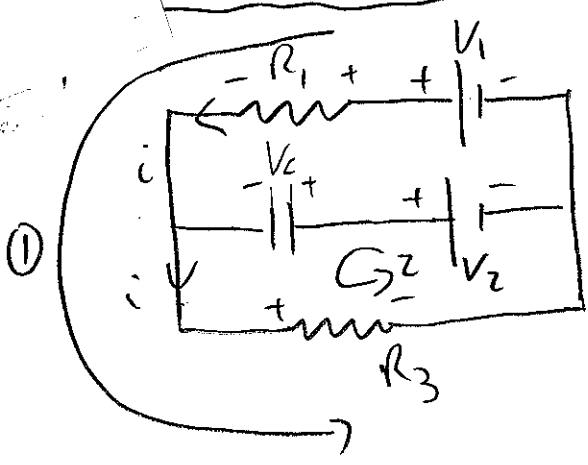


see exam III

①



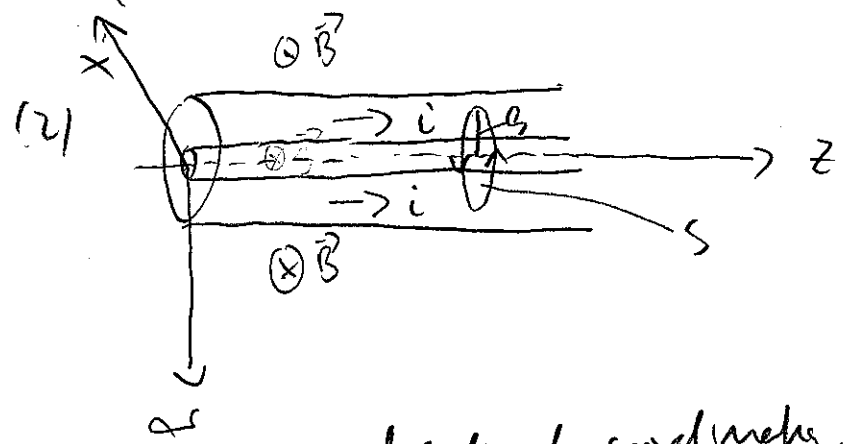
$$1: V_1 - iR_1 - iR_3 = 0$$

$$2: V_2 - V_c - iR_3 = 0$$

no current through capacitor in the steady state!

$$\textcircled{1} \Rightarrow i = \frac{V_1}{R_1 + R_3}; \quad V_c = V_2 - iR_3 = V_2 - \frac{R_3}{R_1 + R_3} V_1$$

$$Q = CV_c = C \left(V_2 - \frac{R_3}{R_1 + R_3} V_1 \right)$$



In our usual cylindrical coordinates, for \vec{B} we expect the form

$$\vec{B} = B_\phi(\rho) \hat{e}_\phi$$

To find B_ϕ we use Ampère's Law

$$\oint_{\partial S} \vec{B} \cdot d\vec{r} = \int_S d\vec{S} \cdot \vec{j} \mu_0$$

For S we use a circular disc \parallel to the xy plane with radius ρ : Then we know, with ϕ as the parameter of ∂S

$$\vec{r} = s \vec{e}_s = s (\cos\varphi \vec{i}_x + \sin\varphi \vec{i}_y)$$

$$d\vec{r} = ds \vec{e}_s = ds (-\sin\varphi \vec{i}_x + \cos\varphi \vec{i}_y)$$

$$\oint_{\partial S} d\vec{r} \cdot \vec{B} = \int_0^{2\pi} ds \int_0^a s \vec{e}_\varphi B_\varphi(s) \vec{i}_\varphi = 2\pi s B_\varphi(s)$$

The current density vector is piecewise

$$\vec{j} = \begin{cases} 0 & \text{for } 0 \leq s < a \\ \frac{i \vec{i}_z}{\pi(b^2 - a^2)} & \text{for } a \leq s < b \\ 0 & \text{for } s > b \end{cases}$$

Thus

$$\int_S d\vec{S} \cdot \vec{j} = \begin{cases} 0 & \text{for } 0 \leq s < a \\ \frac{i\pi(s^2 - a^2)}{\pi(b^2 - a^2)} = i \frac{s^2 - a^2}{b^2 - a^2} & \text{for } a \leq s < b \\ i & \text{for } s > b \end{cases}$$

$$\Rightarrow B_\varphi(s) = \begin{cases} 0 & \text{for } 0 \leq s < a \\ \frac{\mu_0 i}{2\pi s} \frac{s^2 - a^2}{b^2 - a^2} & \text{for } a \leq s < b \\ \frac{\mu_0 i}{2\pi s} & \text{for } s > b \end{cases}$$

(a) For a particle with $\vec{v} = v \vec{i}_z$ along $s=0$ we have

$$\vec{F}_{\text{mag}} = q \vec{v} \times \vec{B} = 0$$

because at $s=0: \vec{B} = 0$

(b) For the particle outside the cylinder we have

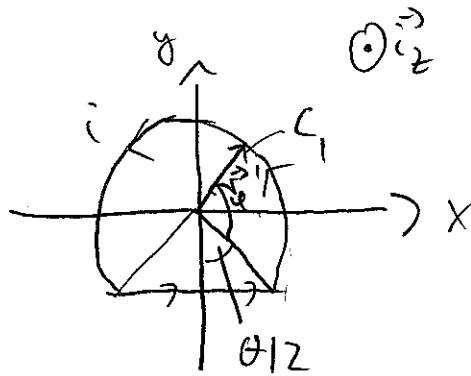
$$\vec{v} = v_1 \vec{i}_z + v_2 \vec{i}_s \quad ; \quad \vec{r} = s \vec{e}_s + z \vec{i}_z \quad (\text{with } s > b)$$

$$\begin{aligned} \vec{F}_{\text{mag}} &= q \vec{v} \times \vec{B} = q (v_1 \vec{i}_z + v_2 \vec{i}_s) \times B_\varphi \vec{i}_\varphi \\ &= -q v_1 B_\varphi \vec{i}_s + q v_2 B_\varphi \vec{i}_z \end{aligned}$$

$$\vec{F}_{mg} = \frac{4\mu_0 i}{2\pi R} (-\nu_1 \vec{i}_y + \nu_2 \vec{i}_z)$$

(3)

(3)



Biot-Savart law

(a) along circle

$$\vec{B}_1 = \frac{\mu_0 i}{4\pi R} \int d\vec{r} \times \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3} ; \vec{r} = 0$$

$C_1: \vec{r}' = R \vec{i}_y$ with ϕ (measured as in the figure!)
 $\phi \in (0, 2\pi - \theta)$

$$d\vec{r}' = R \vec{i}_\phi \frac{d\phi}{2\pi - \theta}$$

$$\vec{B}_1 = -\frac{\mu_0 i}{4\pi} \int_0^{2\pi - \theta} d\phi R \vec{i}_\phi \times \frac{R \vec{i}_y}{R^3}$$

$$= -\frac{\mu_0 i}{4\pi R} (-\vec{i}_z) \int_0^{2\pi - \theta} d\phi = +\frac{\mu_0 i}{4\pi R} (2\pi - \theta) \vec{i}_z$$

(b) along straight line

$\vec{r}' = -R \cos(\frac{\theta}{2}) \vec{i}_y + x \vec{i}_x$ with x from $-R \sin(\frac{\theta}{2}) \rightarrow R \sin(\frac{\theta}{2})$

$$d\vec{r}' = dx \vec{i}_x$$

$$\vec{B}_2 = \frac{\mu_0 i}{4\pi} \int_{-R \sin(\frac{\theta}{2})}^{R \sin(\frac{\theta}{2})} dx \vec{i}_x \times \frac{R \cos(\frac{\theta}{2}) \vec{i}_y}{[x^2 + R^2 \cos^2(\frac{\theta}{2})]^{3/2}}$$

$$\vec{B}_2 = \frac{\mu_0 i}{4\pi} R \cos\left(\frac{\theta}{2}\right) \vec{i}_z \int_0^{R \sin(\theta/2)} dx' \frac{1}{(x'^2 + R^2 \cos^2(\frac{\theta}{2}))^{3/2}}$$

$$= \frac{\mu_0 i}{2\pi} \frac{R^2 \sin(\frac{\theta}{2})}{R^3 \cos(\frac{\theta}{2})} \vec{i}_z$$

$$\vec{B}_2 = \frac{\mu_0 i \tan(\frac{\theta}{2})}{2\pi R} \vec{i}_z$$

$$\vec{B} = \left[\frac{\mu_0 i}{4\pi R} (2\pi - \theta) + \frac{\mu_0 i \tan(\frac{\theta}{2})}{2\pi R} \right] \vec{i}_z$$

$$(4) \vec{B} = -\vec{i}_z (\alpha + \beta x)$$

$$(a) \Phi_{\vec{B}} = - \int_0^H dy \int_0^D dx (\alpha + \beta x)$$

$$= -H \left(\alpha D + \frac{\beta}{2} D^2 \right)$$

$$(b) \oint_{\partial S} d\vec{r} \cdot \vec{E} = +V_c = -\dot{\Phi}_B - L \frac{di}{dt} = + \frac{Q}{C}$$

$$-\dot{\Phi}_B = H C_1 D$$

$$\frac{Q}{C} = -L \dot{Q} + H C_1 D$$

$$(c) \text{ For } L=0 \Rightarrow Q = C H C_1 D$$

There is a constant EMF induced and thus the charge is constant or same

For the forced case ($L \neq 0$) we have to add the solution of the homogeneous equation:

$$\ddot{Q}_H = -\frac{Q_H}{LC} \Rightarrow Q_H = a \cos(\omega t + \phi_0)$$

$$\text{with } \omega = \frac{1}{\sqrt{LC}}$$

If for $t=0$; $Q=0$ (initial condition) we have

$$\phi_0 = 0; a = -CH_0 D$$

$$Q(t) = CH_0 D [1 - \cos(\omega t)]$$

(5)