

Dissipative effects in mixtures

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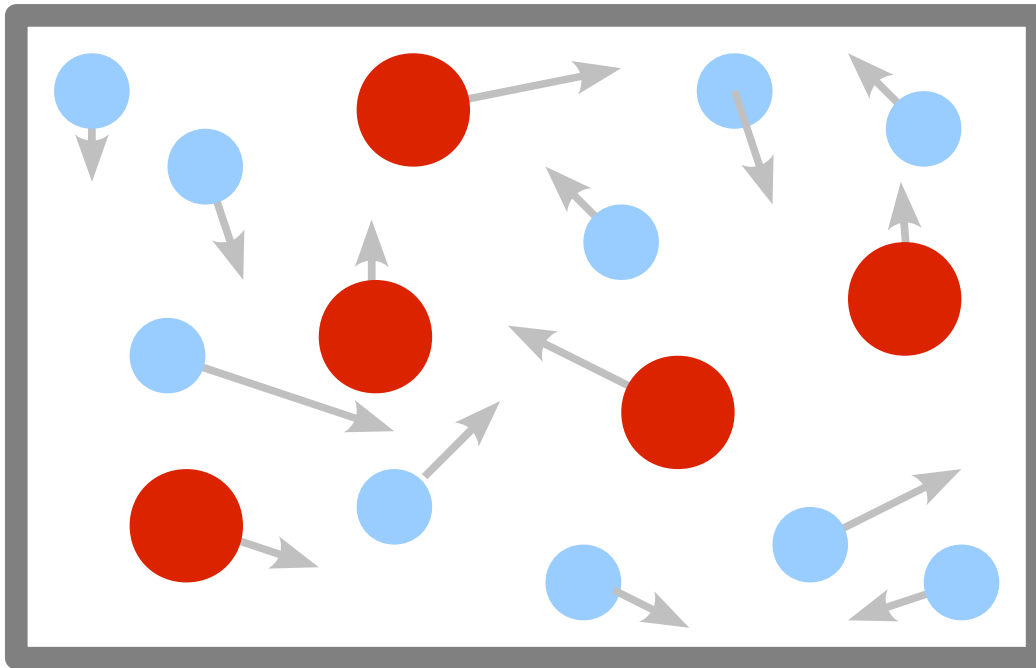
Motivation

Comparison of hydrodynamic calculations with experimental data
→ extraction of η/s , *EoS* ...

From η/s we learn about the inner dynamics in the medium

Motivation

Can we apply standard, one-component hydrodynamics to describe dissipative effects in a mixture?



Interaction cross sections

$$\sigma_{11}, \sigma_{12}, \sigma_{22}$$

*Two distinct
mean-free path scales*

QGP or Hadron Gas are mixtures

Dissipative hydro

Ideal hydrodynamics:

$$\partial_\mu T^{\mu\nu} = 0 \quad T^{\mu\nu} \text{ is isotropic locally} \leftrightarrow \text{Momentum distribution is isotropic}$$

$$\partial_\mu N^\mu = 0$$

EoS

Israel-Stewart hydrodynamics:

$$\partial_\mu T^{\mu\nu} = 0$$

$$\partial_\mu N^\mu = 0$$

$$\dot{\pi}^{\mu\nu} = -\frac{\pi^{\mu\nu}}{\tau_\pi} + \frac{\sigma^{\mu\nu}}{\beta_2} + \pi^{\mu\nu} \frac{T}{\beta_2} \partial_\alpha \left(\frac{\beta_2}{2T} u^\alpha \right)$$



Relaxation time $\tau_\pi = 2\beta_2\eta$

$\pi^{\mu\nu}$ describes anisotropy of the momentum space distribution

Dissipative hydro in a mixture

$\partial_\mu T^{\mu\nu} = 0$ *total energy is conserved*

$\partial_\mu N^\mu = 0$ *total particle number is conserved*

$$\dot{\pi}^{\mu\nu} = \sum \dot{\pi}_i^{\mu\nu}$$



*Same as the standard Israel-Stewart for each mixture component?
Then, what is the viscosity of **mixture components**?*

Dissipative hydro in a mixture

For our derivation we assume:

- particle numbers are conserved (no radiative processes)
- $T_1 = T_2 = T$ (rather strong assumption, but we really need this one)
- the global frame u^μ is not very different from the frames u_i^μ of the components



Need to check these two in transport calculations

Dissipative hydro in a mixture

For our derivation we assume:

- particle numbers N_i are conserved (no radiative processes)
 - $T_1 = T_2 = T$ (rather strong assumption, but we really need this one)
 - the global frame u^μ is not very different from the frames u^μ_i of the components
-
- we take isotropic scatterings, $d\sigma/d\Omega$ independent of Ω



can translate cross section to viscosity $\eta = \frac{6}{5} \frac{T}{\sigma}$

Dissipative hydro in a mixture

For two components with a given n_1/n_2 and cross sections σ_{11} , σ_{12} , σ_{22}

$$\dot{\pi}_1 = -\pi_1 \cdot \left(\frac{5}{9} n_1 \sigma_{11} + \frac{7}{9} n_2 \sigma_{12} \right) + \pi_2 \cdot \frac{2}{9} n_1 \sigma_{12} + \text{gradient terms}$$

$$\dot{\pi}_2 = -\pi_2 \cdot \left(\frac{5}{9} n_2 \sigma_{22} + \frac{7}{9} n_1 \sigma_{12} \right) + \pi_1 \cdot \frac{2}{9} n_2 \sigma_{12} + \text{gradient terms}$$

Two relaxation times per equation
Compare with the Israel-Stewart Eq.:

$$\dot{\pi} = -\pi \cdot \frac{5}{9} n \sigma + \text{gradient terms}$$

Dynamics in a mixture

Let's check the relaxation part of the equations

$$\dot{\pi}_1 = -\pi_1 \cdot \left(\frac{5}{9} n_1 \sigma_{11} + \frac{7}{9} n_2 \sigma_{12} \right) + \pi_2 \cdot \frac{2}{9} n_1 \sigma_{12} + \text{gradient terms}$$

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Dynamics in a mixture

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$$n_1, n_2$$

$$T_1 = T_2 = T$$

$$\sigma_{11}, \sigma_{12}, \sigma_{22}$$

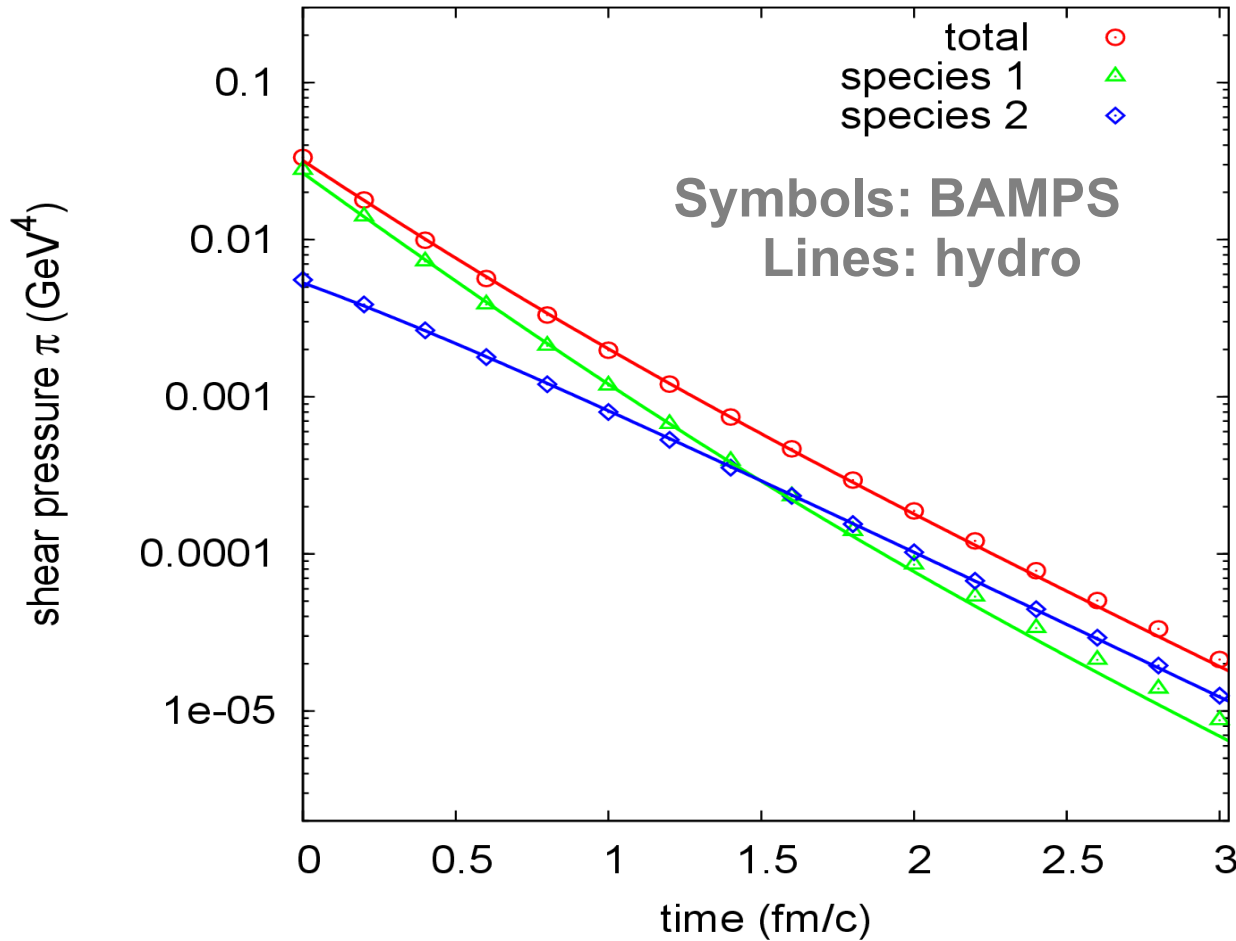
$$\pi_1, \pi_2$$

BAMPS BOX

$$f_i = d \cdot e^{-E_i/T} \cdot \left(1 + \frac{3}{8 e_i T^2} \cdot \pi_i \cdot \left(\frac{1}{2} p_T^2 - p_z^2 \right) \right)$$

Grad's Formalism

Mixture in BAMPS



$$n_1/n_2 = 5$$

$$\text{initial } \pi_1/\pi_2 = n_1/n_2$$

$$\sigma_{11} = 4 \text{ mb}$$

$$\sigma_{12} = 2 \text{ mb},$$

$$\sigma_{22} = 1 \text{ mb}$$

$$T = 0.4 \text{ GeV}$$

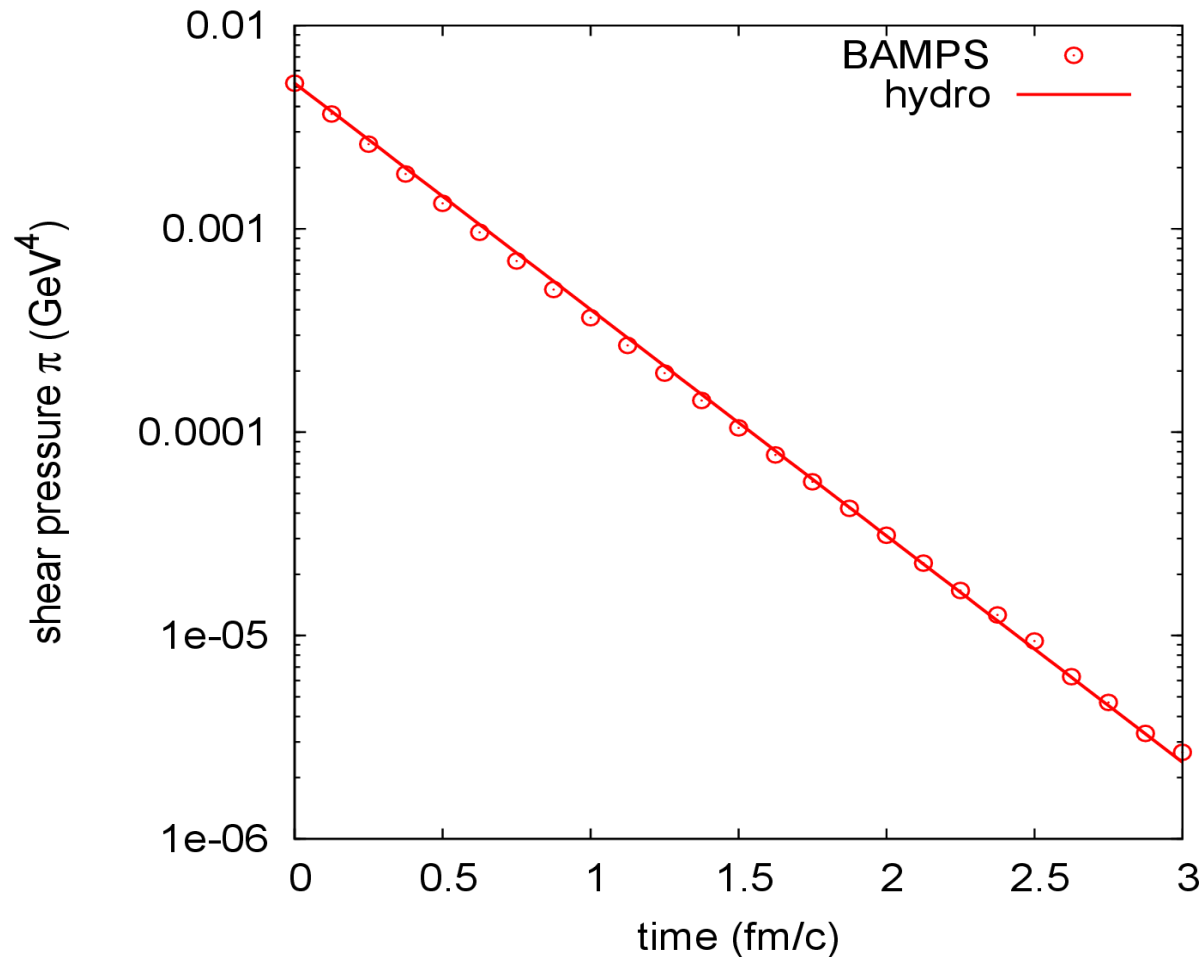
$$\dot{\pi}_1 = -\pi_1 \cdot \left(\frac{5}{9} n_1 \sigma_{11} + \frac{7}{9} n_2 \sigma_{12} \right) + \pi_2 \cdot \frac{2}{9} n_1 \sigma_{12}$$

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Relaxation in BAMPS

In the standard **one-component** Israel-Stewart hydrodynamics:

$$\dot{\pi} = -\frac{\pi}{\tau_{\pi}} = -\pi \cdot \frac{5}{9} n \sigma \longrightarrow \pi = \pi(0) \cdot e^{-\tau/\tau_{\pi}}$$



$$\sigma = 3.4 \text{ mb}$$
$$T = 0.4 \text{ GeV}$$

Green-Kubo

Application of Green-Kubo formula in BAMPS: *C.Wesp et al, arXiv:1106.4306*

$$\eta = \frac{V}{T} \int_0^{\infty} C(\tau) d\tau$$

Auto-correlation function

$$C(\tau) = \frac{1}{3} \left(\langle \pi^{xy}(0) \pi^{xy}(\tau) \rangle + \langle \pi^{xz}(0) \pi^{xz}(\tau) \rangle + \langle \pi^{yz}(0) \pi^{yz}(\tau) \rangle \right)$$

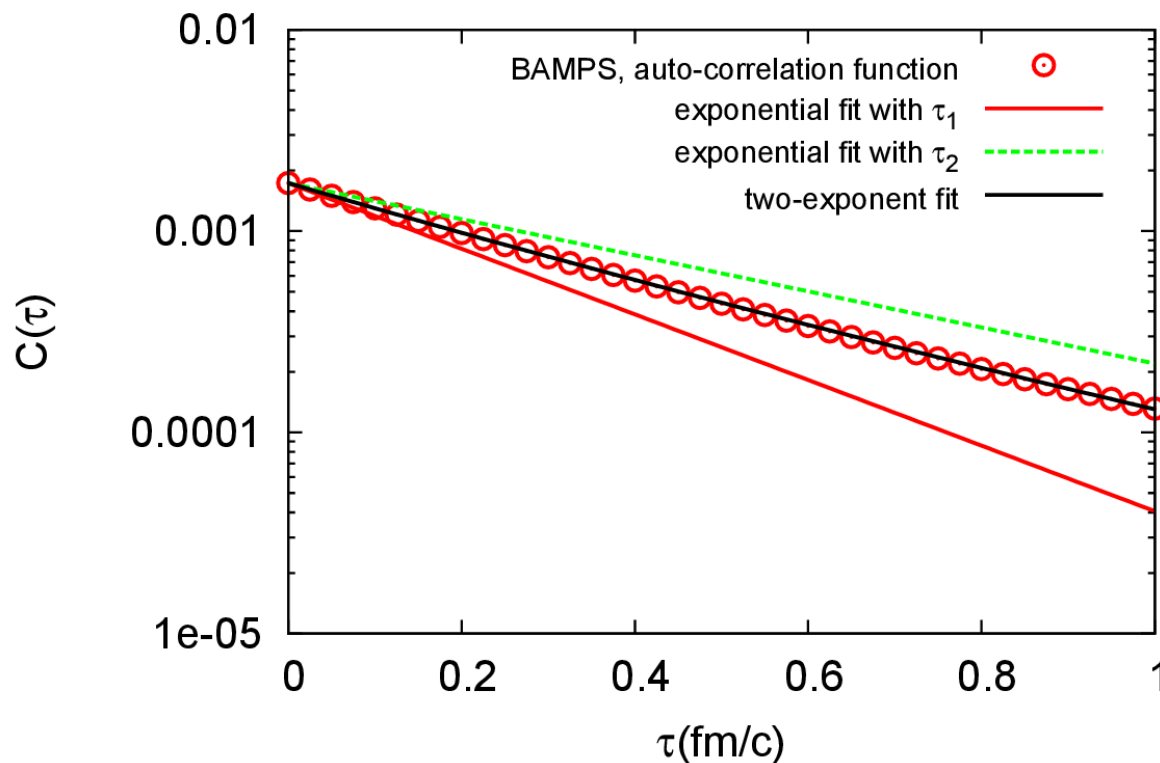
τ = correlation time

Green-Kubo

Application of Green-Kubo formula in BAMPS: *C.Wesp et al, arXiv:1106.4306*

$$\eta = \frac{V}{T} \int_0^\infty C(\tau) d\tau$$

Mixture in BAMPS



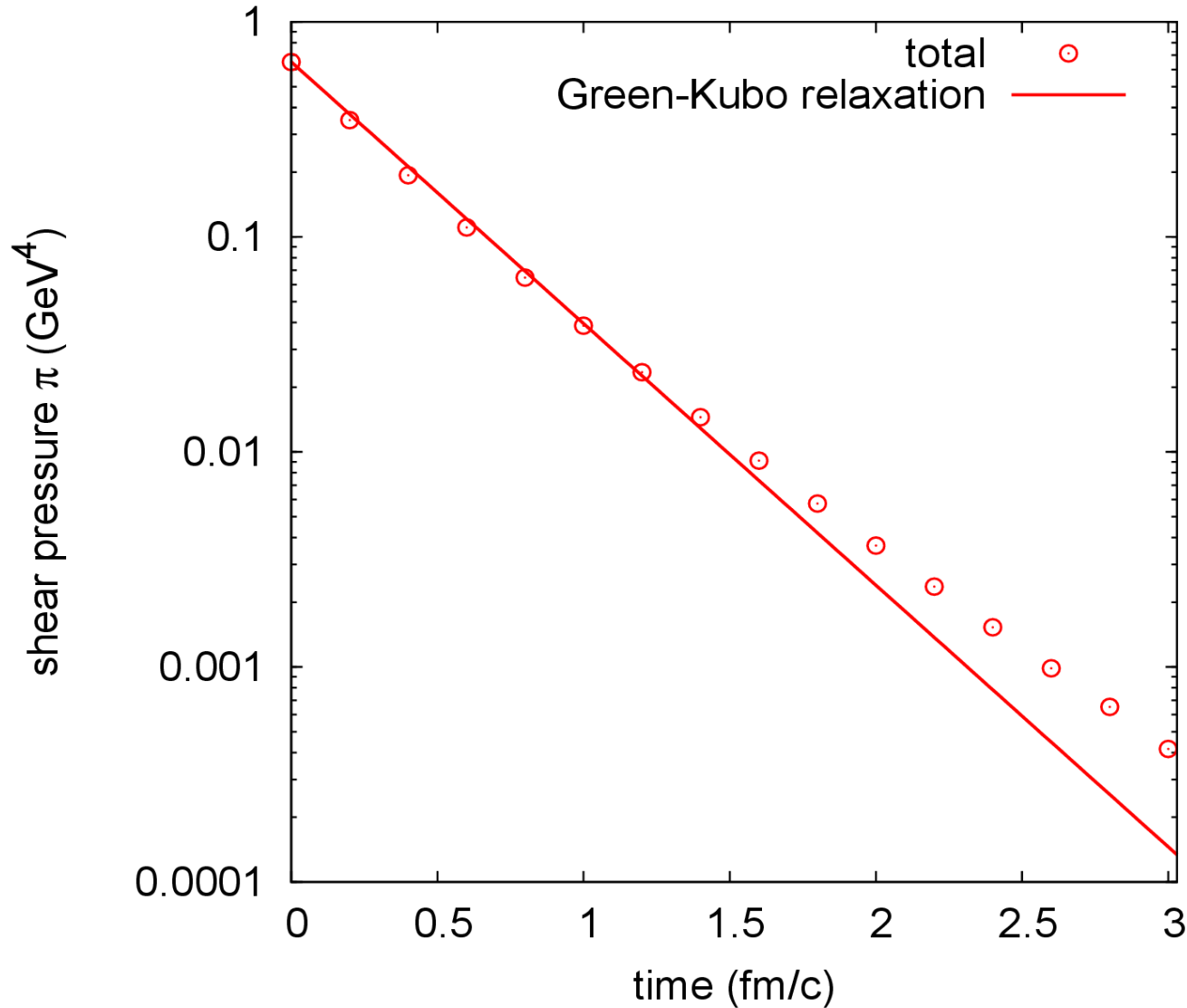
$$C(\tau) = C_1 \cdot e^{-\tau/\tau_1} + C_2 \cdot e^{-\tau/\tau_2}$$

$$\tau_1 \approx 0.2 \text{ fm/c}$$

$$\tau_2 \approx 0.4 \text{ fm/c}$$

$$\eta = 0.062 \text{ GeV}^3$$

Mixture in BAMPS



$$n_1 / n_2 = 5$$

$$\text{initial } \pi_1 / \pi_2 = n_1 / n_2$$

$$\sigma_{11} = 4 \text{ mb}$$

$$\sigma_{12} = 2 \text{ mb},$$

$$\sigma_{22} = 1 \text{ mb}$$

$$T = 0.4 \text{ GeV}$$

$$\text{Green-Kubo viscosity } \eta = 0.062 \text{ GeV}^3$$

Dynamics in a mixture

$$\dot{\pi}_1 = -\pi_1 \cdot \left(\frac{5}{9} n_1 \sigma_{11} + \frac{7}{9} n_2 \sigma_{12} \right) + \pi_2 \cdot \frac{2}{9} n_1 \sigma_{12}$$

$$\dot{\pi}_2 = -\pi_2 \cdot \left(\frac{5}{9} n_2 \sigma_{22} + \frac{7}{9} n_1 \sigma_{12} \right) + \pi_1 \cdot \frac{2}{9} n_2 \sigma_{12}$$

$$\pi(\tau) = A \cdot e^{-\tau/\tau_1} + B \cdot e^{-\tau/\tau_2}$$



$$\pi(\tau) = \pi_0 \cdot e^{-\tau/\tau_\pi}$$

$$\eta(\tau) = \frac{2}{5} \cdot e \cdot \left(\frac{r(\tau)}{1+r(\tau)} \cdot \lambda_1^{-1} + \frac{1}{1+r(\tau)} \cdot \lambda_2^{-1} \right)$$

with

$$r(\tau) = \frac{\pi_1(\tau)}{\pi_2(\tau)} = A(n, \sigma) \cdot \tanh(\tau \cdot B(n, \sigma) + C(n, \sigma, \pi_0))$$

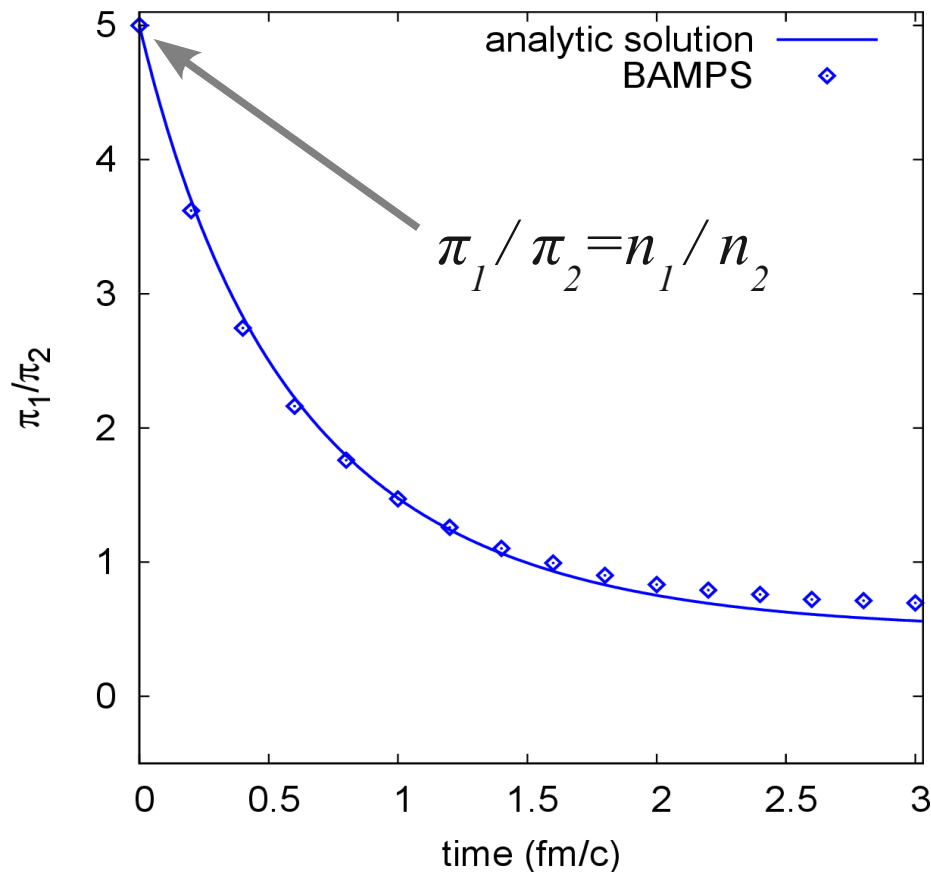
$$\lambda_1^{-1} = n_1 \sigma_{11} + n_2 \sigma_{12}$$

$$\lambda_2^{-1} = n_2 \sigma_{22} + n_1 \sigma_{12}$$

Mixture in BAMPS

$$\eta(\tau) = \frac{2}{5} \cdot e \cdot \left(\frac{r(\tau)}{1+r(\tau)} \cdot \lambda_1^{-1} + \frac{1}{1+r(\tau)} \cdot \lambda_2^{-1} \right)$$

$$r(\tau) = \frac{\pi_1(\tau)}{\pi_2(\tau)} = A(n, \sigma) \cdot \tanh(\tau \cdot B(n, \sigma) + C(n, \sigma, \pi_0))$$

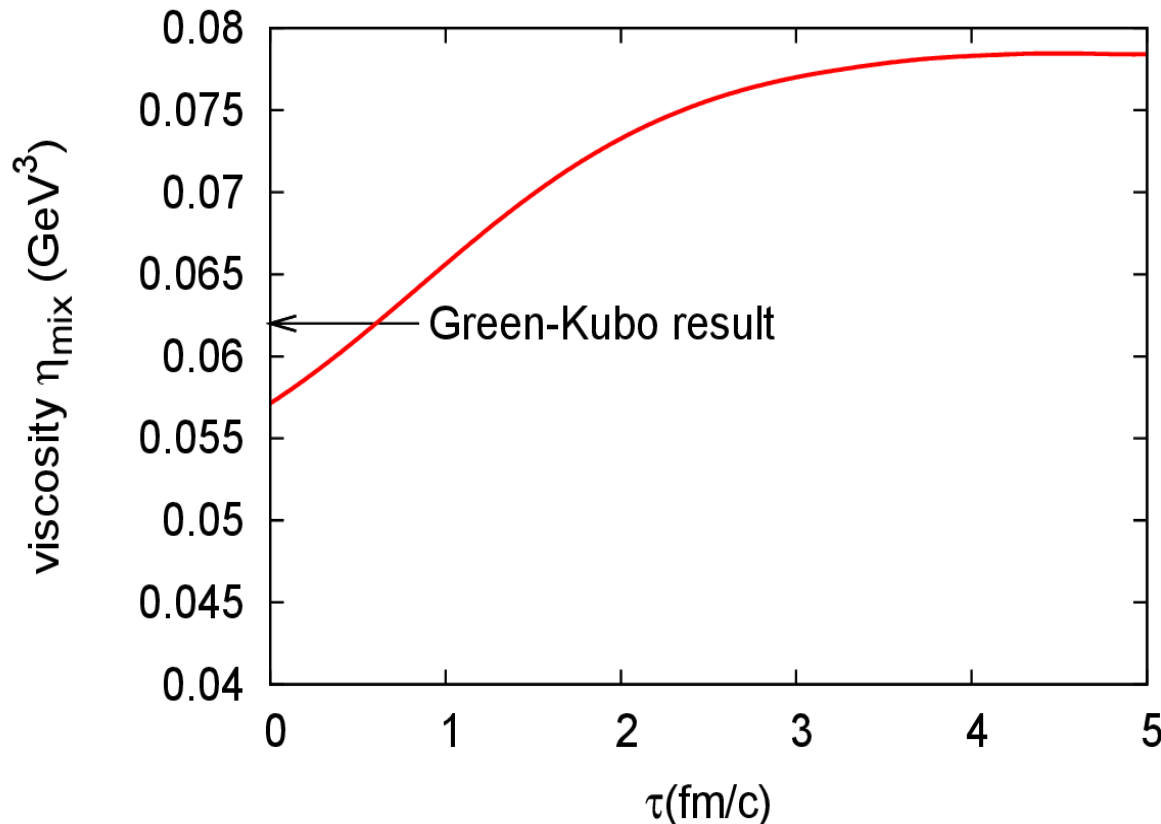


Existence of a characteristic stationary value

Dynamics in mixtures

$$\eta(\tau) = \frac{2}{5} \cdot e \cdot \left(\frac{r(\tau)}{1+r(\tau)} \cdot \lambda_1^{-1} + \frac{1}{1+r(\tau)} \cdot \lambda_2^{-1} \right)$$

$$r(\tau) = \frac{\pi_1(\tau)}{\pi_2(\tau)} = A(n, \sigma) \cdot \tanh(\tau \cdot B(n, \sigma) + C(n, \sigma, \pi_0))$$



mixture with mean free path scales $\lambda_1 \sim 0.2$ fm and $\lambda_2 \sim 0.4$ fm

Dynamics in mixtures

$$\eta(\tau) = \frac{2}{5} \cdot e \cdot \left(\frac{r(\tau)}{1+r(\tau)} \cdot \lambda_1^{-1} + \frac{1}{1+r(\tau)} \cdot \lambda_2^{-1} \right)$$

$$r(\tau) = \frac{\pi_1(\tau)}{\pi_2(\tau)} = A(n, \sigma) \cdot \tanh(\tau \cdot B(n, \sigma) + C(n, \sigma, \pi_0))$$

From the results so far we can conclude

- Existence of a characteristic time-dependence of the viscosity in a mixture
- Applicability of one-component hydrodynamics to a mixture depends on the chosen initial conditions

Initializing hydrodynamic calculations

$$\dot{\pi}^{\mu\nu} = -\frac{\pi^{\mu\nu}}{\tau_\pi} + \frac{\sigma^{\mu\nu}}{\beta_2} + \pi^{\mu\nu} \frac{T}{\beta_2} \partial_\alpha \left(\frac{\beta_2}{2T} u^\alpha \right)$$

In one-component hydrodynamic calculations the standard choices are

$$\pi^{\mu\nu}(\tau_0) = 0 \quad \vee \quad \pi^{\mu\nu}(\tau_0) = 2\eta\sigma^{\mu\nu}$$

...but there is no clear prescription what's the right choice.

For a mixture this would mean

$$\pi_i^{\mu\nu}(\tau_0) = 0 \quad \rightarrow \quad \dot{\pi}_i^{\mu\nu}(\tau_0) = \frac{\sigma^{\mu\nu}}{\beta_i} \quad \rightarrow \quad \frac{\pi_1(\tau_0 + d\tau)}{\pi_2(\tau_0 + d\tau)} = \frac{\beta_2}{\beta_1} = \frac{e_1}{e_2} = \frac{n_1}{n_2}$$

Which also means, that the characteristic time-dependence of shear viscosity must be taken into account

Conclusions and Outlook

- standard one-component hydrodynamics in general cannot be applied to describe dissipative effects in mixtures
- It is only in case the initial conditions are chosen properly that one-component description can be applied
- Green-Kubo formalism is not reliable for mixtures – additional time-modulation must be taken into account

$$\eta/s = \eta/s(T) \rightarrow \eta/s = \eta/s(T) * f(t)$$

Conclusions and Outlook

- Most reliable way to check these conclusions:
Kinetic transport calculations → BAMPS

See how evolution of an expanding “QGP” with $\sigma_{gg}, \sigma_{gq}, \sigma_{qq} \sim 1/T^2$
can be reproduced by one-component calculations.

How the cross section (i.e. η/s) must be chosen in one-component case?
Can any hint of the time-dependence of η be seen?

Mixture in BAMPS → (isochronous) freeze-out → flow observables
vs

One-component fluid in BAMPS → (isochronous) freeze-out → flow observables

Work in progress