

# **Thermodynamics at strong coupling from Holographic QCD**

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**APC, U. Paris VII**

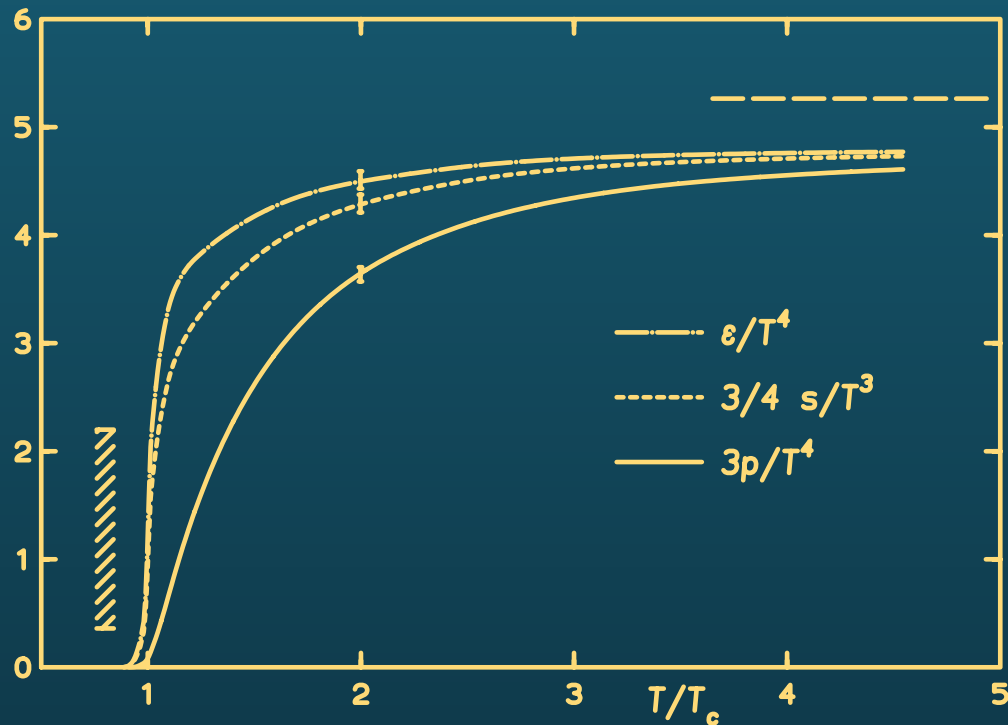
**Excited QCD**

**Les Houches, February 23 2011**

**Work with E. Kiritsis, U. Gursoy, L. Mazzanti, G. Michalogiorgakis, '07-'10**

# Motivation: Thermodynamics of 4D YM

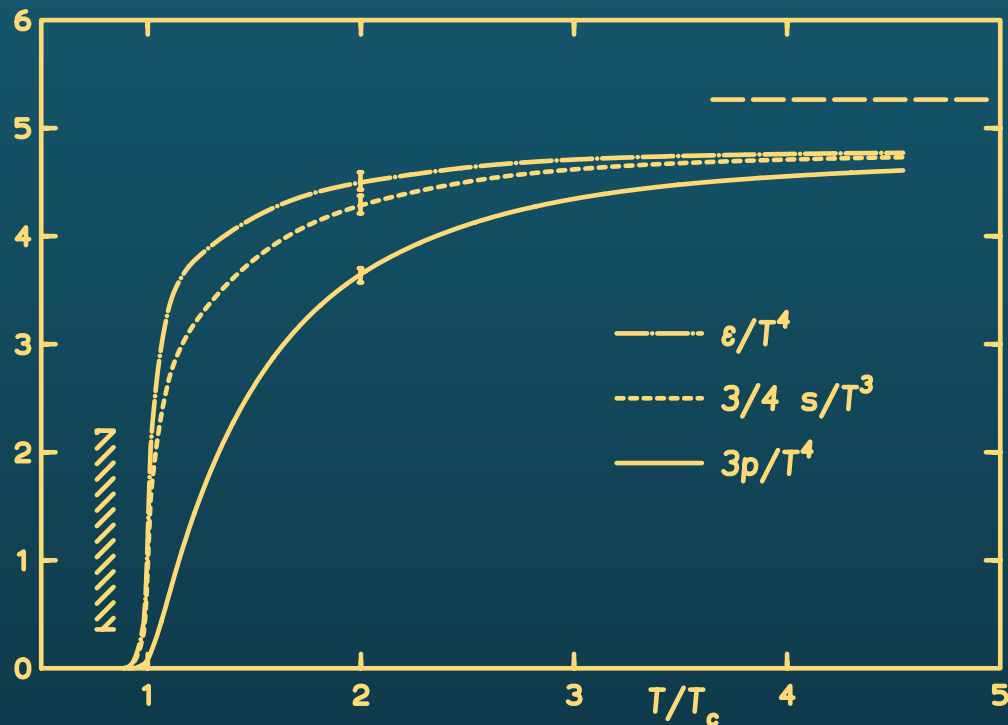
4D Pure YM theory exhibits a **first order deconfining phase transition** around  $T_c \sim 260 MeV$ .



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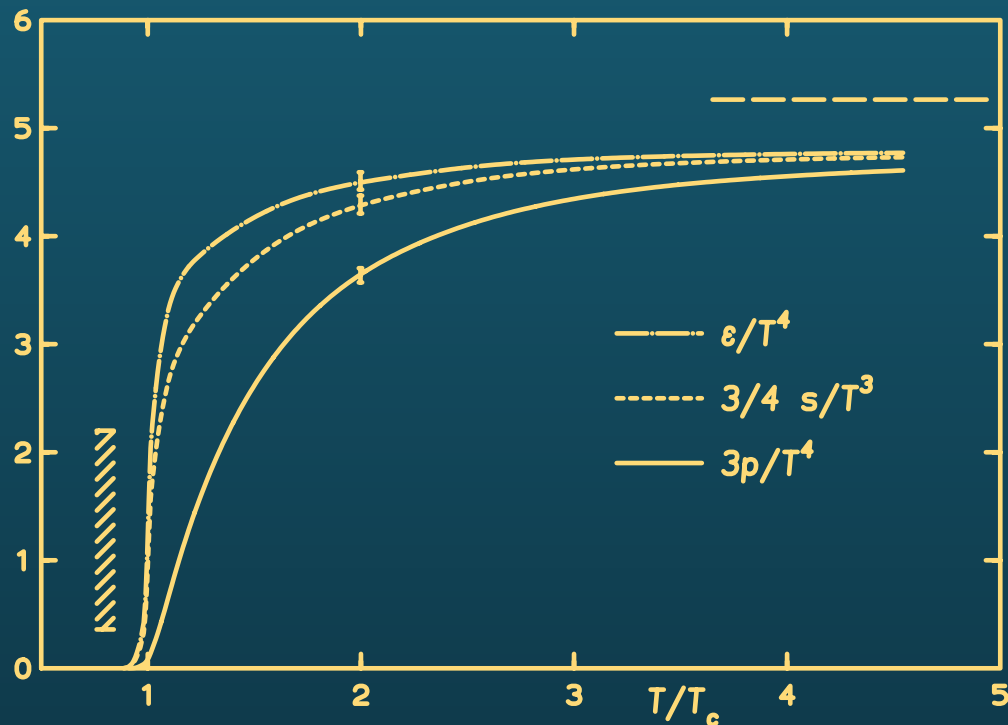


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- Can AdS/CFT techniques reproduce the phase transition and the YM equation of state found on the lattice?
- If yes, can it compute hydrodynamic and transport coefficient of the QCD plasma?

# Outline

- AdS/CFT 101
- Holographic description of deconfinement transition
  - AdS examples
  - 5D Phenomenological models

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( = being creative)

# Gauge/Gravity duality

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I will restrict the discussion to the **5D**  $(x^\mu, r)$  directions, and consider general **non-conformal backgrounds**.

# Field/Operator correspondence

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Gravity description valid at large  $N$ , large  $g^2 N$

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The **partition function** is given by the **gravity action evaluated at its extremum** (solution of the of the gravity-side field equation)

$$\mathcal{Z}(\beta) = e^{-S_{grav}[g_0, \Phi_0]}$$

where  $g_0$  and  $\Phi_0$  are periodic with period  $\beta$ .

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- From  $\mathcal{F}_i(T)$  we can compute other equilibrium quantities (entropy, pressure, speed of sound) using standard thermodynamic formulae.

# Plasma/Black Hole correspondence

Holographic description of a deconfined plasma in thermal equilibrium at temperature  $T \equiv 1/\beta$ :



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- Due to AdS boundary conditions, **BH is stable**.
- Polyakov loop gets a vev, signaling deconfinement.
- Spectrum of quasinormal modes.

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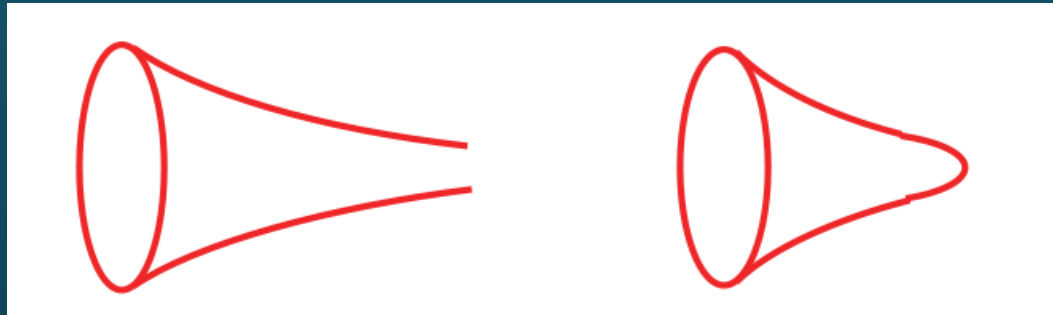
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Thermal gas in  $AdS$

$AdS$  Black Hole



$$ds_{TG}^2 = \frac{\ell}{r^2} [dr^2 + d\tau^2 + d\vec{x}^2] \quad ds_{BH}^2 = \frac{\ell}{r^2} \left[ \frac{dr^2}{f(r)} + f(r)d\tau^2 + d\vec{x}^2 \right]$$

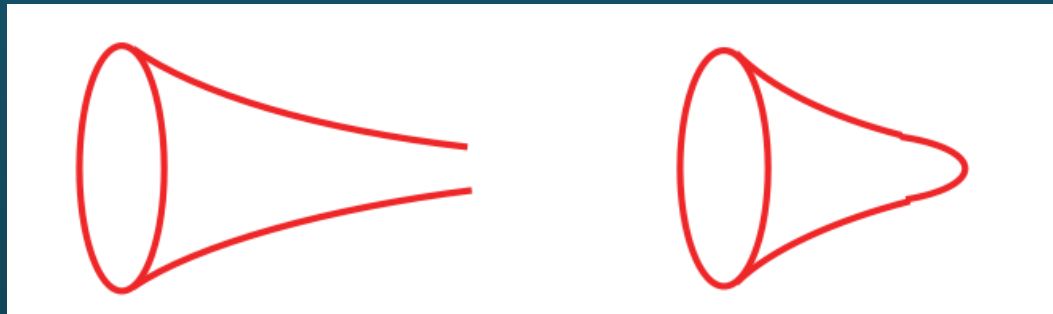
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$\mathcal{F}_{AdS} > \mathcal{F}_{BH}$  for all  $T > 0$ : BH always dominates;  
Thermodynamics of a conformal gas,  $\mathcal{F} = c T^4$

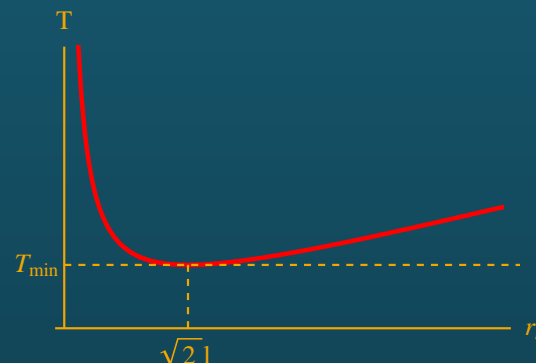
$$\mathcal{N} = 4 \text{ SYM on } S^3 \times S^1 \Leftrightarrow \text{Global AdS}_5$$

To get a phase transition at nonzero  $T$  we need to introduce a **scale**,  
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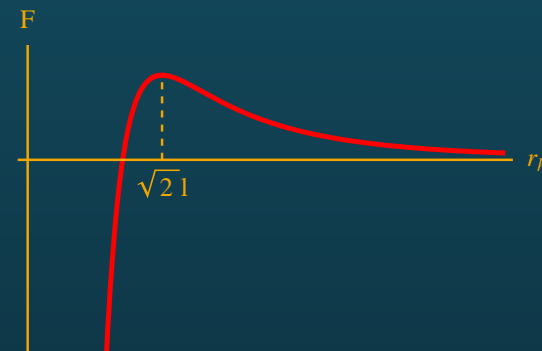
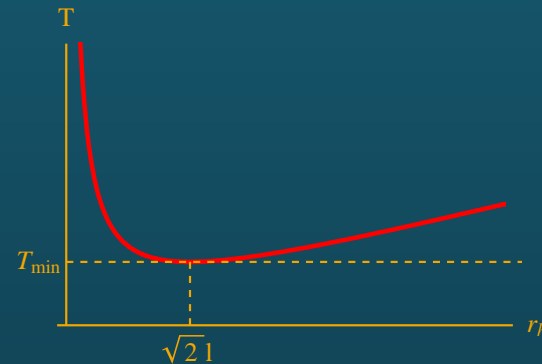




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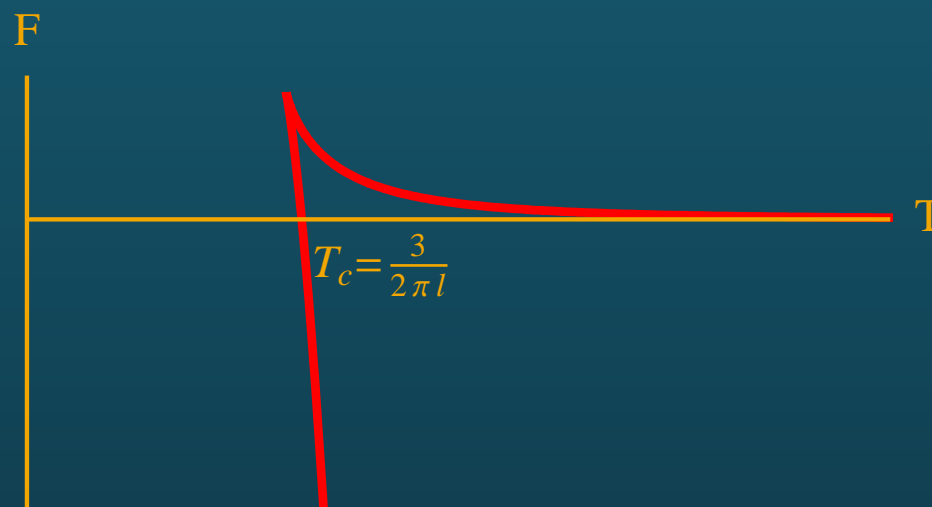
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- “Large” BH dominant for  
 $T_c > T_{min}$



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Hawking-Page transition.

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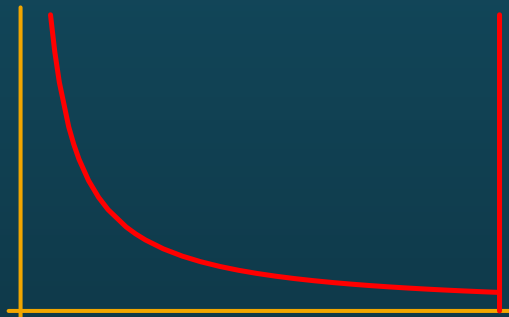
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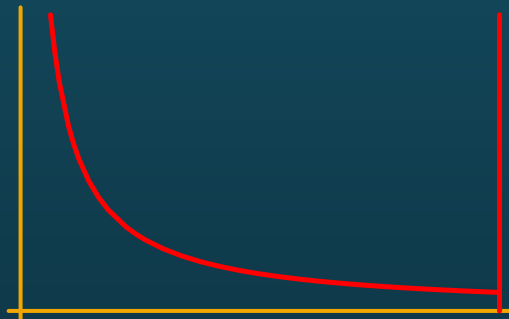


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- still “too” conformal, except for the transition;
- issue of boundary conditions at the IR cutoff

# 5D Einstein-Dilaton Theories

A simple class of models that displays **realistic** thermodynamics

$$S_E = -M_p^3 N_c^2 \int d^5x \sqrt{-g} \left[ R - \frac{4}{3} (\partial\Phi)^2 - V(\Phi) \right]$$

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- Black holes (**Deconfined phase - Gluon Plasma**):

$$ds^2 = b^2(r) \left( \frac{dr^2}{f(r)} + f(r) d\tau^2 + dx_i^2 \right), \quad \Phi = \Phi(r), \quad 0 < r < r_h$$

$$b(r) \xrightarrow{r \rightarrow 0} b_0(r) \sim \ell/r, \quad f(r) \xrightarrow{r \rightarrow 0} 1, \quad f(r_h) = 0,$$

# Dilatonic Black holes

Gürsoy, Kiritsis, Mazzanti, F.N. '08

Consider a potential that in the IR (large  $\Phi$ ) asymptotes:

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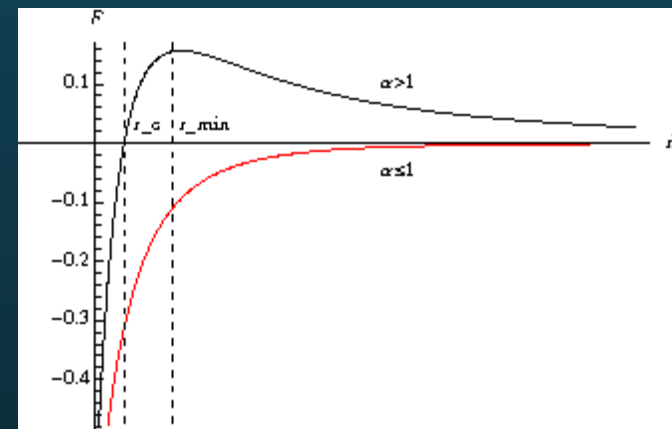
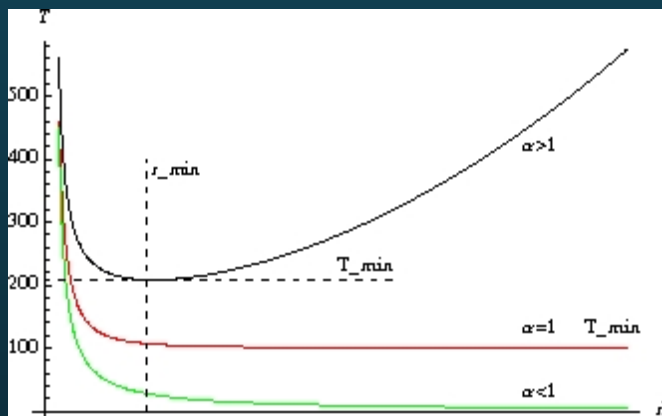
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Situation as in the global  $AdS$  case ( $\alpha > 1$ ) or Poincaré  $AdS$  ( $\alpha < 1$ )

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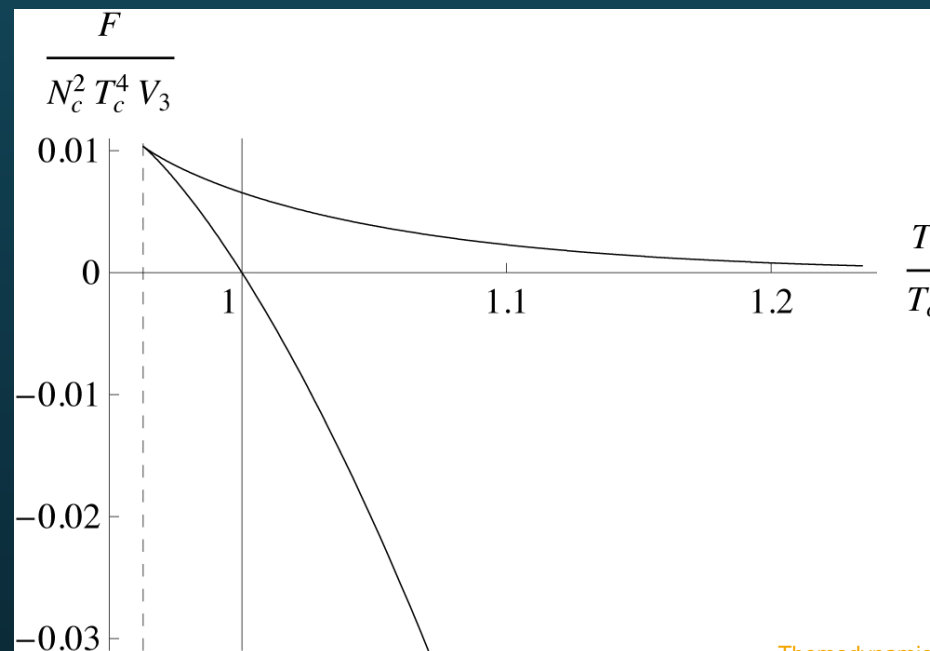
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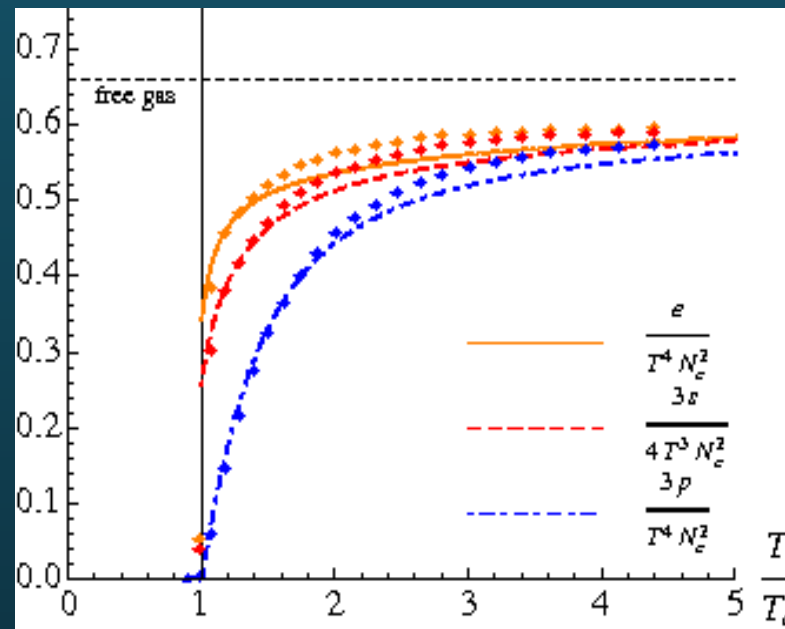
- **confining case: Hawking-Page transition**



# Explicit Model

Gürsoy, Kiritsis, Mazzanti, F.N. '09

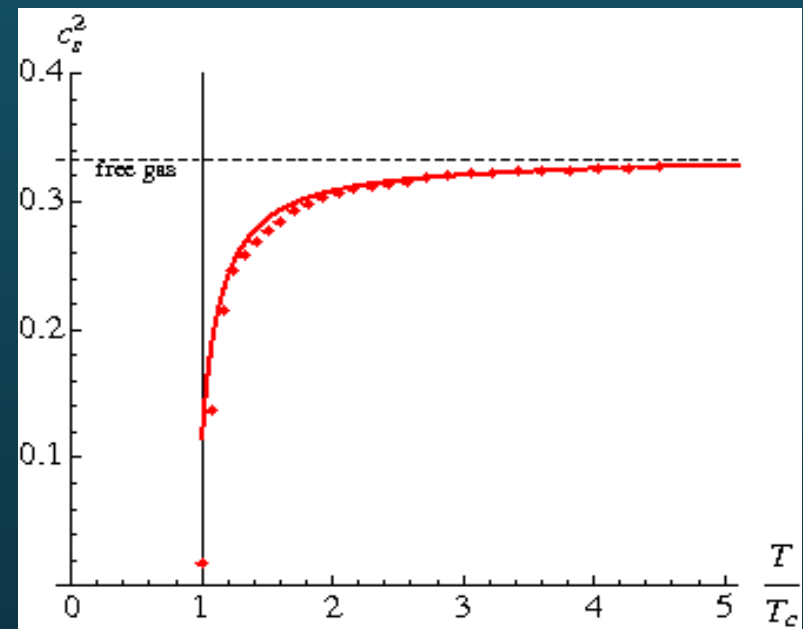
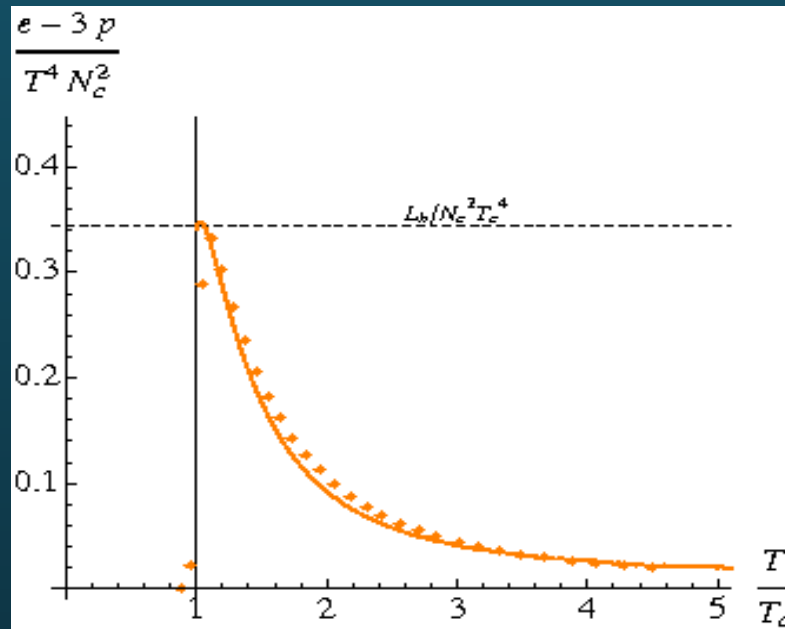
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- The potential  $V(\Phi)$  is supposed to be fixed phenomenologically as a way of parametrizing some of the unknown features of the full holographic dual.

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- Phenomenological 5D models can be used obtain quantitatively accurate results
- Natural to extend the description beyond equilibrium (e.g. diffusion)