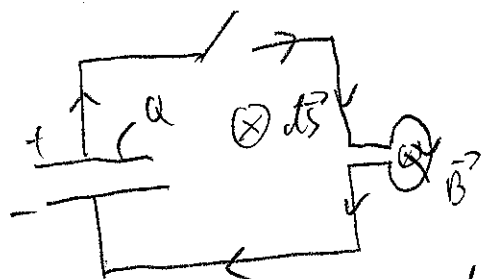


Solutions of problems chpt. 12

(1)



$$\oint \vec{E} \cdot d\vec{r} = -\frac{Q}{C} = -L \frac{di}{dt}$$

$$i = -\frac{dQ}{dt} \Rightarrow -L \frac{di}{dt} = -L \frac{d^2Q}{dt^2} = \frac{Q}{C}$$

$$\frac{d^2Q}{dt^2} = -\frac{1}{LC} Q$$

$$Q(t) = A \cos(\omega t) + B \sin(\omega t)$$

$$\dot{Q}(t) = -A\omega \sin(\omega t) + B\omega \cos(\omega t) = -i(t)$$

$$\ddot{Q}(t) = -A\omega^2 \cos(\omega t) - B\omega^2 \sin(\omega t) = -\frac{1}{LC} (A \cos(\omega t) + B \sin(\omega t))$$

$$\Rightarrow \omega^2 = \frac{1}{LC} \Rightarrow \omega = \frac{1}{\sqrt{LC}}$$

Initial conditions:

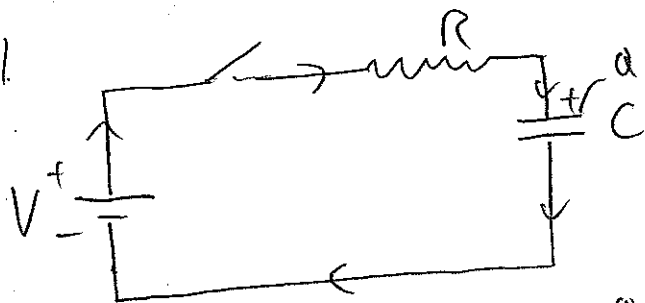
$$Q(t=0) = Q_0; \quad i(t=0) = 0$$

$$\Rightarrow A = Q_0 \quad \Rightarrow B = 0$$

$$Q(t) = Q_0 \cos(\omega t)$$

$$i(t) = -\dot{Q}(t) = Q_0 \omega \sin(\omega t)$$

(2)



$$\oint d\vec{r} \vec{E} = -V + Ri + \frac{Q}{C} = 0; \quad i = \frac{dQ}{dt}$$

$$\Rightarrow R \frac{dQ}{dt} + \frac{Q}{C} = V$$

Homogeneous equation

$$R \dot{q}_H + \frac{q}{C} = 0 \Rightarrow \dot{q}_H = -\frac{q}{RC}$$

$$q_H(t) = A \exp\left(-\frac{t}{RC}\right)$$

Particular solution of inhomogeneous eq.

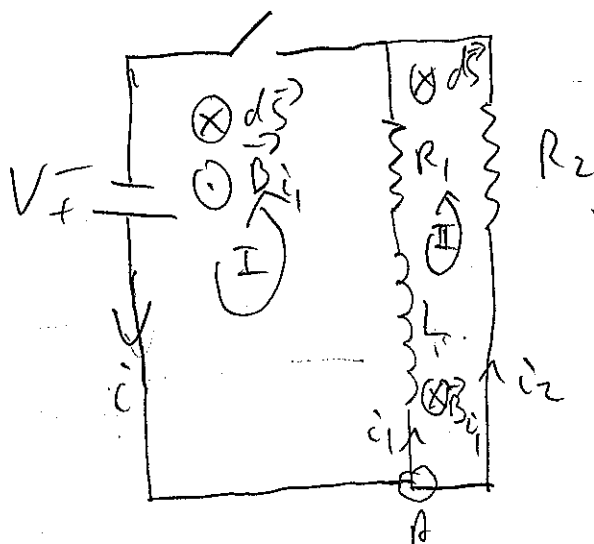
$$q_I = \text{const} \Rightarrow \frac{dq_I}{dt} = 0 \Rightarrow q_I = CV$$

$$\Rightarrow q(t) = q_I + q_H(t) = CV + A \exp\left(-\frac{t}{RC}\right)$$

$$q(0) = q_0 \Rightarrow CV + A = q_0 \Rightarrow A = q_0 - CV$$

$$\Rightarrow q(t) = CV + (q_0 - CV) \exp\left(-\frac{t}{RC}\right)$$

(3)



$$i = i_1 + i_2$$

$$\textcircled{1}: \oint d\vec{r} \cdot \vec{E} = -V + i_1 R_1 = -L \frac{di_1}{dt}$$

$$\textcircled{2}: \oint d\vec{r} \cdot \vec{E} = R_2 i_2 - R_1 i_1 = +L \frac{di_1}{dt}$$

Add $\textcircled{1}$ and $\textcircled{2}$

$$-V + R_2 i_2 = 0 \Rightarrow i_2 = \frac{V}{R_2}$$

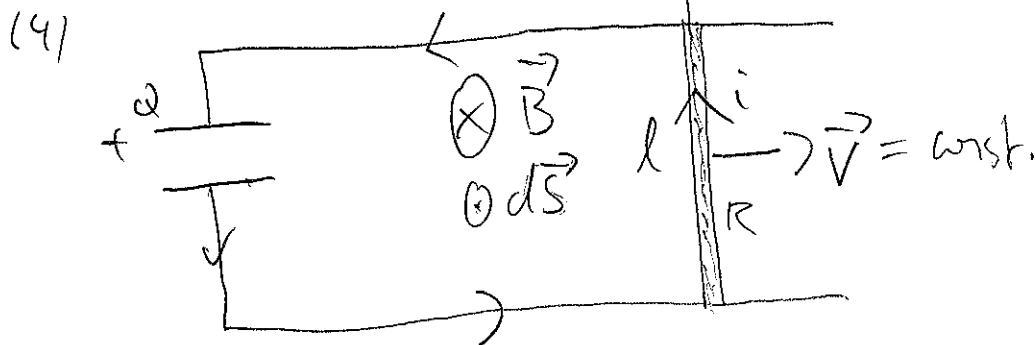
For i_1 we solve $\textcircled{1}$. Now we know the particular solution for the RL circuit and can work with energy

$$i_1(t) = \frac{V}{R_1} + A \exp\left(-\frac{R_1}{L} t\right)$$

Initial condition $i_1(t=0) = 0$

(3)

$$\Rightarrow i_1(t) = \frac{V}{R_1} [1 - \exp(-\frac{R_1}{L} t)]$$



No self-inductance \Rightarrow

$$\oint d\vec{r} \cdot \vec{E} = iR + \frac{d\mathcal{Q}}{dt} = l v B$$

$$i = \frac{d\mathcal{Q}}{dt} = \dot{\mathcal{Q}}$$

$$\Rightarrow \dot{\mathcal{Q}} + \frac{\mathcal{Q}}{RC} = \frac{l v B}{R}$$

$$\mathcal{Q}(t) = A \exp(-\frac{t}{RC}) + l v B C$$

$$\mathcal{Q}(t=0) = 0$$

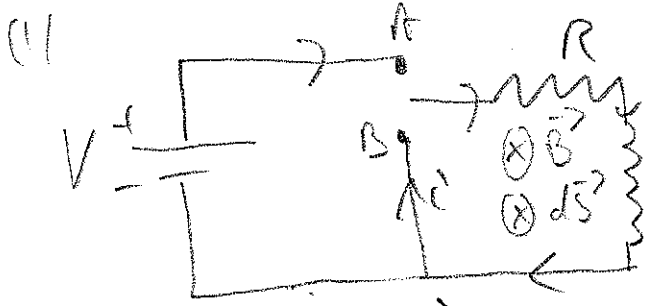
$$\mathcal{Q}(t) = l v B C [1 - \exp(-\frac{t}{RC})]$$

$$(5) \quad i \dot{\mathcal{Q}} + \frac{\mathcal{Q}}{C} = l v B - \frac{d}{dt}(L i)$$

Here L changes with time either, because the geometry of the loop changes. We cannot calculate it with our means.

Exercises - Chapter 12

(1)



If switch is in position B, we have

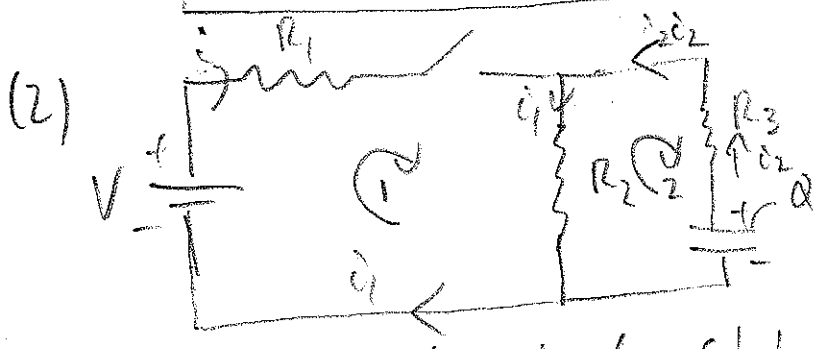
$$\oint \vec{E} \cdot d\vec{r} = Ri = -L \frac{di}{dt}$$

$$\Rightarrow \text{diff eq is } \exp\left(-\frac{R}{L}t\right)$$

The initial condition is determined by the steady state when the switch is in position A for a long time:

$$i_0 = \frac{V}{R}$$

$$\Rightarrow i(t) = \frac{V}{R} \exp\left(-\frac{R}{L}t\right)$$



First we solve the steady state problem when the switch is closed for a long time. Then $i_2 = 0$ and $i = i_1$

Then $\oint \vec{E} \cdot d\vec{r} = -V + R_1 i_1 + R_2 i_1 = 0$

$$\Rightarrow i_1 = \frac{V}{R_1 + R_2}$$

$$\textcircled{2} \Rightarrow \oint \vec{d}\vec{r} \vec{E} = \frac{Q}{C} - R_3 i_2 - R_2 i_1 = 0$$

because of $i_1 = 0 \Rightarrow Q = C R_2 i_2 = C R_2 \frac{V}{R_1 + R_2}$

If the switch is opened, $i = 0 = i_1 - i_2 \Rightarrow i_1 = i_2$
and loop $\textcircled{2}$ gives

$$\frac{Q}{C} = (R_2 + R_3) i_2$$

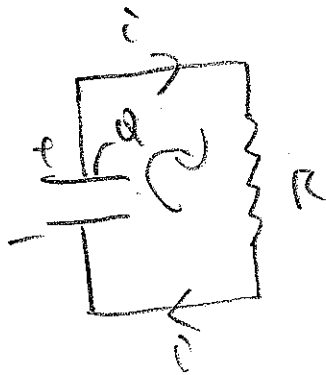
but in our direction of i_2 : $i_2 = -\frac{dQ}{dt} = -\dot{Q}$

$$\Rightarrow \dot{Q} = -\frac{Q}{C(R_2 + R_3)}$$

$$\Rightarrow Q(t) = Q_0 \exp\left(-\frac{t}{C(R_2 + R_3)}\right)$$

$$Q(t) = \frac{C R_2 V}{R_1 + R_2} \exp\left[-\frac{t}{C(R_2 + R_3)}\right]$$

(3)



$$\oint \vec{d}\vec{r} \vec{E} = iR - \frac{Q}{C} = 0$$

$$i = -\frac{dQ}{dt} = -\dot{Q}$$

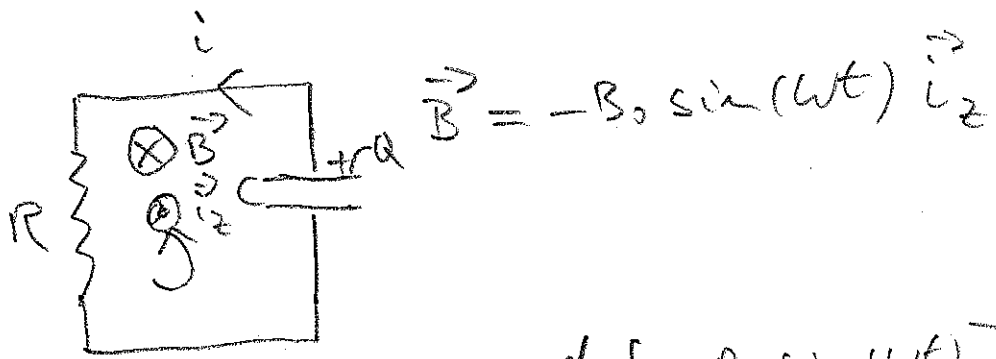
$$\Rightarrow \dot{Q} = -\frac{Q}{RC}$$

$$Q(t) = Q_0 \exp\left(-\frac{t}{RC}\right)$$

$$i(t) = -\dot{Q}(t) = \frac{Q_0}{RC} \exp\left(-\frac{t}{RC}\right)$$

(4)

(3)



$$\oint \vec{E} \cdot d\vec{r} = -\frac{dQ}{C} + iR = \frac{d}{dt} [A B_0 \sin(\omega t)]$$

$$= A B_0 \omega \cos(\omega t)$$

$$i = -\frac{dQ}{dt}$$

$$\Rightarrow R\dot{Q} + \frac{Q}{C} = -A B_0 \omega \cos(\omega t)$$

Homogeneous equation

$$\dot{Q}_H = -\frac{Q_H}{RC} \Rightarrow Q_H(t) = a \exp\left(-\frac{t}{RC}\right)$$

Inhomogeneous equation

Ansatz: $Q_I(t) = \alpha \cos(\omega t) + \beta \sin(\omega t)$

$$\dot{Q}_I(t) = -\alpha \omega \sin(\omega t) + \beta \omega \cos(\omega t)$$

$$R[-\alpha \omega \sin(\omega t) + \beta \omega \cos(\omega t)] + \frac{1}{C}[\alpha \cos(\omega t) + \beta \sin(\omega t)] = -A B_0 \omega \cos(\omega t)$$

$$\Rightarrow -R\alpha\omega + \frac{\beta}{C} = 0 \Rightarrow \beta = R\omega C\alpha$$

$$\frac{\alpha}{C} + \beta\omega R = -A B_0 \omega$$

$$\alpha \left(\frac{1}{C} + R^2\omega^2 C\right) = -A B_0 \omega$$

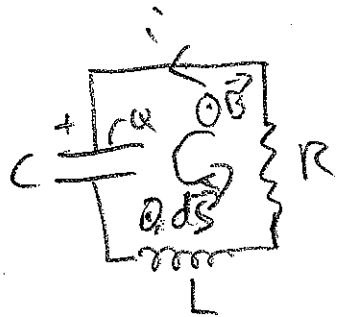
$$A = - \frac{A B_0 \omega C}{1 + (R\omega C)^2}$$

$$B = - \frac{R A \omega^2 C^2 B_0}{1 + (R\omega C)^2}$$

$$Q(t) = A \cos(\omega t) + B \sin(\omega t) - A \exp\left(-\frac{t}{RC}\right)$$

(from initial condition $Q(t=0) = 0$).

(5) The full story is as follows



$$\oint d\vec{r} \cdot \vec{E} = \frac{Q}{C} + Ri = -L \frac{di}{dt}$$

$$i = + \frac{dQ}{dt} = \dot{Q}$$

$$\Rightarrow L \ddot{Q} + R \dot{Q} + \frac{Q}{C} = 0$$

We expect a damped harmonic behaviour. Thus the right ansatz is

~~$$Q(t) = \exp(-\beta t) [A \cos(\omega t) + B \sin(\omega t)]$$~~

We must have two linearly independent functions and two integration constants, because we have a 2nd order linear ODE.

Now:

~~$$Q = -\beta \exp(-\beta t) [A \cos(\omega t) + B \sin(\omega t)] + \exp(-\beta t) [-A \omega \sin(\omega t) + B \omega \cos(\omega t)]$$~~

~~$$\begin{aligned} \ddot{Q} &= \exp(-\beta t) [(B\omega - \beta A) \cos(\omega t) - (B\beta + A\omega) \sin(\omega t)] \\ &+ \exp(-\beta t) [-\omega(B\omega - \beta A) \sin(\omega t) + \omega(B\beta + A\omega) \cos(\omega t)] \end{aligned}$$~~

$$Q(t) = \exp(-\beta t) q(t)$$

$$\dot{Q}(t) = -\beta \exp(-\beta t) q + \exp(-\beta t) \dot{q}$$

$$= \exp(-\beta t) (\dot{q} - \beta q)$$

$$\ddot{Q}(t) = -\beta \exp(-\beta t) (\dot{q} - \beta q) + \exp(-\beta t) (\ddot{q} - \beta \dot{q})$$

$$= \exp(-\beta t) [\ddot{q} - 2\beta \dot{q} + \beta^2 q]$$

$$\Rightarrow \exp(-\beta t) [L(\ddot{q} - 2\beta \dot{q} + \beta^2 q) + R(\dot{q} - \beta q) + \frac{q}{C}] = 0$$

This is simplifying, if the coefficient in front of \dot{q} is made to vanish:

$$-2\beta L + R = 0 \Rightarrow \beta = \frac{R}{2L}$$

Then the remaining ODE for q reads

$$L\ddot{q} + (L\beta^2 - \beta R + \frac{1}{C})q = 0$$

$$L\ddot{q} + (\frac{R^2}{4L} - \frac{R^2}{2L} + \frac{1}{C})q = 0$$

$$L\ddot{q} + (\frac{1}{C} - \frac{R^2}{4L})q = 0$$

$$\Rightarrow \ddot{q} = -(\frac{1}{LC} - \frac{R^2}{4L^2})q$$

Now we have to consider three cases

$$(a) \frac{1}{LC} - \frac{R^2}{4L^2} > 0.$$

Then we have

$$\ddot{q} = -\omega^2 q \text{ with } \omega = \sqrt{\frac{1}{LC} - \frac{R^2}{4L^2}}$$

The general solution in this case is

$$q(t) = A \cos(\omega t) + B \sin(\omega t)$$

$$(b) \frac{1}{LC} - \frac{R^2}{4L^2} = 0$$

$$\Rightarrow \ddot{q} = 0 \Rightarrow q(t) = At + B$$

$$(c) \frac{1}{LC} - \frac{R^2}{4L^2} < 0$$

$$\Rightarrow \ddot{q} = -\lambda^2 q \text{ with } \lambda = \sqrt{\frac{R^2}{4L^2} - \frac{1}{LC}}$$

Then the general solution is

$$q(t) = A \exp(+\lambda t) + B \exp(-\lambda t)$$

Solution for original equation

case (a)

$$q(t) = \exp\left(-\frac{Rt}{2L}\right) [A \cos(\omega t) + B \sin(\omega t)]$$

$$\text{with } \omega = \sqrt{\frac{1}{LC} - \frac{R^2}{4L^2}}$$

Initial conditions

$$q(t=0) = q_0$$

$$\dot{q}(t=0) = 0$$

$$q(t=0) = A = q_0$$

$$\dot{q}(t) = -\frac{R}{2L} \exp\left(-\frac{Rt}{2L}\right) [A \cos(\omega t) + B \sin(\omega t)]$$

$$\exp\left(-\frac{Rt}{2L}\right) [-A\omega \sin(\omega t) + B\omega \cos(\omega t)]$$

$$\dot{q}(t=0) = -\frac{R}{2L} A + B\omega = 0 \Rightarrow B = \frac{R}{2L\omega}$$

$$q(t) = \exp\left(-\frac{Rt}{2L}\right) \left[q_0 \cos(\omega t) + \frac{R}{2L\omega} \sin(\omega t) \right]$$

Case (A)

$$q(t) = \exp\left(-\frac{Rt}{2L}\right) (At + B)$$

$$q(t=0) = q_0 = B$$

$$\dot{q}(t) = \exp\left(-\frac{Rt}{2L}\right) \left[-\frac{R}{2L}(At + B) + A\right]$$

$$\dot{q}(t=0) = A - \frac{R}{2L}B = 0 \Rightarrow A = \frac{R}{2L}B = \frac{R}{2L}q_0$$

$$q(t) = q_0 \exp\left(-\frac{Rt}{2L}\right) \left[\frac{R}{2L}t + 1\right]$$

Case (C)

$$q(t) = \exp\left(-\frac{Rt}{2L}\right) [A \exp(\lambda t) + B \exp(-\lambda t)]$$

$$\text{with } \lambda = \sqrt{\frac{R^2}{4L^2} - \frac{1}{LC}}$$

$$q(t) = A \exp\left[\left(\lambda - \frac{R}{2L}\right)t\right] + B \exp\left[-\left(\lambda + \frac{R}{2L}\right)t\right]$$

Note that this is always damped, because

$$-\frac{R}{2L} \pm \lambda < 0$$

$$\dot{q}(t) = A \left(\lambda - \frac{R}{2L}\right) \exp\left[\left(\lambda - \frac{R}{2L}\right)t\right] - B \left(\lambda + \frac{R}{2L}\right) \exp\left[-\left(\lambda + \frac{R}{2L}\right)t\right]$$

$$q(t=0) = A + B = q_0$$

$$\dot{q}(t=0) = \lambda(A - B) - \frac{R}{2L}(A + B) = 0$$

$$A - B = \frac{R}{2L\lambda} q_0$$

$$A = \frac{1}{2} a_0 \left(1 + \frac{R}{2L\lambda} \right)$$

$$B = \frac{1}{2} a_0 \left(1 - \frac{R}{2L\lambda} \right)$$