

Reonances in the medium

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based on

HvH, PRD **65**, 025010 (2001); PhD thesis 2000

Ivanov, Knoll, Voskresensky, NPA **657**, 413 (1999); NPA **672**, 313 (2000)

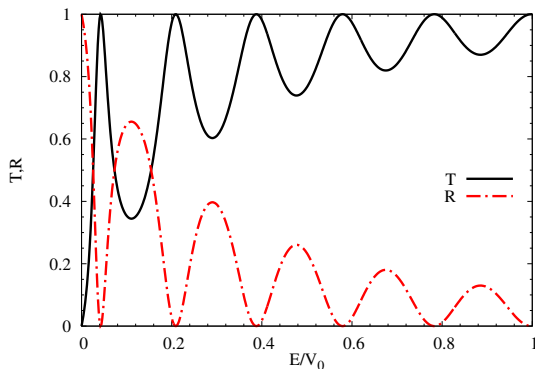
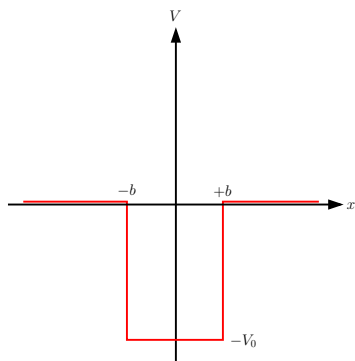
Knoll, Ivanov, Voskresensky, Ann. Phys. **293**, 126 (2001)



- 1 Warm-up: Resonance in quantum mechanics
- 2 Φ -derivable approximations
- 3 Transport equations
- 4 Summary

What's a resonance?

- quantum mechanics 101: Particle in a potential pot
- wave packet with energy around transmission-resonance peak
- nearly no reflection



What's a resonance?

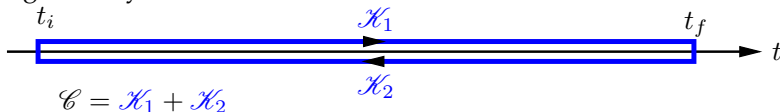
- quantum mechanics 101: Particle in a potential pot
- wave packet with energy around resonance peak
- nearly no reflection; stays a while in pot

Schwinger-Keldysh real-time formalism

- calculate **expectation values** of observables
- statistical operator defines state at **initial time, t_i** \Rightarrow “in-in formalism”
- time evolution

$$\langle O \rangle (t) = \text{Tr} \left[\hat{\rho}(t_i) \underbrace{\mathcal{T}_a \left\{ \exp \left[+i \int_{t_i}^t dt' \mathbf{H}_I(t') \right] \right\}}_{\text{anti time-ordered}} \mathbf{O}_I(t) \underbrace{\mathcal{T}_c \left\{ \exp \left[-i \int_{t_i}^t dt' \mathbf{H}_I(t') \right] \right\}}_{\text{time-ordered}} \right].$$

- Schwinger-Keldysh real-time contour:



Baym's Φ functional

- write generating functional for Green's functions as **path integral**
- introduce **local** and **bilocal** sources

$$Z[J, K] = N \int D\phi \exp \left[iS[\phi] + i \{J_1 \phi_1\}_1 + \left\{ \frac{i}{2} K_{12} \phi_1 \phi_2 \right\}_{12} \right]$$

- generating functional for **connected Green's functions**

$$W[J, K] = -i \ln Z[J, K]$$

- functional Legendre transform

$$\Gamma[\varphi, G] = W[J, K] - \{ \varphi_1 J_1 \}_1 - \frac{1}{2} \{ (\varphi_1 \varphi_2 + iG_{12}) K_{12} \}_{12}$$

- loop expansion

$$\Gamma[\varphi, G] = S_0[\varphi] + \frac{i}{2} \text{Tr} \ln(-iG^{-1}) + \frac{i}{2} \left\{ D_{12}^{-1} (G_{12} - D_{12}) \right\}_{12} \\ + \Phi[\varphi, G] \Leftarrow \text{all closed 2PI interaction diagrams}$$

$$D_{12}^{-1} = -\square - m^2$$

Baym's Φ functional

- equations of motion

$$\frac{\delta\Gamma}{\delta\varphi_1} = -J_1 - \{K_{12}\varphi_2\}_2 \stackrel{!}{=} 0, \quad \frac{\delta\Gamma}{\delta G_{12}} = -\frac{i}{2}K_{12} \stackrel{!}{=} 0$$

- mean field

$$-\square\varphi - m^2\varphi := j = -\frac{\delta\Phi}{\delta\varphi}$$

- “full” propagator $G \Rightarrow$ Dyson equation:

$$-i(D_{12}^{-1} - G_{12}^{-1}) := -i\Sigma = 2\frac{\delta\Phi}{\delta G_{21}}$$

- retarded Green's function for homogeneous system in momentum space

$$G_{\text{ret}}(p) = \frac{1}{p^2 - m^2 - \Sigma_{\text{ret}}(p)}$$

- spectral function

$$A(p) = -2 \text{Im} G_{\text{ret}}(p) = -2 \frac{\text{Im} \Sigma_{\text{ret}}(p)}{[p^2 - m^2 - \text{Re} \Sigma_{\text{ret}}(p)]^2 + [\text{Im} \Sigma_{\text{ret}}(p)]^2}$$

Properties of Φ -derivable approximations

- truncations of Φ functional \Rightarrow Φ -derivable approximations

$$\begin{aligned}
 i\Phi &= \text{[diagram: cross with 4 external lines]} + \text{[diagram: tadpole with 2 external lines]} + \text{[diagram: self-energy loop]} + \frac{1}{2} \text{[diagram: bubble with 2 external lines]} + \frac{1}{2} \text{[diagram: bubble with 2 internal lines]} + \frac{1}{3} \dots \\
 -ij &= \text{[diagram: cross with 3 external lines]} + \text{[diagram: tadpole with 1 external line]} + \text{[diagram: bubble with 1 external line]} + \dots \\
 -i\Sigma &= \text{[diagram: tadpole with 1 external line]} + \text{[diagram: tadpole with 1 external line]} + \text{[diagram: bubble with 2 external lines]} + \text{[diagram: bubble with 2 external lines]} + \dots
 \end{aligned}$$

- conservation laws for expectation values of **conserved quantities**
- in **thermal equilibrium** $i\Gamma = \ln Z$
- thermodynamic consistency: **bulk properties like pressure, energy, entropy** in accordance with dynamics
- same result from partition sum as from Green's functions!
- " Φ derivability" sufficient and necessary scheme!**

Transport equations

- start from Φ -derivable Dyson equation for Green's function

$$\begin{aligned}(\square_1 - \square_2)D^{12}(x_1, x_2) &= \int_{\mathcal{C}} dx_3 [\Sigma(x_1^1, x_3)D(x_3, x_2^2) - D(x_1^1, x_3)\Sigma(x_3, x_2^2)] \\ &= \text{Coll}(x_1^1, x_2^2)\end{aligned}$$

- assume smallness of space-time gradients in “collective macroscopic” variable $R = (x_1 + x_2)/2$
- Wigner transform of any two-point function, F

$$F(x_1, x_2) = \int \frac{d^4p}{(2\pi)^4} \exp[-ip \cdot (x_1 - x_2)] \tilde{F}\left(\frac{x_1 + x_2}{2}, p\right).$$

- assume space-time gradients wrt. R to be “small” \Rightarrow gradient expansion \Rightarrow “coarse graining”

$$2p \cdot \frac{\partial}{\partial R} iD^{12}(R, p) = \text{Coll}^{12}(R, p)$$

Gradient expansion of collision term

$$2p \cdot \frac{\partial}{\partial R} iD^{12}(R, p) = \text{Coll}^{12}(R, p)$$

- if Φ beyond pure two-point level \Rightarrow **memory + spatial correlations**
- simplify further by introducing $\text{Coll}_{\text{loc}}^{12}$:
diagrams evaluated at reference point R
- usual momentum Feynman rules with $D_{12}(R, p)$
- to have exact conservation laws add 1st-order ∂_R correction

$$D\left(\frac{x_i + x_j}{2}, p\right) \simeq D(R, p) + \frac{1}{2}[(x_i + x_j) - R] \cdot \frac{\partial}{\partial R} D(R, p)$$

- for **local** Green's functions and self-energies

$$iD^{12}(R, p) = f(R, p)A(R, p), \quad A(R, p) = -2 \text{Im} D_{\text{ret}}(R, p)$$

- as in equilibrium with **off-equilibrium phase-space distrib.** $f(R, p)$
- usual local **Dyson equation** for retarded Green's function

$$D_{\text{ret}}(R, p) = \frac{1}{p^2 - m^2 - \Sigma_{\text{ret}}(R, p)}$$

Diagrammar for gradient expansion

$$\overline{\overline{i \quad j}} = \frac{1}{2}(\partial_i + \partial_j)G(i, j) \rightarrow \partial_X G(X, p),$$

$$\overleftarrow{\overleftarrow{i \quad j}} = -i(x_i - x_j) \rightarrow -(2\pi)^4 \frac{\partial}{\partial p}$$

- arbitrary two-point function $M(x_1, x_2)$ with internal points x_3, \dots

$$\diamond\{M(1, 2)\} = \diamond \begin{array}{c} \square \\ \bullet \quad \bullet \\ 1 \quad 2 \end{array} \equiv \begin{array}{c} \square \\ \diamond \\ \bullet \quad \bullet \\ 1 \quad 2 \end{array} = \begin{array}{c} \overset{3}{\text{---}} \overset{4}{\text{---}} \\ \text{---} \text{---} \\ \bullet \quad \bullet \\ 1 \quad 2 \end{array} + \begin{array}{c} \overset{3}{\text{---}} \overset{4}{\text{---}} \\ \text{---} \text{---} \\ \bullet \quad \bullet \\ 1 \quad 2 \end{array}$$

$$M'(x_1, x_2; x_3, x_4) = \frac{\delta M(x_1, x_2)}{\delta iG(x_4, x_3)}$$

- collision term** \Rightarrow convolution integral

$$\diamond\{C(X, p)\} = \diamond \begin{array}{c} \diamond \quad \diamond \\ \bullet \quad \bullet \end{array}$$

$$= \begin{array}{c} \text{---} \text{---} \\ \text{---} \text{---} \\ \bullet \quad \bullet \end{array} + \begin{array}{c} \text{---} \text{---} \\ \text{---} \text{---} \\ \bullet \quad \bullet \end{array} + \begin{array}{c} \diamond \quad \diamond \\ \bullet \quad \bullet \end{array} + \begin{array}{c} \diamond \quad \diamond \\ \bullet \quad \bullet \end{array}$$

$$= \{A(X, p), B(X, p)\} + A(X, p)\diamond\{B(X, p)\} + \diamond\{A(X, p)\}B(X, p).$$

Diagrammar for gradient expansion

- **transport equation in Kadanoff-Baym form**

$$p \cdot \frac{\partial}{\partial R} iD^{12}(R, p) + \left\{ \text{Re } \Sigma_{\text{ret}}, iD^{12} \right\}_{\text{pb}} + \left\{ i\Sigma^{12}, \text{Re } D_{\text{ret}} \right\}_{\text{pb}} = C_{\text{loc}}^{12} + C_{\text{mem}}^{12}$$

- then **Noether currents exactly conserved also after gradient expansion**
- problem: 2nd Poisson bracket (“back-flow term”) cannot be represented in test-particle Monte Carlo
- Botermans-Malfliet ansatz

$$i\Sigma^{12}(R, p) = -f(R, p)\Gamma(R, p), \quad \Gamma(R, p) = -2 \text{Im } \Sigma_{\text{ret}}$$

- valid up to 1st-order gradients
- **Caveat:** in conservation laws from BM ansatz

$$A(R, p) \rightarrow \mathcal{B}(R, p) := \frac{\partial}{\partial p_0} \left[2 \text{Im } \ln(D_{\text{ret}}^{-1}) - \text{Re } G_{\text{ret}}\Gamma \right]$$

- for narrow resonances (BW approximation) $\mathcal{B} \simeq \frac{1}{2}A^2\Gamma$
- for test-particle off-shell method \Rightarrow see W. Cassing’s talk
- **Caveat:** possible trouble with **tachyons**
 - transition to semi-class. particle picture \leftrightarrow WKB/eikonal approximation
 - particle velocity \Rightarrow group velocity **superluminal around resonance**
 - no trouble in wave picture (see Sommerfeld+Brillouin 1913!)

Application: Lifetime of an “off-shell resonance”

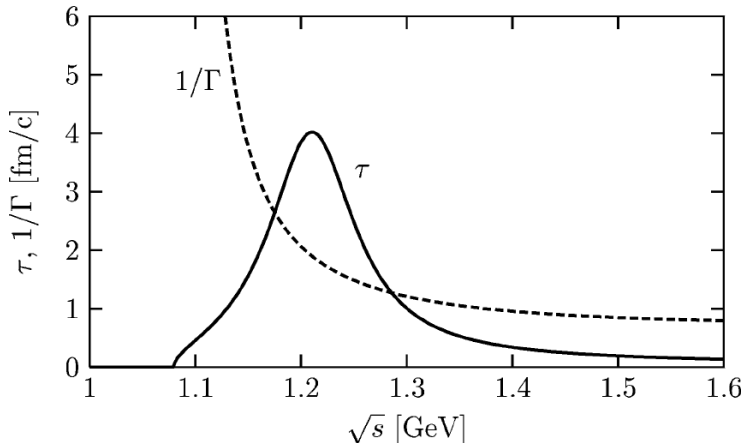
- in transport codes: resonances propagated as particles
- subject to decay with probability $\exp(-\Gamma\Delta t)$
- but $\Gamma = \Gamma(M)$ (vacuum) or even $\Gamma = \Gamma(M, \vec{p})$ (in med)
- in virial expansion (formally expansion of D around $D_{\text{vac}} \Rightarrow$ “thermodynamics” in terms of S matrix [Dashen, Ma, Bernstein 1969])
- correct lifetime from KB equations [Leupold, NPA 695, 377 (2001)]

$$\tau = 2p_0 \mathcal{B}_{\text{vac}} = \frac{\partial \delta}{\partial p_0}$$

- also from resonant wave propagation [Danielewicz, Pratt, PRC 53, 249 (1996)]
 \Rightarrow “delay time”: $\partial \delta / \partial E$

Application: Lifetime of an “off-shell resonance”

- example: $\Delta(1232)$ (from [Leupold, NPA 695, 377 (2001)])



- propagation of instable resonances great challenge for transport
- start from **self-consistent Φ derivable approximations**
- approximate **Kadanoff-Baym equations**
for Wigner transformed **single-particle GF**
- gradient expansion \Rightarrow coarse-grained dynamics
 \Rightarrow semi-classical transport equations
 \Rightarrow **positive phase-space distributions**
- Kadanoff-Baym form: exact conservation laws for Noether currents for
complete 1st-order gradient expansion
- Botermans-Malfliet form: feasibility as **test-particle MC**
- finite width \Rightarrow “off-shell potential”
- Caveat: **danger of superluminal particles**;
pragmatically solved in GiBUU, pHSD (...?)
- has intuitive physical interpretation (at least in simplifying limits)