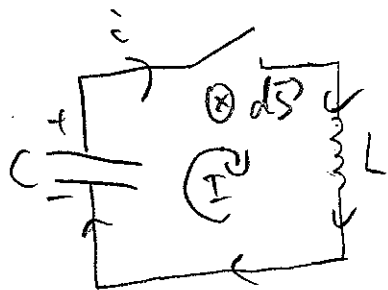


The LC circuit

①

(a) The simple oscillator



Suppose the capacitor is charged with a certain charge Q_0 at $t=0$, and there is no current flowing

Now we close the switch at $t=0$. To analyze this problem we assume a current remaining as in the circuit and use loop \textcircled{I} & Faraday's law. The surface vector $d\vec{S}$ must point into the page (RHR!) The magnetic field induced by the current also points in. Thus we have

$$\vec{\Phi}_B = + Li$$

$$\Rightarrow \oint_{dS} d\vec{r} \cdot \vec{E} = -\frac{Q}{C} = -\frac{d}{dt}(Li) = -L \frac{di}{dt}$$

Since i takes charge from the + plate of the capacitor we have:

$$i = -\frac{dQ}{dt} \Rightarrow \frac{di}{dt} = -\frac{d^2Q}{dt^2}$$

$$\Rightarrow \frac{d^2Q}{dt^2} = -\frac{1}{LC} Q = -\omega_0^2 Q$$

$$\omega_0 = \frac{1}{\sqrt{LC}} \quad (\text{Thomson formula})$$

The general solution for the 2nd-order ode is

$$Q(t) = A \cos(\omega_0 t) + B \sin(\omega_0 t)$$

The constants A and B are determined by the initial conditions:

$$Q(0) = Q_0$$

$$\dot{Q}(0) = 0$$

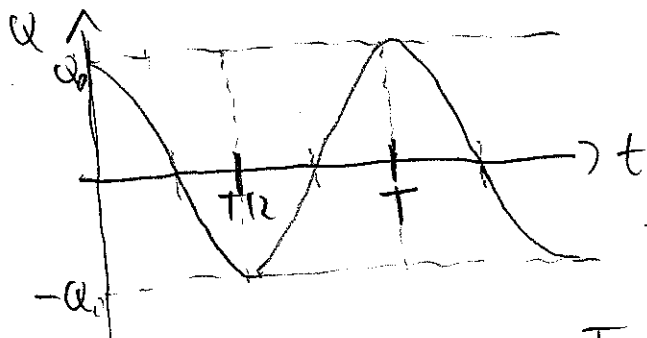
$$\Rightarrow Q(0) = A = Q_0 ; \dot{Q}(0) = -\frac{dQ}{dt} \Big|_{t=0} = -B\omega_0 = 0$$

$$\Rightarrow B = 0$$

$$\Rightarrow Q(t) = Q_0 \cos(\omega_0 t)$$

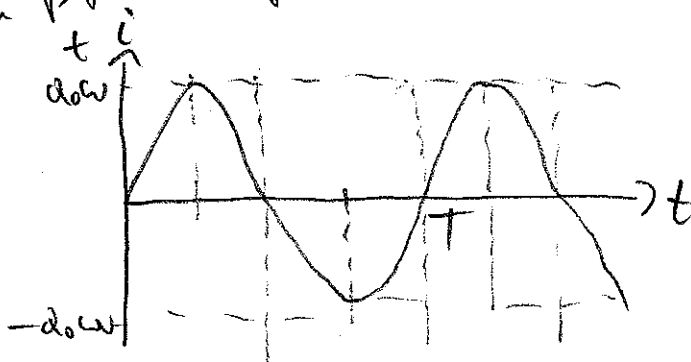
i.e. the charge oscillates back and forth. The current is

$$i(t) = -\frac{dQ}{dt} = Q_0 \omega_0 \sin(\omega_0 t)$$



$$T : \omega T = 2\pi \Rightarrow T = \frac{2\pi}{\omega_0} = 2\pi\sqrt{LC}$$

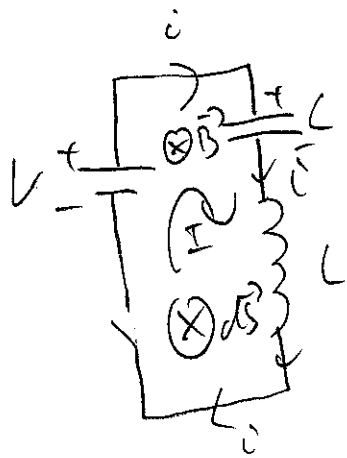
is the period of the oscillations



The current has of course the same period.

(b) Battery in series

(3)



In the same way as before we get

$$-V + \frac{Q}{C} = -L \frac{di}{dt}$$

$$\text{but now } i = + \frac{dQ}{dt}$$

$$\Rightarrow L \frac{d^2 Q}{dt^2} + \frac{Q}{C} = V$$

$$\Rightarrow \frac{d^2 Q}{dt^2} + \frac{Q}{LC} = \frac{V}{L}$$

We only need a particular solution for the in-hom. equation, since we solved the homogeneous one already above. "ANSWER of the right-hand side"

$$Q = A_1 = \text{const} \Rightarrow \frac{A_1}{LC} = \frac{V}{L} \Rightarrow A_1 = CV$$

The general solution thus is (with $\omega = \frac{1}{\sqrt{LC}}$ as given)

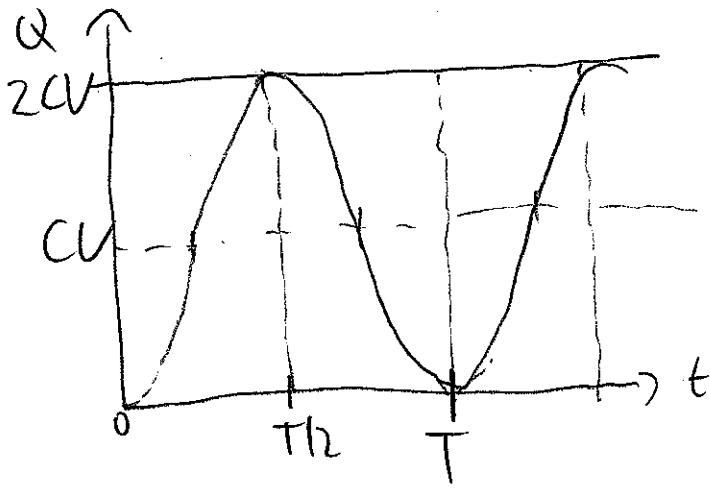
$$Q(t) = CV + A \cos(\omega t) + B \sin(\omega t)$$

Initial conditions: $t=0: Q(0)=0, \dot{Q}(0)=0$

$$\Rightarrow \begin{aligned} CV + A &= 0 \\ B\omega &= 0 \end{aligned} \Rightarrow A = -CV, B=0$$

$$Q(t) = CV [1 - \cos(\omega t)]$$

So we have a peculiar behavior like this



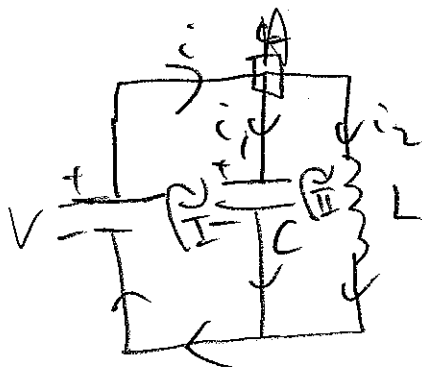
The current is

$$i(t) = \frac{dQ}{dt} = CV \omega \sin(\omega t)$$

as above.

It's of course unphysical since in reality we always have a finite resistance, and the oscillations are damped away such that you end with a steady charge from a unit with the capacitor "fully charged" to $Q_{t \rightarrow \infty} = CV$.

(c) Battery in parallel



in initial condition

$$Q = \dot{Q} = i_1 = i_2 = 0$$

$$I: -V + \frac{Q}{C} = 0 \quad (\text{no } L \text{ in this loop})$$

$$II: -\frac{Q}{C} = -L \frac{di_2}{dt}$$

$$III: \dot{Q} = i_1 + i_2$$

(5)

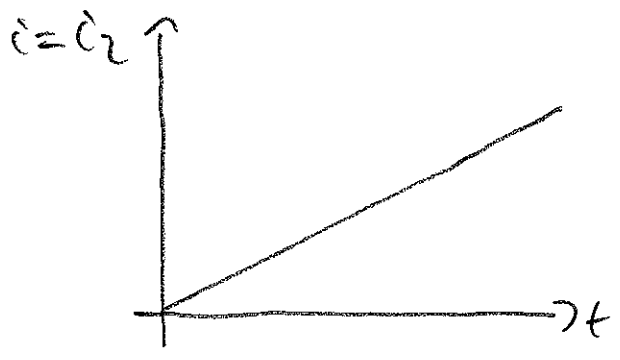
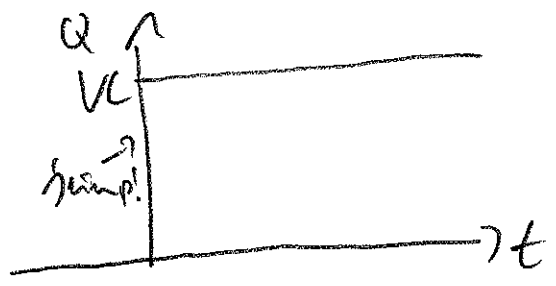
$$\dot{q}_1 = \frac{dq}{dt} = 0 \text{ due to (I)!} \Rightarrow Q = \text{const.} = VC$$

$$\Rightarrow V = L \frac{di_2}{dt} \Rightarrow i_2 = \frac{V}{L} t + A$$

$$i_2(0) = 0 \Rightarrow i_2 = \frac{V}{L} t$$

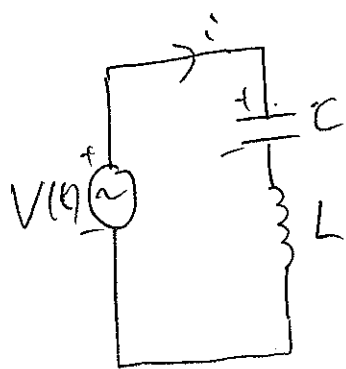
$$i = i_2$$

Thus we find



The jump comes from the neglect of the self-inductance of loop ①. The rise of $i_2 \rightarrow \infty$ for $t \rightarrow \infty$ can also never happen, because of Kirchhoff's law. We shall discuss the more realistic case later on this lecture.

cd) AC Voltage (Series)



$$V(t) = V_0 \cos(\omega t)$$

We only need to plug $V(t)$ on the equation instead of the constant V :

$$\frac{d^2 q}{dt^2} + \frac{q}{LC} = \frac{V_0}{L} \cos(\omega t)$$

and we only need a new particular solution for the inhomogeneous equation. Here it's clear that

$$q(t) = A_2 \cos(\omega t)$$

should work:

$$[-A_2 \omega^2 + A_2 \omega_0^2] \cos(\omega t) = \frac{V_0}{L} \cos(\omega t)$$

$$\Rightarrow \Rightarrow A_2 = \frac{V_0}{L} \frac{1}{\omega_0^2 - \omega^2}$$

This works only for $\omega \neq \omega_0$. We'll see later what to do in this case

The general solution of the full eq. is thus

$$Q(t) = A \cos(\omega_0 t) + B \sin(\omega_0 t) + \frac{V_0}{L} \frac{1}{\omega_0^2 - \omega^2} \cos(\omega t)$$

Initial conditions

$$Q(0) = A + \frac{V_0}{L} \frac{1}{\omega_0^2 - \omega^2} \stackrel{!}{=} 0 \Rightarrow A = -\frac{V_0}{L} \frac{1}{\omega_0^2 - \omega^2}$$

$$\dot{Q}(0) = \omega_0 B = 0 \Rightarrow B = 0$$

$$Q(t) = \frac{V_0}{L} \frac{1}{\omega_0^2 - \omega^2} [\cos(\omega t) - \cos(\omega_0 t)]$$

The current is

$$i(t) = \frac{dQ}{dt} = -\frac{V_0}{L} \frac{1}{\omega_0^2 - \omega^2} [\omega \sin(\omega t) - \omega_0 \sin(\omega_0 t)]$$

The resonance case ($\omega = \omega_0$)

For the same initial conditions we can take the limit $\omega \rightarrow \omega_0$

$$Q(t) = \lim_{\omega \rightarrow \omega_0} \frac{V_0}{L} \frac{1}{\omega_0^2 - \omega^2} [\cos(\omega t) - \cos(\omega_0 t)]$$

de L'Hospital's formula for $\frac{0}{0}$ limit

$$\stackrel{!}{=} \lim_{\omega \rightarrow \omega_0} \frac{V_0}{L} \frac{1}{+\omega} [-t \sin(\omega t)] \text{ in } \omega_0 t$$

$$\text{so } \frac{dQ}{dt} = \frac{V_0}{2\omega_0 L} t \sin(\omega_0 t)$$

$$i(t) = \dot{Q}(t) = \frac{V_0}{2\omega_0 L} [\omega_0 t \cos(\omega_0 t) + \sin(\omega_0 t)]$$

$$\dot{Q}(t) = \frac{V_0}{2L}$$

Check whether this solves the equation:

(8)

$$\ddot{Q}(t) = \frac{V_0}{2\omega_0 L} \left[2\omega_0 \cos(\omega_0 t) - \omega_0^2 t \sin(\omega_0 t) \right]$$

$$\ddot{Q} + \omega_0^2 Q = \frac{V_0}{L} \cos(\omega_0 t) \quad (2)$$

Both, current and charge on the capacitor oscillate with an increasingly growing amplitude ("resonance catastrophe" / "ply"). In reality this never happens, because of a finite resistance

(2) AC (parallel)

$$\frac{di_2}{dt} = \frac{V_0}{L} \cos(\omega t) \Rightarrow i_2 = \frac{V_0 \omega}{\omega L} \sin(\omega t) + A$$

$$i_2(\omega) = 0 \Rightarrow A = 0 \Rightarrow i_2 = \frac{V_0}{\omega L} \sin(\omega t)$$

$$Q = CV_0 \cos(\omega t)$$

$$i_1 = \frac{dQ}{dt} = -V_0 \omega C \sin(\omega t)$$

$$i = i_1 + i_2 = V_0 \left(\frac{1}{\omega L} - \omega C \right) \sin(\omega t)$$

$$i = \frac{V_0}{\omega L} (1 - \omega^2 LC) \sin(\omega t)$$

$$i = \frac{V_0}{\omega} \left(1 - \frac{\omega^2}{\omega_0^2} \right) \sin(\omega t)$$

For $\omega \rightarrow \omega_0$ we get

$$i = 0$$

$$i_2 = -i_1 = \frac{V_0}{\omega_0 L} \sin(\omega_0 t)$$

$$Q = CV_0 \cos(\omega_0 t)$$