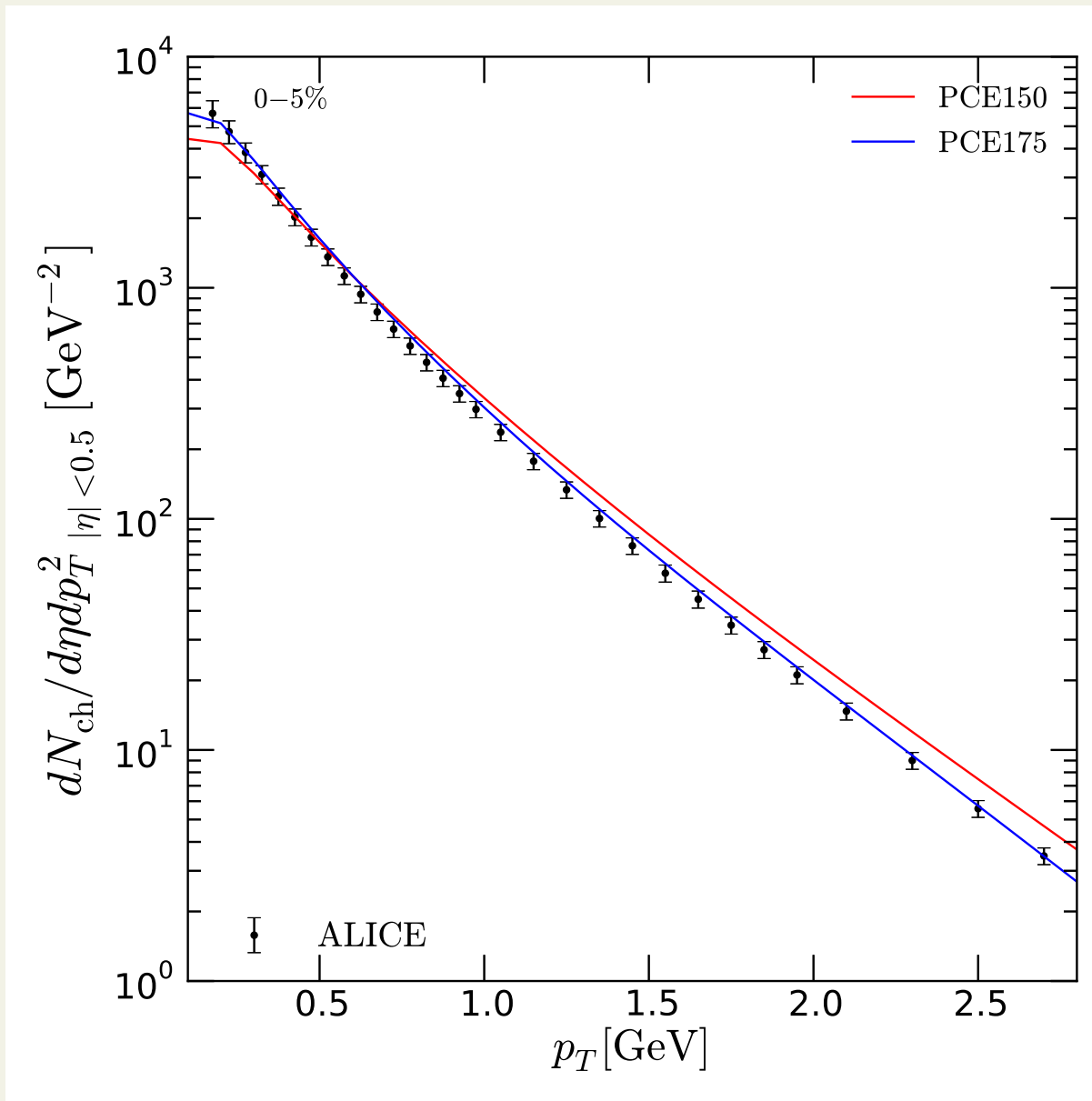
$N\bar{N}$ -annihilations hardly change the EoS

Equation of state



separate strange particle or hyperon freeze-out affects the EoS!

Charged hadron p_T spectrum at LHC

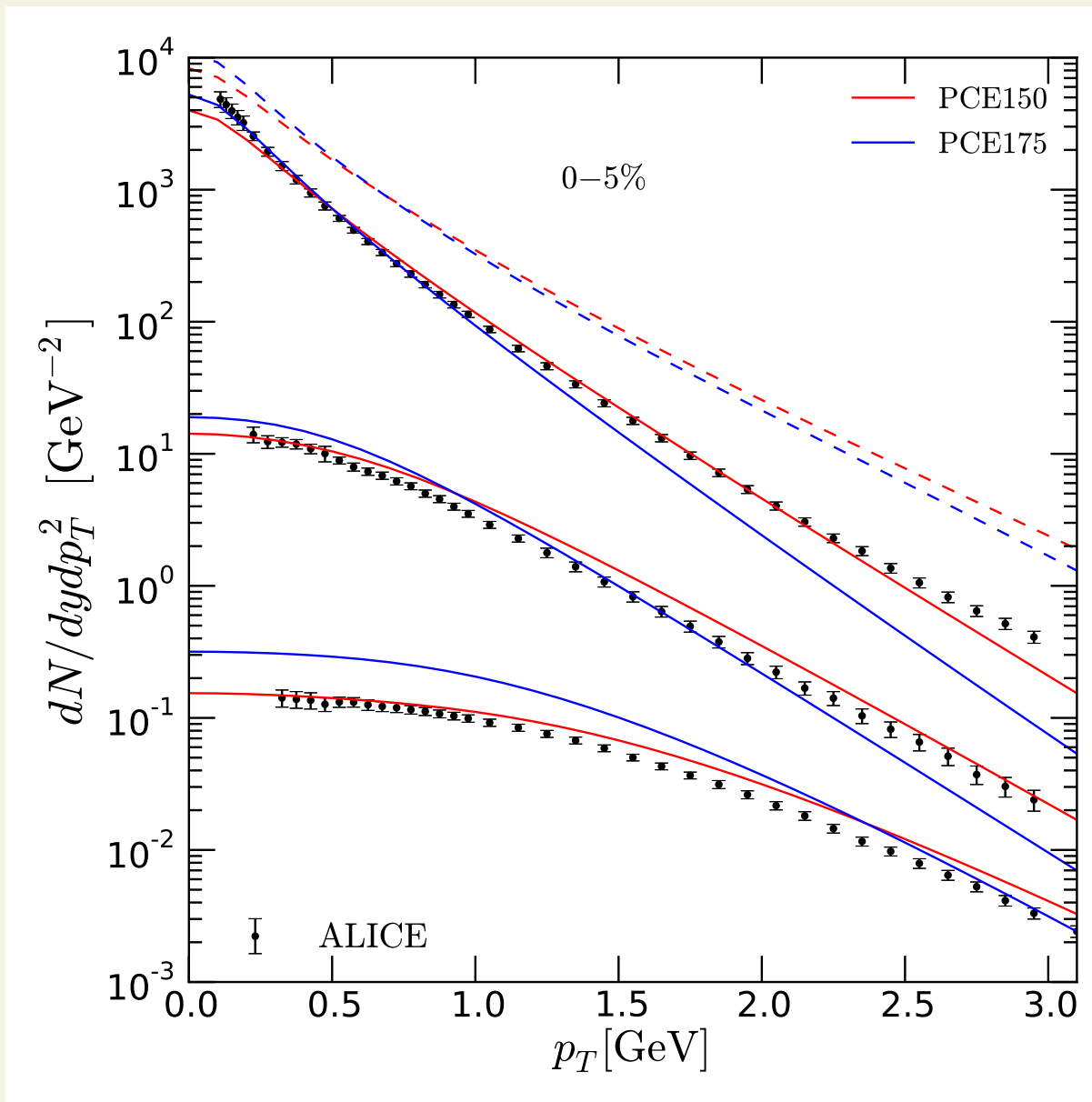


PCE150:
fit to π , K , p yields
no fit to spectrum

PCE175:
no fit to yields
fits the spectrum

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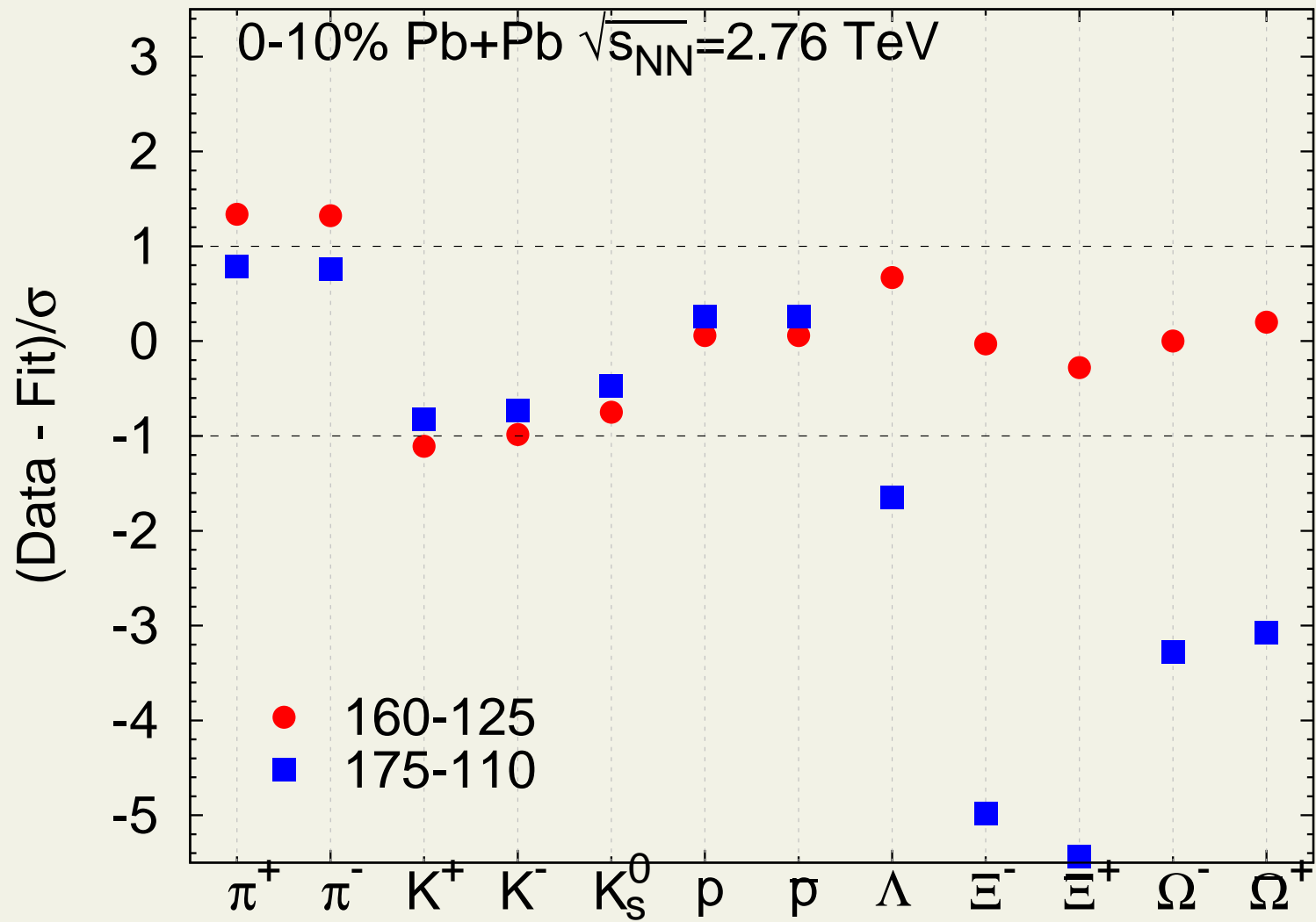
Identified particle p_T spectrum at LHC



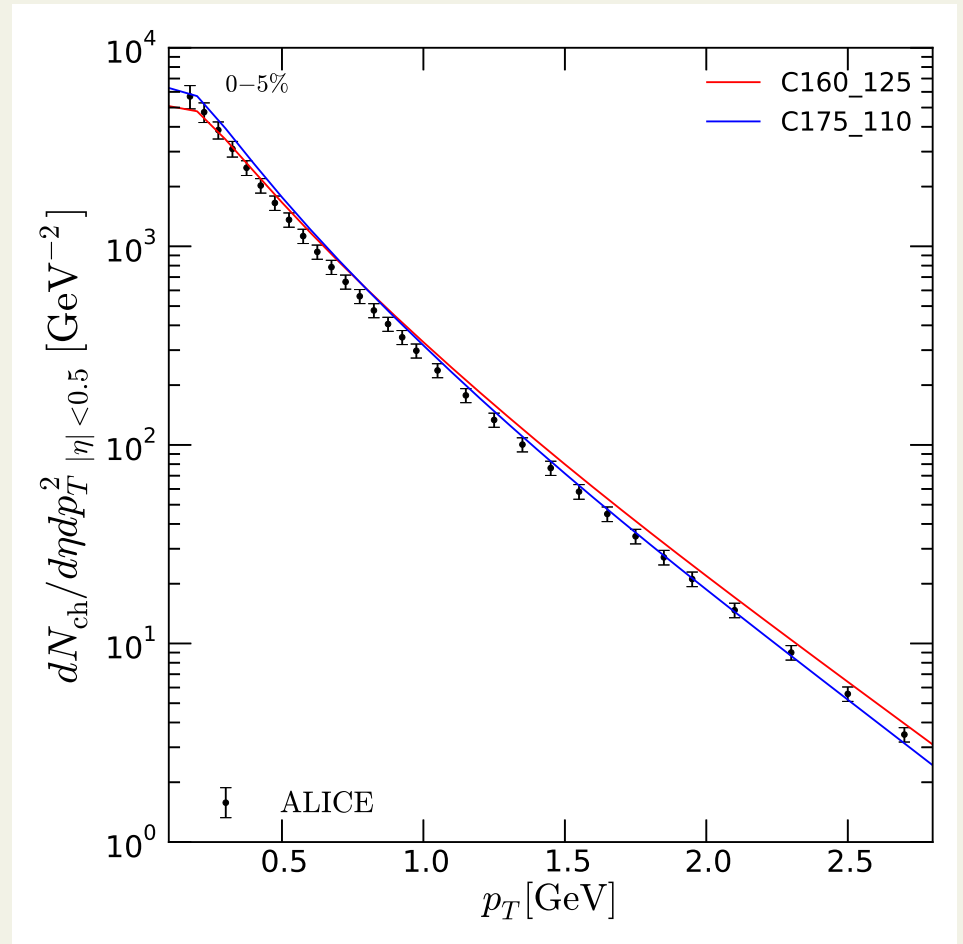
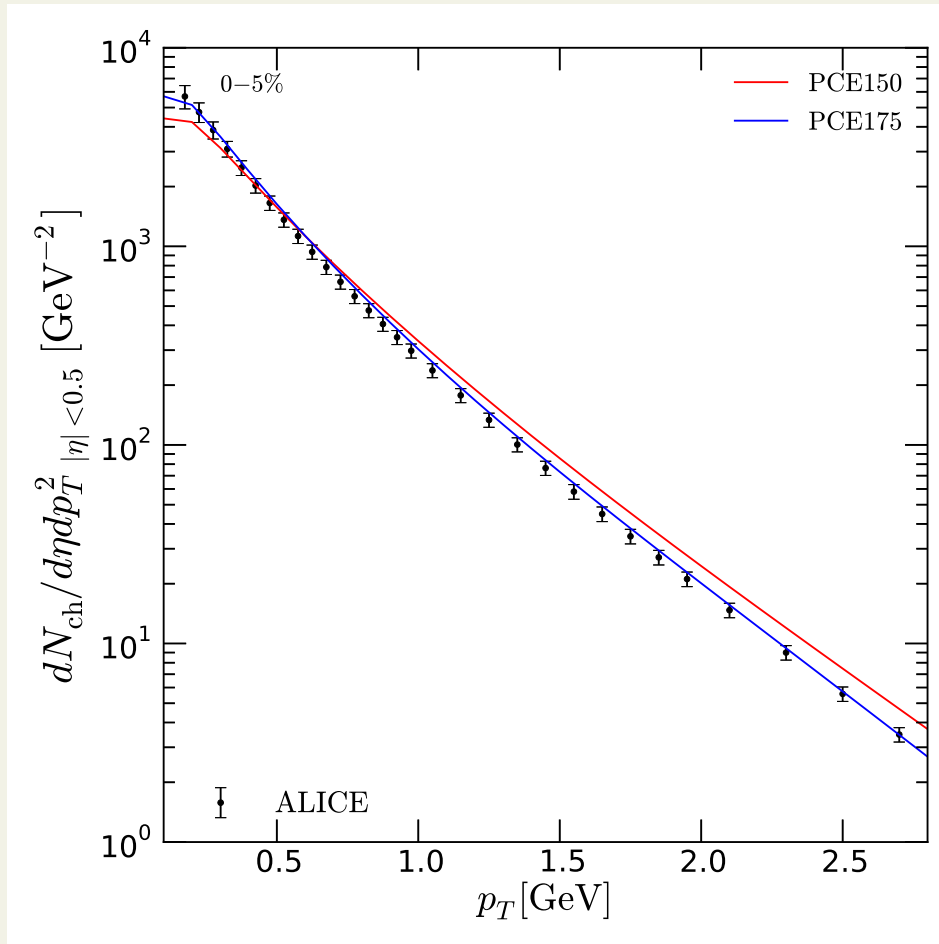
PCE150:
fits pions at high p_T

PCE175:
fits pions at low p_T

©H. Niemi



Charged hadron p_T spectrum at LHC



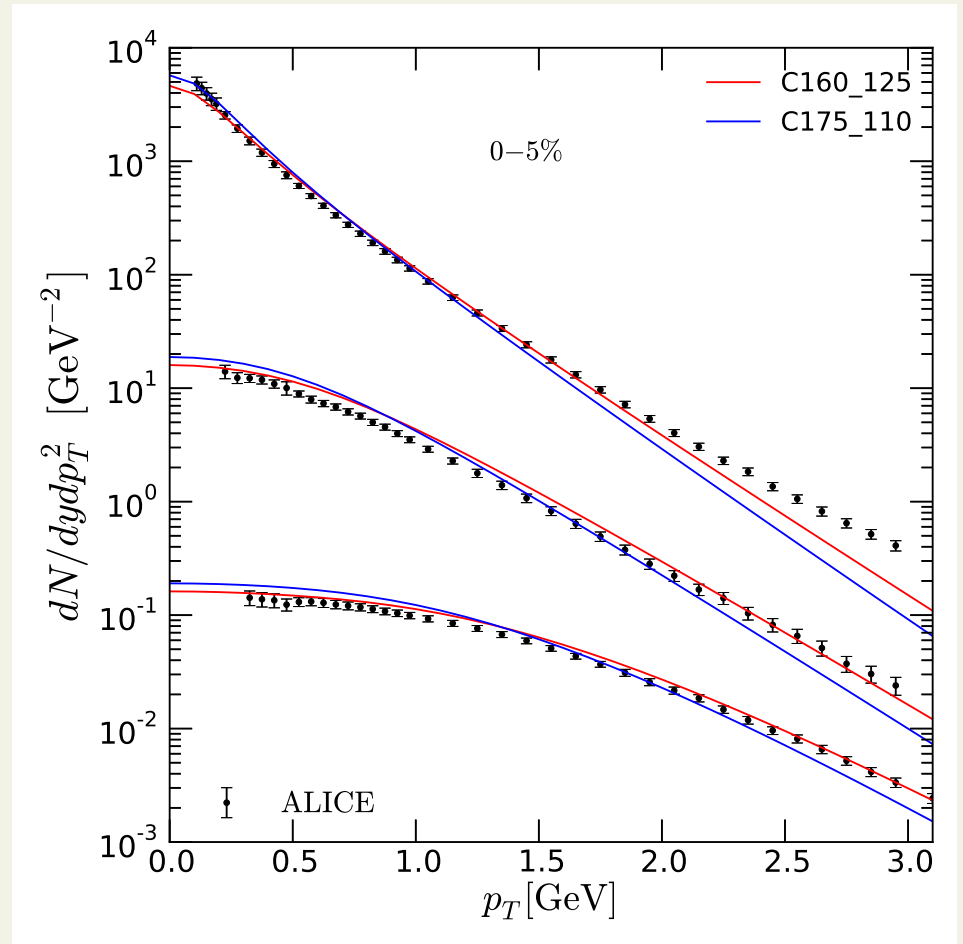
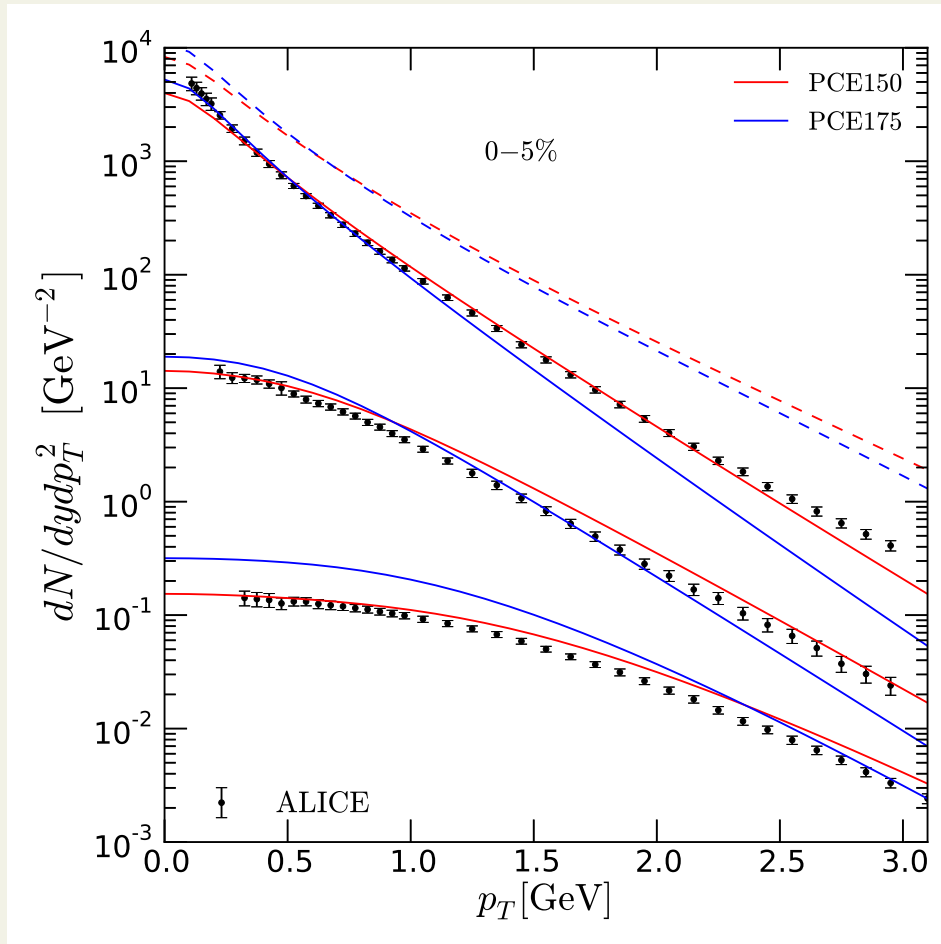
PCE150: $T_{\text{chem}} = 150 \text{ MeV}$

PCE175: $T_{\text{chem}} = 175 \text{ MeV}$

C160_125: $T_1 = 160 \text{ MeV}, T_2 = 125 \text{ MeV}$

C175_110: $T_1 = 175 \text{ MeV}, T_2 = 110 \text{ MeV}$

Identified particle p_T spectrum at LHC



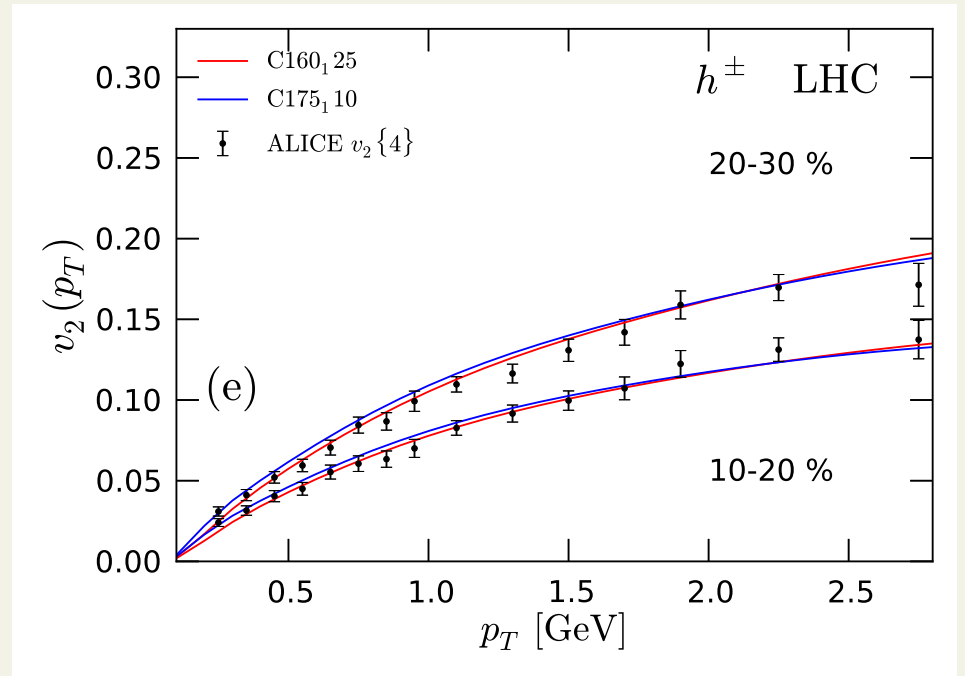
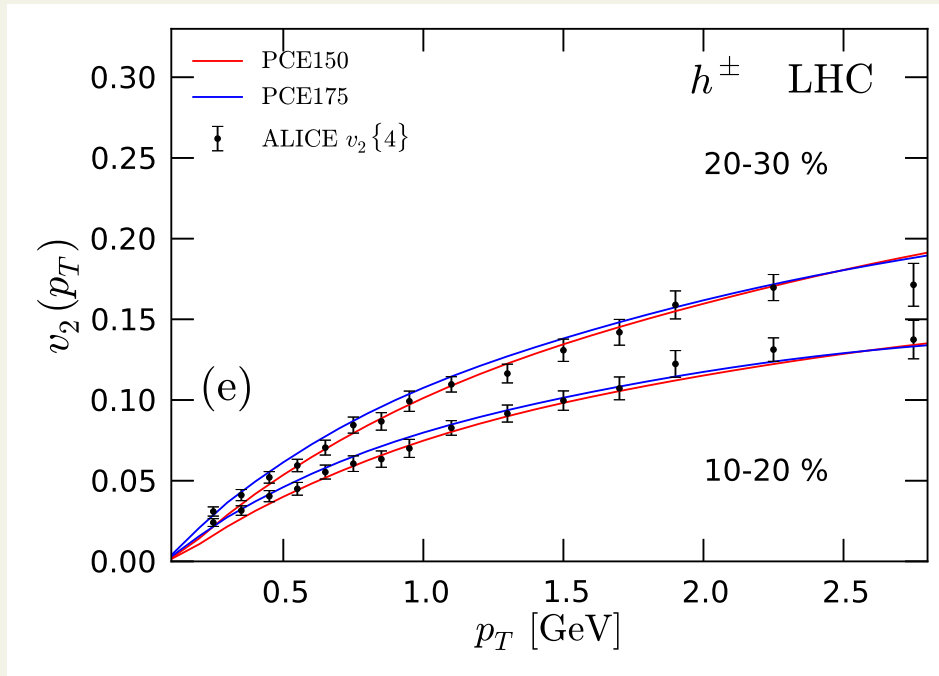
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$v_2(p_T)$ at LHC



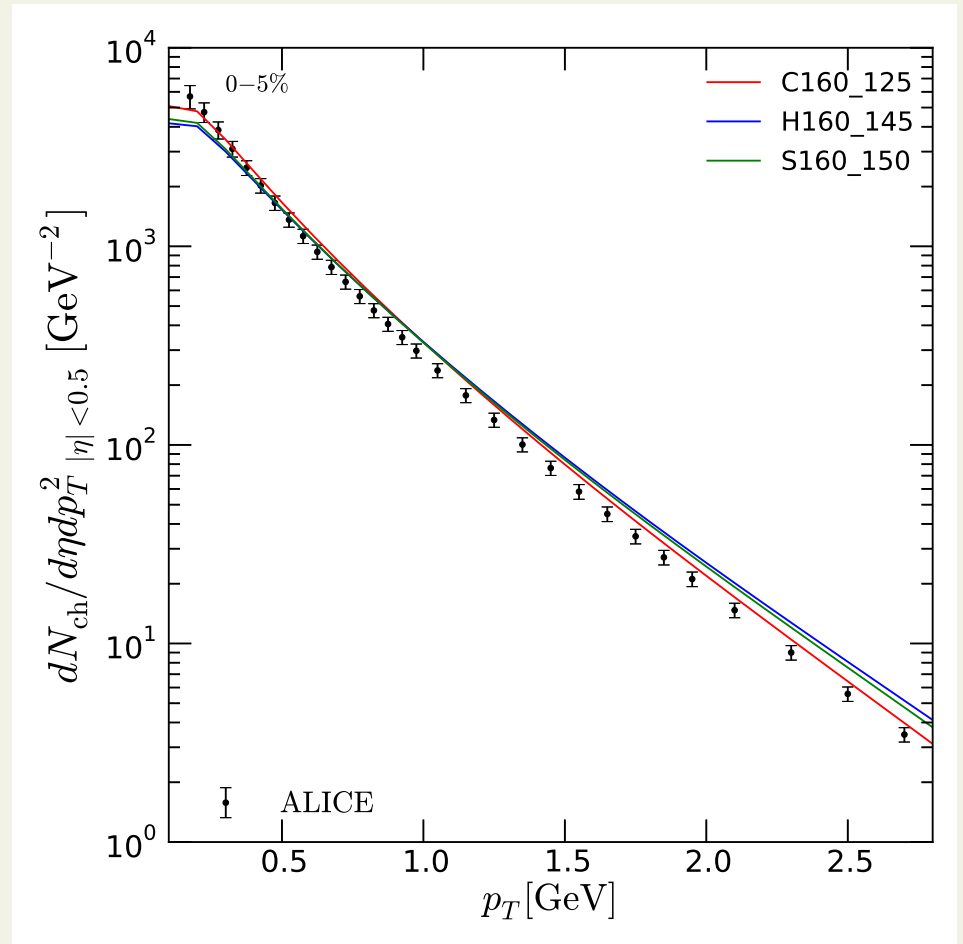
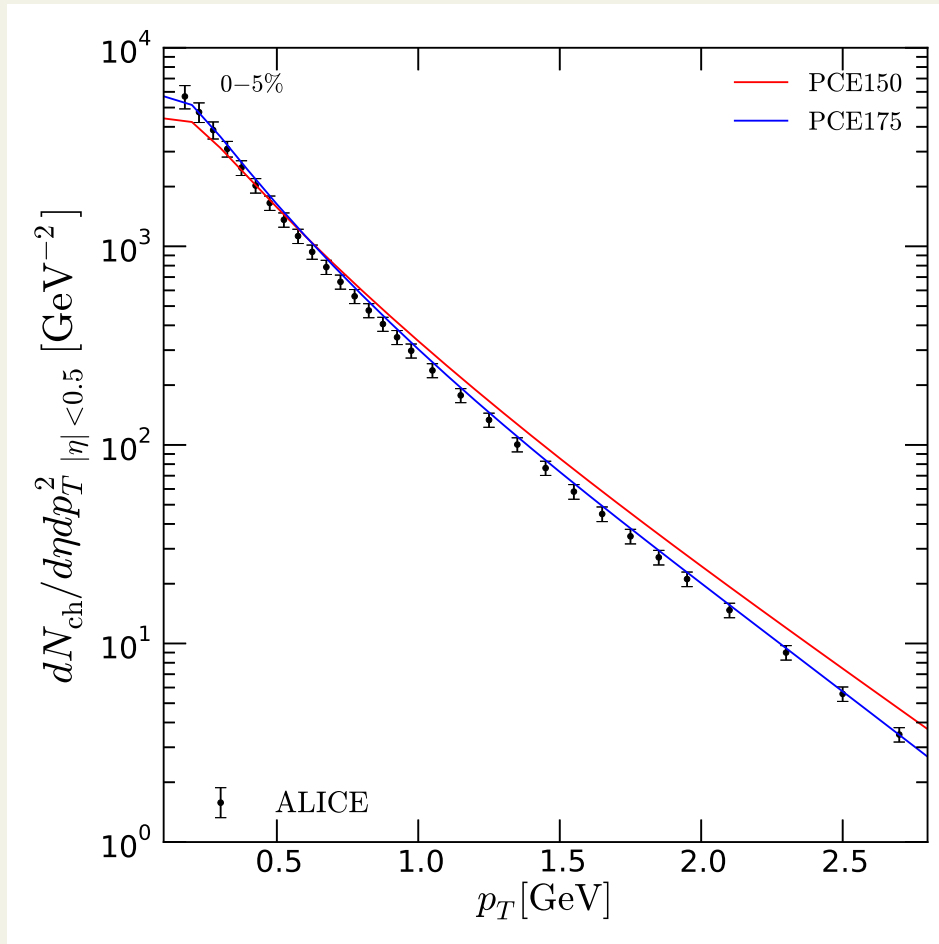
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Charged hadron p_T spectrum at LHC



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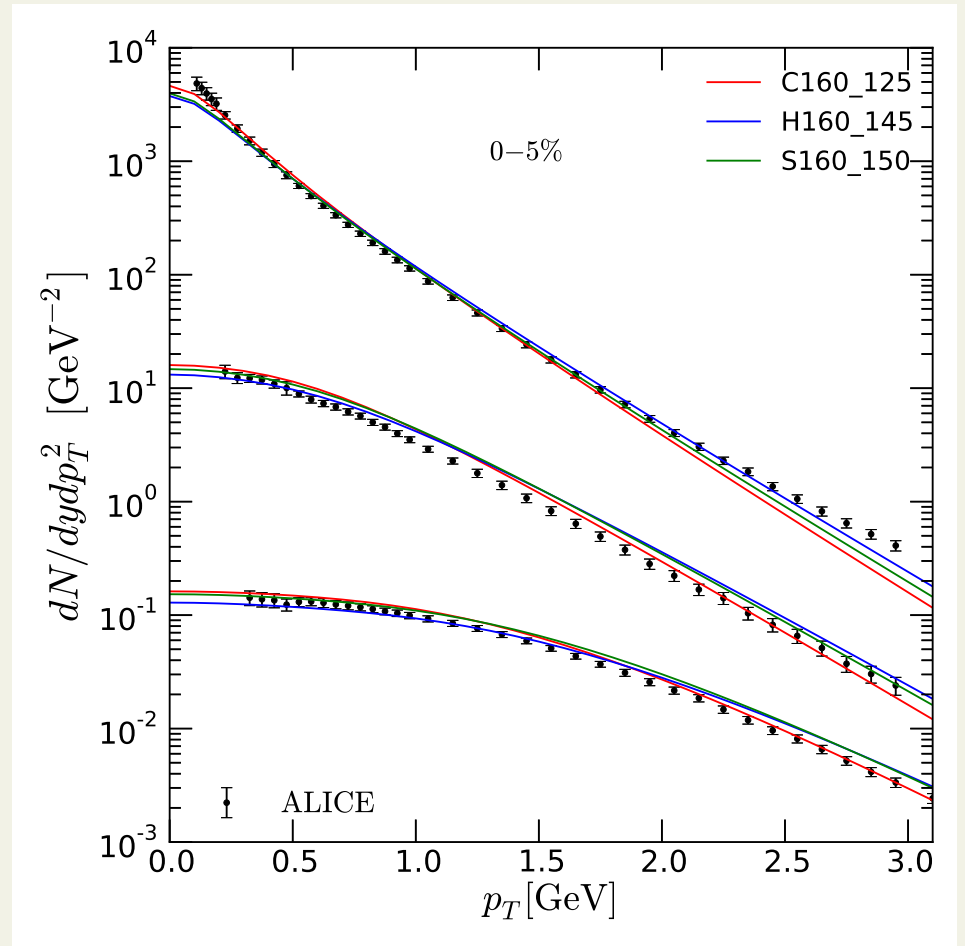
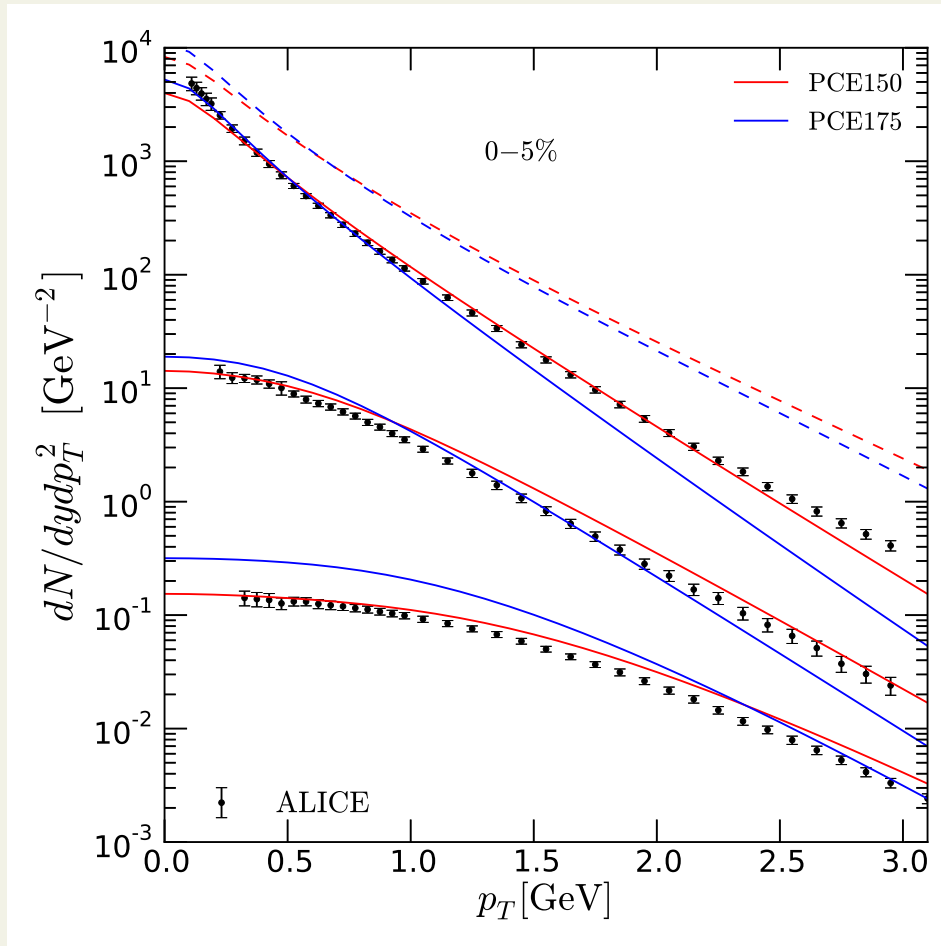
PCE175: $T_{\text{chem}} = 175 \text{ MeV}$

C160_125: $T_1 = 160 \text{ MeV}, T_2 = 125 \text{ MeV}$

H160_145: $T_{\text{hyperon}} = 160 \text{ MeV}, T_{\text{rest}} = 145 \text{ MeV}$

S160_150: $T_{\text{strange}} = 160 \text{ MeV}, T_{\text{rest}} = 150 \text{ MeV}$

Identified particle p_T spectrum at LHC



PCE150: $T_{\text{chem}} = 150 \text{ MeV}$

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Conclusions

- **thermal models are very simple**

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- thermal models do not give satisfactory fit to the ALICE data
- $N\bar{N}$ -annihilations lead to nice fit

Conclusions

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- thermal models do not give satisfactory fit to the ALICE data
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Conclusions

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- possible to construct an EoS to fit the observed yields