

Frankfurt, 13.06.2022

Höhere Quantenmechanik
Summer term 2022

Exercise sheet 9
(Submission date: Until 20.06.2022 12:00)

Evaluation der Lehrveranstaltung am 21.06.22
https://itp.uni-frankfurt.de/~gros/Vorlesungen/QM_2/QM2_evaluation_22.pdf

Exercise 1: spin operators for fermions (12 Points)

We consider a fermionic particle with spin $S = 1/2$ (e.g., an electron), which can be described by the annihilation and creation operators c_α and c_α^\dagger . The index α can take the values $\alpha = 1 = \uparrow$ and $\alpha = 2 = \downarrow$, which indicate, respectively, whether the spin of the particle points up or down (with respect to a certain arbitrary S_z axis). The fermionic operators satisfy the anticommutation relation

$$[c_\alpha, c_\beta]_+ = 0 \quad [c_\alpha^\dagger, c_\beta^\dagger]_+ = 0 \quad [c_\alpha, c_\beta^\dagger]_+ = \delta_{\alpha,\beta} \quad (1)$$

We can define the spin operators for the fermionic particle as

$$S_a = \frac{\hbar}{2} \sum_{\alpha,\beta} c_\alpha^\dagger \sigma_{\alpha,\beta}^a c_\beta \quad (2)$$

where $a = x, y, z$ indicates the spin component and σ^a is the corresponding Pauli matrix. Within this notation, Greek indices run over 1, 2 (or, equivalently, \uparrow, \downarrow), while Latin indices label the x, y, z components of the spin. Note: superscripts/subscripts don't have any special meaning!

- (i) Show that the definition of Eq. (2) satisfies the usual commutation relations of the spin operators, i.e.

$$[S_x, S_y]_- = i\hbar S_z \quad [S_y, S_z]_- = i\hbar S_x \quad [S_z, S_x]_- = i\hbar S_y$$

(7 Points)

Hint 1: you can prove the three equations one by one, or, alternatively, you can pursue a general proof by starting from the expression

$$[S_a, S_b]_- = i\hbar \sum_{c=1}^3 \varepsilon_{a,b,c} S_c,$$

where a, b, c are the component indices (1 for x , 2 for y , 3 for z), and $\varepsilon_{a,b,c}$ is the Levi-Civita symbol.

Hint 2: for three generic operators, A, B, C , the following properties of commutators and anticommutators hold

$$\begin{aligned} [A, BC]_- &= [A, B]_- C + B[A, C]_- \\ [AB, C]_- &= A[B, C]_- + [A, C]_- B \\ [AB, C]_- &= A[B, C]_+ - [A, C]_+ B \end{aligned}$$

For Pauli matrices, the following holds

$$\sigma^a \sigma^b = \delta_{a,b} \mathbb{1} + i \sum_c \varepsilon_{a,b,c} \sigma^c$$

where a, b, c are the component indices (1 for x , 2 for y , 3 for z), and $\mathbb{1}$ the identity matrix.

- (ii) Apply the S_z operator to the state $|\uparrow\rangle = c_\uparrow^\dagger|0\rangle$ or $|\downarrow\rangle = c_\downarrow^\dagger|0\rangle$. What do you get? Comment the result. (2 Points)
- (iii) How do the S_+ and S_- operators look like? Discuss the result. (2 Points)
- (iv) Suppose that the spin- $\frac{1}{2}$ fermionic operators of this exercise create/destroy electrons in a certain s -orbital. The state $|\uparrow\downarrow\rangle = c_\uparrow^\dagger c_\downarrow^\dagger|0\rangle$ corresponds to a filled orbital (no other electrons can be added). What is the result of $S^2|\uparrow\downarrow\rangle$ (with $S^2 = S_x^2 + S_y^2 + S_z^2$)? Don't do any calculation, just use physical arguments. (1 Point)

Exercise 2: one-body operators in second quantization, fields (8 Points)

As discussed in the lecture notes, in second quantization we define annihilation and creation operators by choosing a certain basis of single-particle states. Let us consider identical bosonic particles in this exercise. If we choose the complete orthonormal basis $\{|\alpha\rangle\}$, we can write the annihilation and creation operators as \hat{b}_α and \hat{b}_α^\dagger , which respectively destroy and create a particle in the single-particle state $|\alpha\rangle$. These operators satisfy the commutation relations discussed in the notes. Many-particle states are then obtained by successive application of the creation operator on the vacuum state. Thanks to the commutation properties, the many-particle states satisfy the correct properties for identical particles, i.e., in the case of bosons they are symmetric under the exchange of two particles.

Obviously, the choice of the single-particle basis is not unique. We could choose another basis, $\{|n\rangle\}$, and write the transformation between the annihilation/creation operators in the two bases as follows

$$\begin{aligned} \hat{b}_n &= \sum_\alpha \langle n|\alpha\rangle \hat{b}_\alpha \\ \hat{b}_n^\dagger &= \sum_\alpha (\langle n|\alpha\rangle)^* \hat{b}_\alpha^\dagger = \sum_\alpha \langle \alpha|n\rangle \hat{b}_\alpha^\dagger \end{aligned}$$

Note: in this exercise we use a “hat” to denote operators in order to avoid confusion.

- (i) Consider the special case of the eigenbasis of the position operator \hat{x} , which we denote by $\{|x\rangle\}$. For simplicity, we imagine to be in one dimension. The annihilation and creation operators in the position basis are usually called *field operators* and written as

$$\begin{aligned} \hat{\psi}(x) &= \sum_\alpha \langle x|\alpha\rangle \hat{b}_\alpha = \sum_\alpha \varphi_\alpha(x) \hat{b}_\alpha \\ \hat{\psi}^\dagger(x) &= \sum_\alpha \langle \alpha|x\rangle \hat{b}_\alpha^\dagger = \sum_\alpha \varphi_\alpha^*(x) \hat{b}_\alpha^\dagger \end{aligned}$$

They destroy or create a particle in the position x . Their expansion in terms of \hat{b}_α and \hat{b}_α^\dagger involves the wave functions $\varphi_\alpha(x)$ and $\varphi_\alpha^*(x)$.

Show that the following commutation relations hold

$$[\hat{\psi}(x), \hat{\psi}(x')]_- = 0 \quad [\hat{\psi}^\dagger(x), \hat{\psi}^\dagger(x')]_- = 0 \quad [\hat{\psi}(x), \hat{\psi}^\dagger(x')]_- = \delta(x - x')$$

(2 Points)

- (ii) Single-particle operators can be written in the form of Eq.(6.8) of the lecture notes, i.e. $\hat{O} = \sum_{i=1}^N \hat{o}(i)$, where $\hat{o}(i)$ is an operator that acts only on the i -th particle (N is the total number of particles in the system). The argument i just tells us on which particle the operator \hat{o} acts.

In second quantization, using annihilation/creation operators of a generic basis $\{|\alpha\rangle\}$, single-particle operators \hat{O} are expressed as

$$\hat{O} = \sum_{\alpha, \alpha'} \langle \alpha | \hat{o} | \alpha' \rangle \hat{b}_\alpha^\dagger \hat{b}_{\alpha'} \quad (3)$$

Using this formula, express the position operator (i.e., take $\hat{o} = \hat{x}$) in terms of the annihilation/creation operators of the basis $\{|x\rangle\}$ (i.e., the field operators $\hat{\psi}(x)$ and $\hat{\psi}^\dagger(x)$).

(2 Points)

- (iii) Express the density operator $\hat{N}_r = \sum_{i=1}^N \hat{n}_r(i)$ in terms of the annihilation/creation operators of the basis $\{|x\rangle\}$ (i.e., the field operators $\hat{\psi}(x)$ and $\hat{\psi}^\dagger(x)$). You should use Eq. (3) with

$$\hat{o} = \hat{n}_r = \delta(\hat{x} - r), \quad (4)$$

where r is a certain position in space. (2 Points)

- (iv) We now introduce the basis of plane waves $\{|k\rangle\}$, where $\langle x|k\rangle = \frac{1}{\sqrt{V}} e^{ikx}$ (V being the volume of the system). Express the density operator in terms of creation and annihilation operators of the $\{|k\rangle\}$ basis. Use Eq. (3) with $\hat{o} = \hat{n}_r = \delta(\hat{x} - r)$. Compare the result to the one of the previous point. (2 Points)