Studies of Dilepton Production with UrQMD

Stephan Endres

Frankfurt Institute for Advanced Studies
ITP Uni Frankfurt

May 24th, 2012

H-QM Helmholtz Research School Quark Matter Studies

HGS-HIRe for FAIR

Overview

Overview

1. Introduction
2. Dilepton Production in UrQMD
3. Results
   - Results for HADES Energies
   - Rho Contribution
   - Cross-sections
4. Further Plans
   - Looking at NA60, RHIC and LHC
   - Next Steps

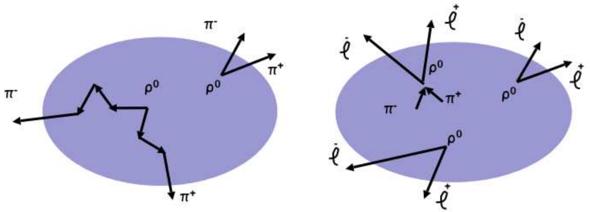
Overview

1. Introduction
2. Dilepton Production in UrQMD
3. Results
   - Results for HADES Energies
   - Rho Contribution
   - Cross-sections
4. Further Plans
   - Looking at NA60, RHIC and LHC
   - Next Steps

Overview

Why Dileptons...?

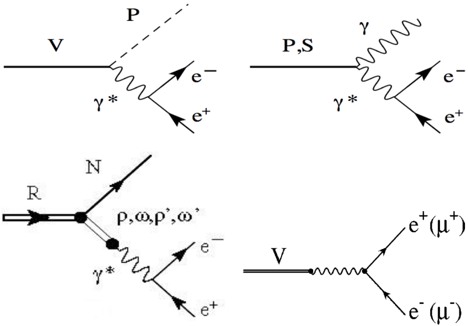
- Dileptons represent a a clean and penetrating probe of hot and dense nuclear matter
- Once produced they do not interact with the surrounding matter
- Aim of studies
  $\Rightarrow$ In-medium modification of vector meson properties
  $\Rightarrow$ Chiral symmetry restoration



Dilepton sources in UrQMD

- Dalitz Decays
  $\Rightarrow \pi^0, \eta, \eta', \omega, \Delta$
  $P \rightarrow \gamma + e^+e^-$
  $V \rightarrow P + e^+e^-$

- Direct Decays
  $\Rightarrow \rho^0, \omega, \phi$



Resonances and Branching Ratios in UrQMD

- Two processes possible in UrQMD

  Collisions

  (e.g. $\pi\pi \to \rho$)

  Resonance decays

  (e.g. $\mathrm{N}^* \to N + \rho$)
- At SIS energies, the resonance excitation and decay is dominant
- Branching ratios are in accordance with PDG

Resonance	Mass	Width	$N\pi$	$N\eta$	$N\omega$	$N\varrho$
N^*_{1440}	1.440	350	0.65			
N^*_{1520}	1.515	120	0.60			0.15
N^*_{1535}	1.550	140	0.60	0.30		
N^*_{1650}	1.645	160	0.60	0.06		0.06
N^*_{1675}	1.675	140	0.40			
N^*_{1680}	1.680	140	0.60			0.10
N^*_{1700}	1.730	150	0.05			0.20
N^*_{1710}	1.710	500	0.16	0.15		0.05
N^*_{1720}	1.720	550	0.10			0.73
N^*_{1900}	1.850	350	0.30	0.14	0.39	0.15
N^*_{1990}	1.950	500	0.12			0.43
N^*_{2080}	2.000	550	0.42	0.04	0.15	0.12
N^*_{2190}	2.150	470	0.29			0.24
N^*_{2220}	2.220	550	0.29		0.05	0.22
N^*_{2250}	2.250	470	0.18			0.25
Δ_{1232}	1.232	115	1.00			
Δ^*_{1600}	1.700	350	0.10			
Δ^*_{1620}	1.675	160	0.15			0.05
Δ^*_{1700}	1.750	350	0.20			0.25
Δ^*_{1900}	1.840	260	0.25			0.25
Δ^*_{1905}	1.880	350	0.18			0.80
Δ^*_{1910}	1.900	250	0.30			0.10
Δ^*_{1920}	1.920	200	0.27			
Δ^*_{1930}	1.970	350	0.15			0.22
Δ^*_{1950}	1.990	350	0.38			0.08

Dalitz Decays

- Dalitz decays can be decomposed into the corresponding decays into a virtual photon and the subsequent decay of the photon via electromagnetic conversion

$$\frac{d\Gamma}{dM^2} = \Gamma_{P\to\gamma\gamma*, V\to P\gamma*} \frac{1}{\pi M^4} M\Gamma_{\gamma*\to ee}$$

- Internal conversion probability of the photon

$$M\Gamma_{\gamma*\to ee} = \frac{\alpha}{3} M^2 \sqrt{1 - \frac{4m_e}{M^2}} \left(1 + \frac{2m_e^2}{M^2}\right)$$

Dalitz Decays

- The widths $\Gamma_{P\to\gamma\gamma^*}$ and $\Gamma_{V\to P\gamma^*}$ can be related to corresponding radiative widths

$$\Gamma_{\mathrm{P}\to\gamma\gamma^*} = 2\Gamma_{\mathrm{P}\to 2\gamma}\left(1-\frac{\mathrm{M}^2}{\mathrm{m}_\mathrm{p}^2}\right)^3 \left|\mathrm{F}_{\mathrm{P}\gamma\gamma^*}(\mathrm{M}^2)\right|^2$$

$$\Gamma_{\mathrm{V}\to\mathrm{P}\gamma^*} = 2\Gamma_{\mathrm{V}\to\mathrm{P}\gamma}\left[\left(1+\frac{\mathrm{M}^2}{\mathrm{m}_\mathrm{V}^2-\mathrm{m}_\mathrm{p}^2}\right)^2-\left(\frac{2\mathrm{m}_\mathrm{V}\mathrm{M}}{\mathrm{m}_\mathrm{V}^2-\mathrm{m}_\mathrm{p}^2}\right)^2\right]^{3/2} \cdot\left|\mathrm{F}_{\mathrm{VP}\gamma^*}(\mathrm{M}^2)\right|^2$$

Form Factors & Direct Decays

- Form factors for the Dalitz decays are obtained from the vector-meson dominance model (VMD). We use the parametrizations by Landsberg and Li/Ko/Brown/Sorge

- The width for the direct decay of a vector meson V to a dilepton pair varies with the dilepton mass like M^{-3}

$$\Gamma_{V\to ee}(M) = \frac{\Gamma_{V\to ee}(m_V)}{m_V}\frac{m_V^4}{M^3}\sqrt{1-\frac{4m_e}{M^2}}(1+\frac{2m_e^2}{M^2})$$

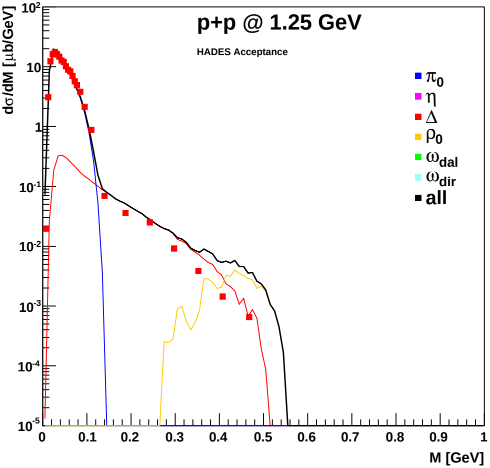
Time Integration Method (Shining)

- Shining in UrQMD applied for Δ, ρ, ω, ϕ and η'
- *Assumption*: Resonance can continuously emit dileptons over its whole lifetime
- Integration of the dilepton emission rate over time
- Collisional broadening of each individual parent resonance is taken into account

$$\frac{dN_{ee}}{dM} = \frac{\Delta N_{ee}}{\Delta M} = \sum_{j=1}^{N_{\Delta M}} \int_{t_i^j}^{t_f^j} \frac{dt}{\gamma} \frac{\Gamma_{ee}(M)}{\Delta M}$$

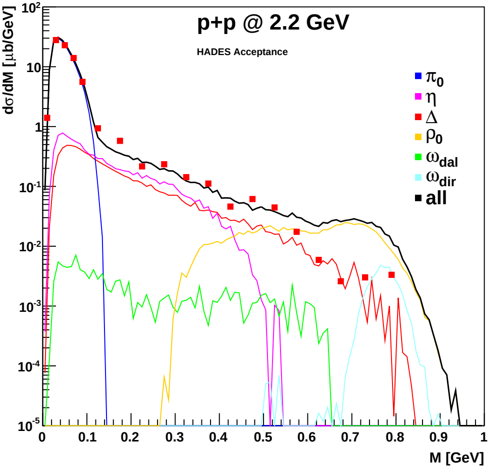
p + p @ 1.25 GeV

- $E_{lab} = 1.25$ GeV just below η threshold
- Small sub-threshold contribution from ρ expected
- Good agreement with data at low masses
- Too many dileptons from 0.3 GeV on $\rightarrow$ especially ρ^0



p + p @ 2.2 GeV

- Above η threshold, energy sufficient to reach pole mass of ω and ρ
- Overestimation of ρ contribution even larger than at 1.25 GeV

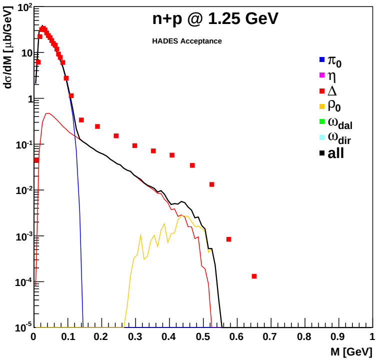


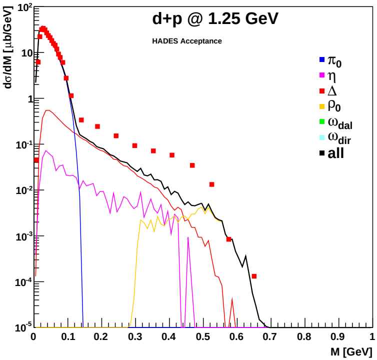
n + p @ 1.25 GeV

- Deuteron beam with 1.25AGeV has been used by HADES besides p+p
- Trigger on forward-going protons in order to select the (quasi-free) np collisions
- Fermi motion of the bound nucleons in the deuteron leads to a smearing of the NN collision energy $\rightarrow$ reaching above η-production threshold

$\Rightarrow$ One can not easily compare data with pure n+p simulations

n + p @ 1.25 GeV





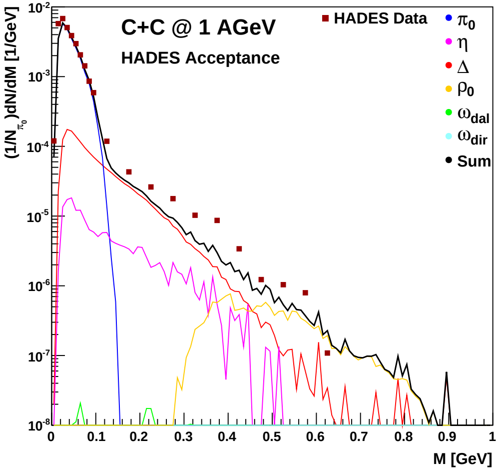
- η and more ρ are produced in d+p, compared to n+p
- However, even for d+p the Yield is underestimated by a factor of about 5 to 10

Study of A+A collisions

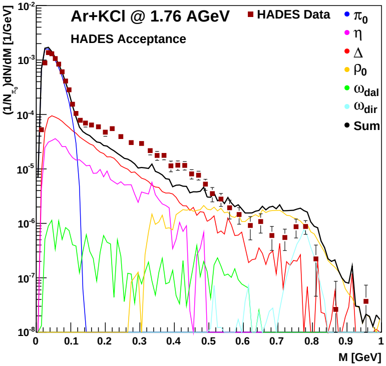
- In nucleus-nucleus collisions, additional effects compared to pp are expected
  - Fermi Momentum
  - Not only p+p, but also p+n and n+n collisions
  - Secondary interactions, depending on system size and energy
- Vector meson spectral functions may be changed in the medium
  - Shift of the pole mass (of the ρ)
  - Resonance melting in the medium
- In UrQMD, no such in-medium effects are implemented

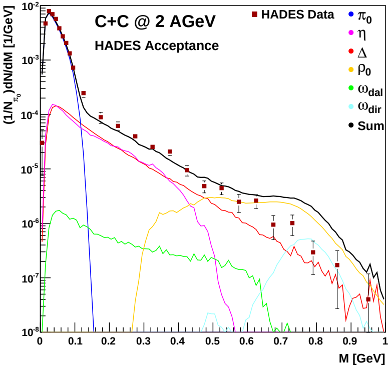
C + C @ 1 AGeV

- Here the ρ contribution fits quite well
- Underestimation below for energies above the pion peak
- Hardly any ω produced, η is relatively small



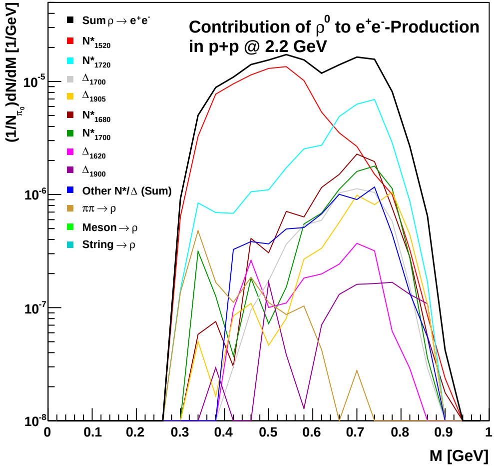
Ar + KCl @ 1.76 AGeV & C + C @ 2 AGeV





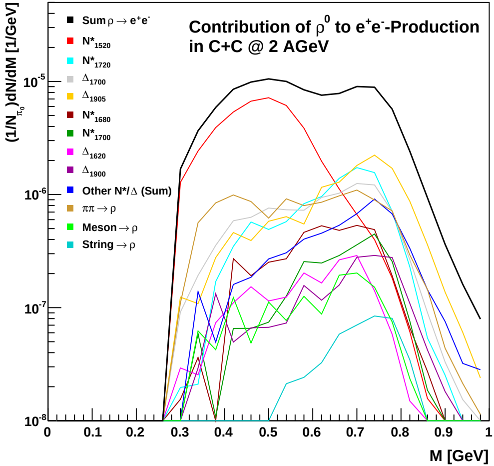
- As in elementary reactions, we get too many dileptons via the ρ^0 resonance, especially in the high-mass tail
- How are the ρ mesons produced?

ρ^0 Contribution



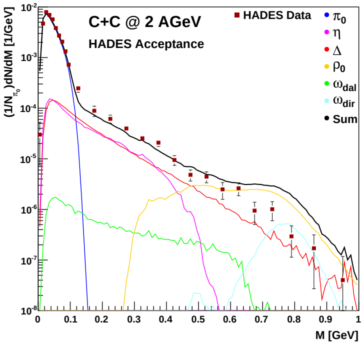
- Main contribution below pole mass by N^*_{1520} resonance
- For pole mass peak via N^*_{1680}, N^*_{1680} and N^*_{1700}
- $\pi\pi \to \rho$ negligible (but contributes!)

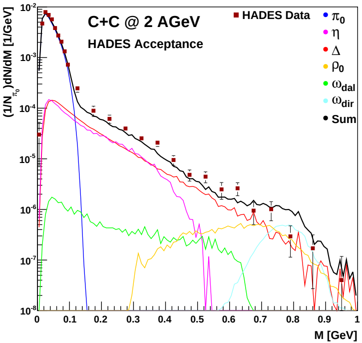
ρ^0 Contribution



- Main contribution below pole mass again by N^*_{1520}
- For pole mass peak via N^*_{1720}, Δ^*_{1700} and Δ^*_{1905}
- $\pi\pi \rightarrow \rho$ makes up already 10%

Comparison with / without Resonance Contribution



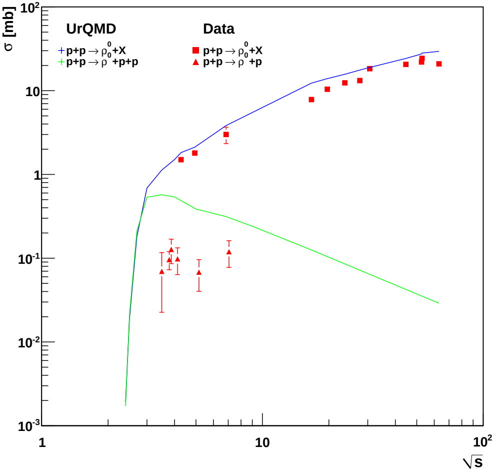


- Even a complete switch-off of all processes Δ / $N^* \to \rho$ gives too many dileptons in the high mass tail
- **Question: Why do we get so much ρ contribution ???**

Why too many dileptons from ρ ?

- All HADES energies are close to thresholds
    - Cross-sections change rapidly with small energy differences
    - Are the cross-sections in order?
- Do in-medium effects not included in UrQMD play a role ?
- Possible Σ channel for a part of what we treat as a ρ?
- Why do we see the large overestimation for invarinat masses above 700 GeV (pole mass)?

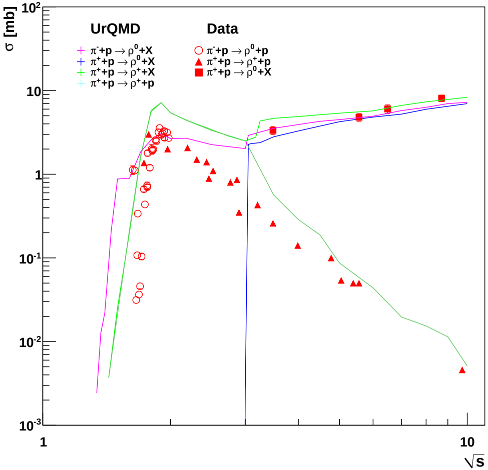
Cross-sections N+N $\rightarrow \rho^0$+X



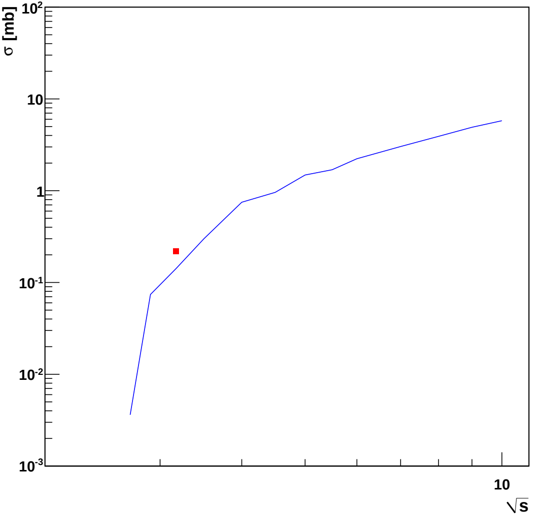
- Good description of inclusive cross-section
- For energies near threshold, σ might be too high, but no data are available for this region
- However, can't explain the overestimation at masses higher than the ρ pole mass

Cross-sections π+N $\rightarrow$ ρ+X

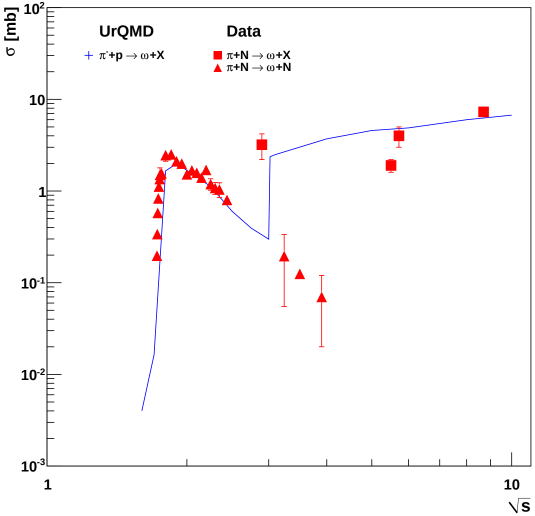
- ρ^+ production overestimated below $\sqrt{s} = 3$ GeV (but not relevant for dileptons)
- ρ^0 from UrQMD fits quite well to data, except for threshold region
- Does experiment just don't see these ρ which are significantly below pol mass?



Cross-sections N+N $\rightarrow \omega$+X



Cross-sections π+N $\rightarrow$ ω+X



Cross-sections

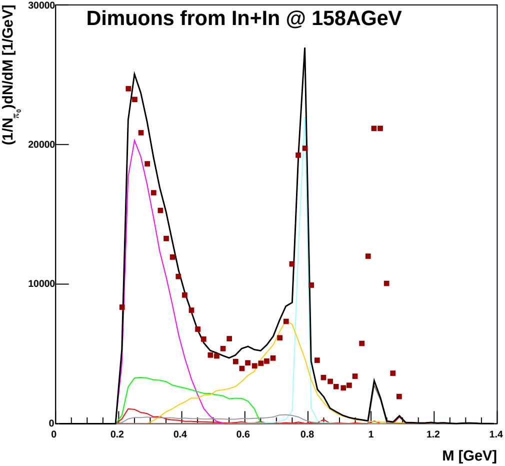
- Other cross-sections might as well play a role, especially R+N or R+R
- Not relevant for elementary, but significant for A+A collision
- Little data is available for these cases $\rightarrow$ Inclusion in UrQMD is is on a vague basis

$\Rightarrow$ Has to be checked next...

Going to ultra-relativistic Energies...

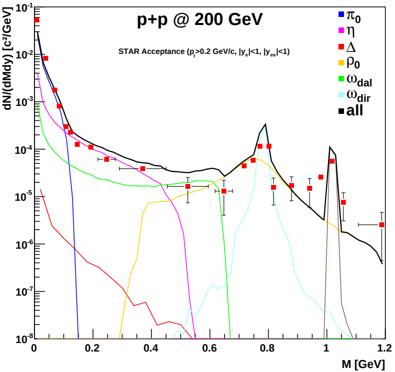
- Going to systems with a hot and dense fireball, does a pure transport approach still work to describe dilepton spectra?
- However, even at RHIC or LHC energies, low mass dileptons from the hadronic phase should dominate clearly over a possible QGP radiation
- In-medium effects are expected to become more dominant

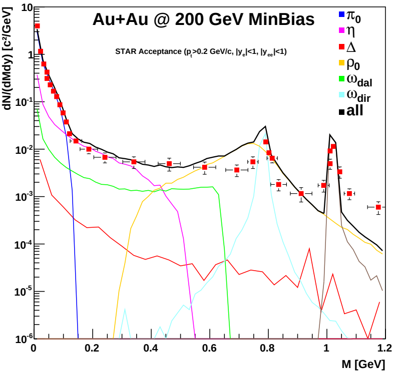
Comparison with NA60 Data



- UrQMD results for Dimuons at SPS energies show good agreement with experimental data
- Only for higher masses the yield is too low, ϕ production is underestimated

Comparison with STAR Data





- Same problem as at low energies: Too many e^+e^- via ρ decay
- However, completely different processes responsible here

Next Steps

- Reduce the overestimation dilepton production in UrQMD, especially via the ρ^0 (Δ and η could be optimezed as well)
- Coarse graining to be done for HADES energies
  - Take local temperature and and baryon chemical potential as functions of space and time
  - Accumulate an ensemble of events and determine local variables via coarse graining
- Dilepton calculation with hybrid model (transport + hydro)
  - Previous work (Dimuons from NA60) by Elvira Santini
  - Proceed with this work and calculate yields for RHIC and LHC ebergies