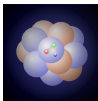


# Heavy-Quark Transport in the QGP

Hendrik van Hees

Justus-Liebig Universität Gießen

October 13, 2009

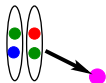


**Institut für  
Theoretische Physik**



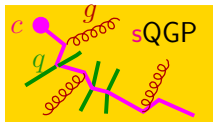
- 1 Heavy-quark interactions in the sQGP
  - Heavy quarks in heavy-ion collisions
  - Heavy-quark diffusion: The Langevin Equation
  - Elastic pQCD heavy-quark scattering
  - Non-perturbative interactions: Resonance Scattering
- 2 Non-photonic electrons at RHIC
- 3 Microscopic model for non-perturbative HQ interactions
  - Static heavy-quark potentials from lattice QCD
  - T-matrix approach
- 4 Summary and Outlook

# Heavy Quarks in Heavy-Ion collisions

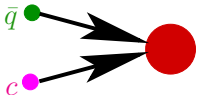


$c, b$  quark

hard production of HQs  
described by PDF's + pQCD (PYTHIA)

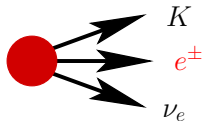


HQ rescattering in QGP: Langevin simulation  
drag and diffusion coefficients from  
microscopic model for HQ interactions in the sQGP



Hadronization to  $D, B$  mesons via  
quark coalescence + fragmentation

V. Greco, C. M. Ko, R. Rapp, PLB **595**, 202 (2004)



semileptonic decay  $\Rightarrow$   
"non-photonic" electron observables  
 $R_{AA}^{e^+e^-}(p_T), v_2^{e^+e^-}(p_T)$

# Relativistic Langevin process

- **Langevin process**: friction force + Gaussian random force
- in the (local) rest frame of the heat bath

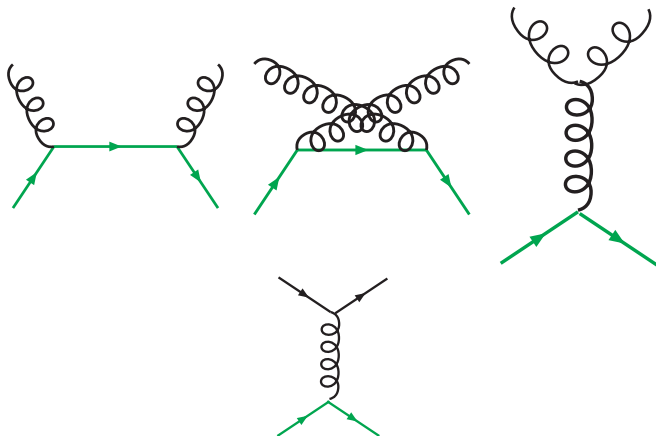
$$d\vec{x} = \frac{\vec{p}}{E_p} dt,$$

$$d\vec{p} = -A\vec{p}dt + \sqrt{2dt}[\sqrt{B_0}P_\perp + \sqrt{B_1}P_\parallel]\vec{w}$$

- $\vec{w}$ : normal-distributed random variable
- $A$ : friction (drag) coefficient
- $B_{0,1}$ : diffusion coefficients
- dependent on **realization of stochastic process**
- to guarantee correct equilibrium limit: Use **Hänggi-Klimontovich calculus**, i.e., use  $B_{0/1}(t, \vec{p} + d\vec{p})$
- Einstein dissipation-fluctuation relation  $B_0 = B_1 = E_p T A$ .
- to implement flow of the medium
  - use **Lorentz** boost to change into local “heat-bath frame”
  - use **update rule** in heat-bath frame
  - boost back into “lab frame”

# Elastic pQCD processes

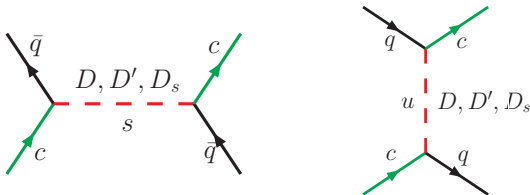
- Lowest-order matrix elements [Cambridge 79]



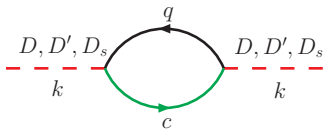
- **Debye-screening mass** for  $t$ -channel gluon exch.  $\mu_g = gT$ ,  $\alpha_s = 0.4$
- not sufficient to understand RHIC data on “non-photonic” electrons

# Non-perturbative interactions: Resonance Scattering

- General idea: Survival of  $D$ - and  $B$ -meson like **resonances** above  $T_c$
- **elastic** heavy-light-(anti-)quark scattering

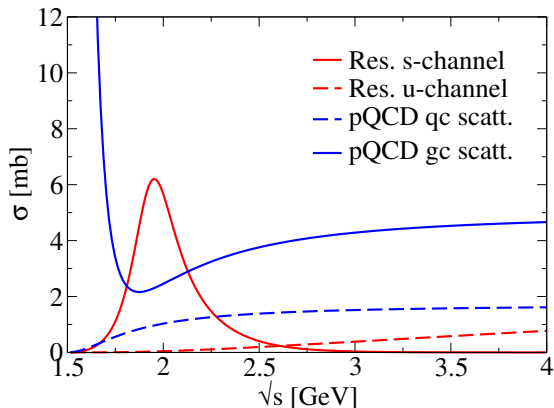


- $D$ - and  $B$ -meson like resonances in sQGP



- parameters
  - $m_D = 2 \text{ GeV}$ ,  $\Gamma_D = 0.4 \dots 0.75 \text{ GeV}$
  - $m_B = 5 \text{ GeV}$ ,  $\Gamma_B = 0.4 \dots 0.75 \text{ GeV}$

# Cross sections



- total pQCD and resonance cross sections: comparable in size
- BUT pQCD forward peaked  $\leftrightarrow$  resonance isotropic
- resonance scattering more effective for friction and diffusion

# Time evolution of the fire ball

- Elliptic **fire-ball** parameterization  
fitted to hydrodynamical flow pattern [Kolb '00]

$$V(t) = \pi(z_0 + v_z t)a(t)b(t), \quad a, b: \text{semi-axes of ellipse,}$$
$$v_{a,b} = v_\infty[1 - \exp(-\alpha t)] \mp \Delta v[1 - \exp(-\beta t)]$$

- **Isentropic expansion**:  $S = \text{const}$  (fixed from  $N_{\text{ch}}$ )
- **QGP Equation of state**:

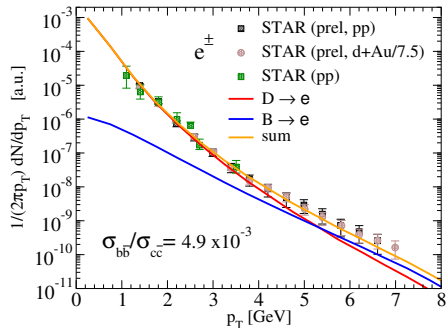
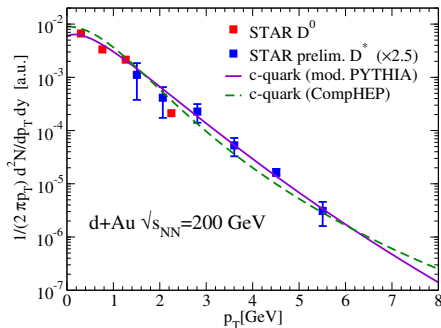
$$s = \frac{S}{V(t)} = \frac{4\pi^2}{90} T^3 (16 + 10.5n_f^*), \quad n_f^* = 2.5$$

- obtain  $T(t) \Rightarrow A(t, p)$ ,  $B_0(t, p)$  and  $B_1 = TEA$
- for semicentral collisions ( $b = 7$  fm):  $T_0 = 340$  MeV,  
QGP lifetime  $\simeq 5$  fm/ $c$ .
- simulate FP equation as **relativistic Langevin process**

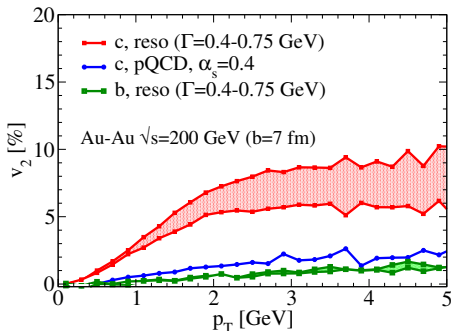
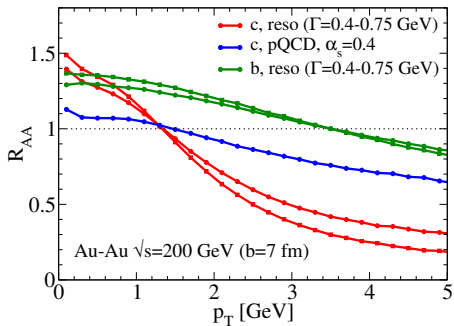


# Initial conditions

- need initial  $p_T$ -spectra of **charm** and **bottom** quarks
  - (modified) PYTHIA to describe exp. **D** meson spectra, assuming  $\delta$ -function fragmentation
  - exp. **non-photonic single- $e^\pm$**  spectra: Fix bottom/charm ratio

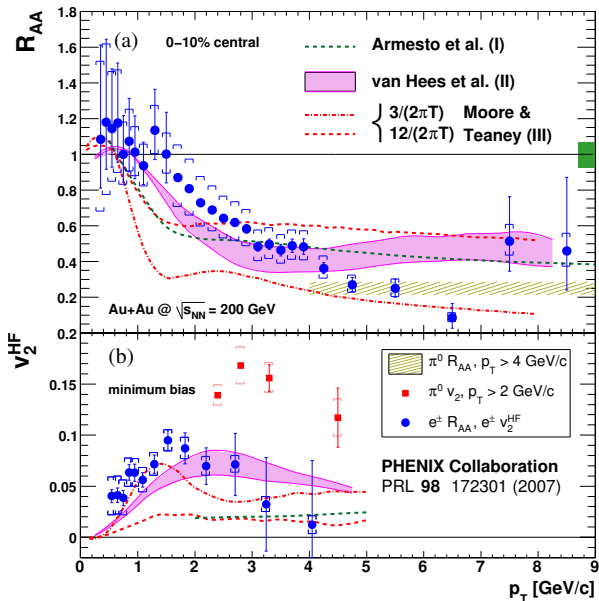


# Spectra and elliptic flow for heavy quarks

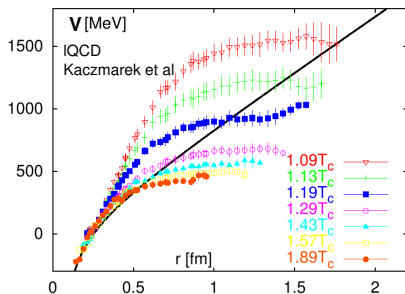


- $\mu_D = gT$ ,  $\alpha_s = g^2/(4\pi) = 0.4$
- resonances  $\Rightarrow$   $c$ -quark thermalization  
without upscaling of cross sections
- Fireball parametrization consistent with hydro

# Comparison to single-electron spectra @ RHIC



# Microscopic model: Static potentials from lattice QCD



- color-singlet free energy from lattice
- use **internal energy**

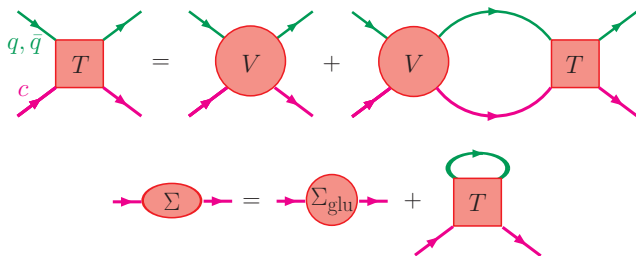
$$U_1(r, T) = F_1(r, T) - T \frac{\partial F_1(r, T)}{\partial T},$$

$$V_1(r, T) = U_1(r, T) - U_1(r \rightarrow \infty, T)$$

- Casimir scaling for other color channels [Nakamura et al 05; Döring et al 07]

$$V_{\bar{3}} = \frac{1}{2}V_1, \quad V_6 = -\frac{1}{4}V_1, \quad V_8 = -\frac{1}{8}V_1$$

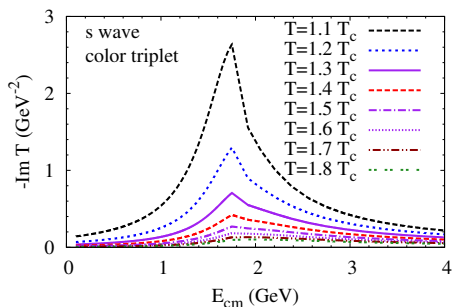
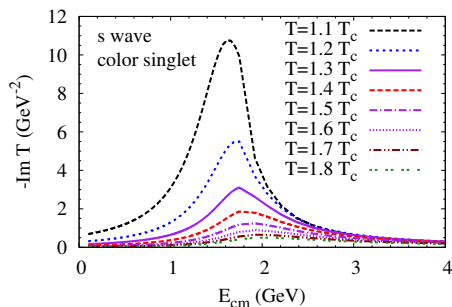
- Brueckner many-body approach for elastic  $Qq, Q\bar{q}$  scattering



- reduction scheme: 4D Bethe-Salpeter  $\rightarrow$  3D Lipmann-Schwinger
- $S$ - and  $P$  waves
- same scheme for light quarks (self consistent!)
- Relation to invariant **matrix elements**

$$\sum |\mathcal{M}(s)|^2 \propto \sum_q d_a (|T_{a,l=0}(s)|^2 + 3|T_{a,l=1}(s)|^2 \cos^2 \theta_{\text{cm}})$$

# T-matrix

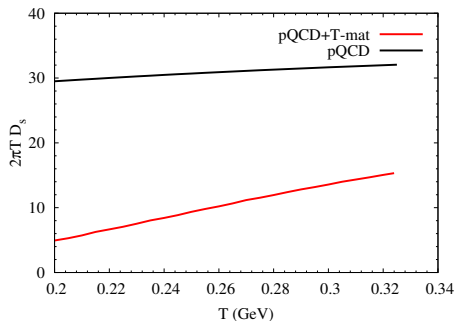
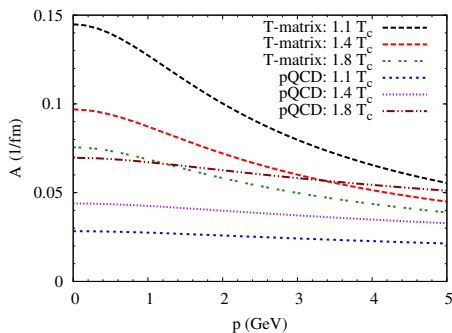


- resonance formation at lower temperatures  $T \simeq T_c$
- melting of resonances at higher  $T$ !  $\Rightarrow$  sQGP
- $P$  wave smaller
- resonances near  $T_c$ : natural connection to quark coalescence

[Ravagli, Rapp 07; Ravagli, HvH, Rapp 08]

- model-independent assessment of elastic  $Qq$ ,  $Q\bar{q}$  scattering
- problems: uncertainties in extracting potential from IQCD  
in-medium potential  $V$  vs.  $F$ ?

# Transport coefficients



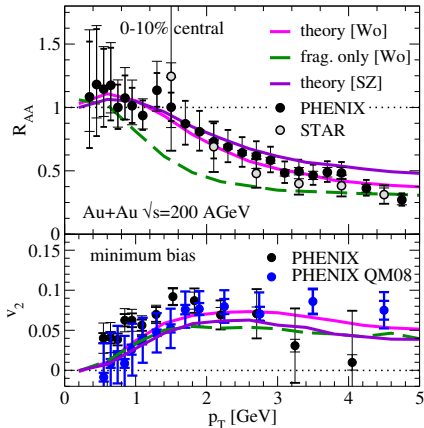
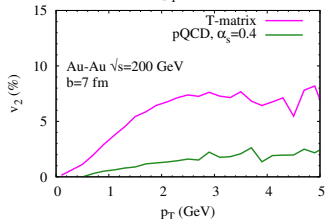
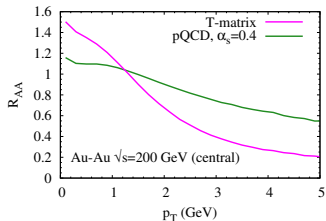
- from **non-pert.** interactions reach  $A_{\text{non-pert}} \simeq 1/(7 \text{ fm}/c) \simeq 4A_{\text{pQCD}}$
- **A decreases with higher temperature**
- higher density (over)compensated by **melting of resonances!**
- spatial diffusion coefficient

$$D_s = \frac{T}{mA}$$

**increases** with temperature

# Non-photonic electrons at RHIC

- same model for bottom
- quark **coalescence**+**fragmentation**  $\rightarrow D/B \rightarrow e + X$



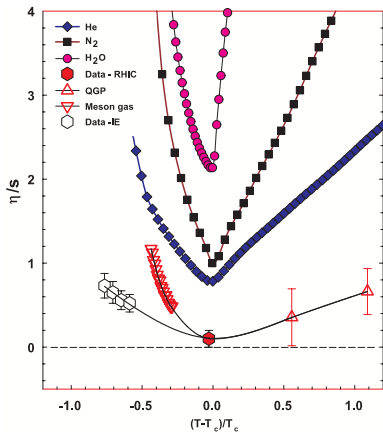
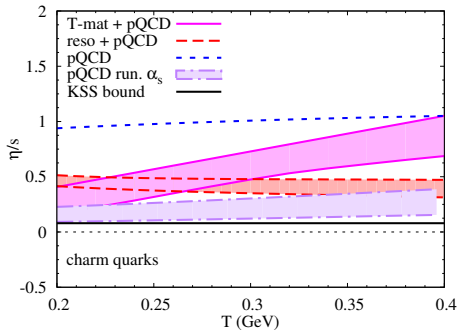
- **coalescence crucial for description of data**
- increases **both**,  $R_{AA}$  and  $v_2 \Leftrightarrow$  “momentum kick” from light quarks!
- “resonance formation” **towards  $T_c$**   $\Rightarrow$  **coalescence natural** [Ravagli, Rapp 07]



# Transport properties of the sQGP

- spatial diffusion coefficient: **Fokker-Planck**  $\Rightarrow D_s = \frac{T}{m_A} = \frac{T^2}{D}$
- measure for coupling strength in plasma:  $\eta/s$

$$\frac{\eta}{s} \simeq \frac{1}{2} T D_s \quad (\text{AdS/CFT}), \quad \frac{\eta}{s} \simeq \frac{1}{5} T D_s \quad (\text{wQGP})$$



[Lacey, Taranenko (2006)]

# Summary and Outlook

## • Summary

- Heavy quarks in the sQGP
- non-perturbative interactions
  - mechanism for strong coupling: resonance formation at  $T \gtrsim T_c$
  - IQCD potentials parameter free
  - res. melt at higher temperatures  $\Leftrightarrow$  consistency betw.  $R_{AA}$  and  $v_2$ !
- also provides “natural” mechanism for quark coalescence
- resonance-recombination model [L. Ravagli, HvH, R. Rapp, Phys. Rev. C **79**, 064902 (2009)]
- problems
  - extraction of  $V$  from lattice data
  - potential approach at finite  $T$ :  $F$ ,  $V$  or combination?

## • Outlook

- include inelastic heavy-quark processes (gluo-radiative processes)
- other heavy-quark observables like charmonium suppression/regeneration