

asymptotic safety beyond the Standard Model

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Frankfurt, 18 May 2017

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standard model

local QFT for fundamental interactions

strong nuclear force

weak force

electromagnetic force

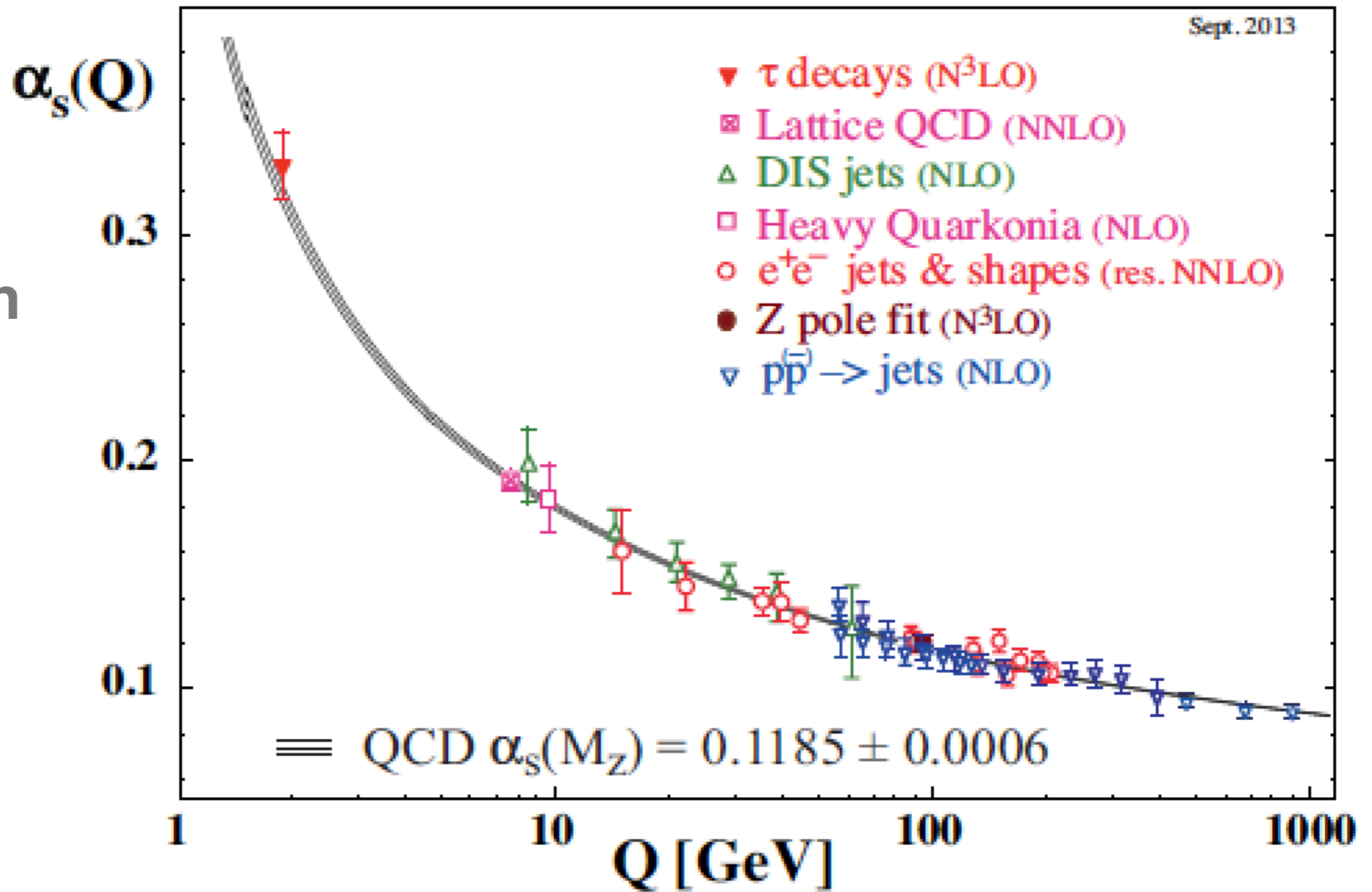
open challenges

what comes **beyond the SM**?

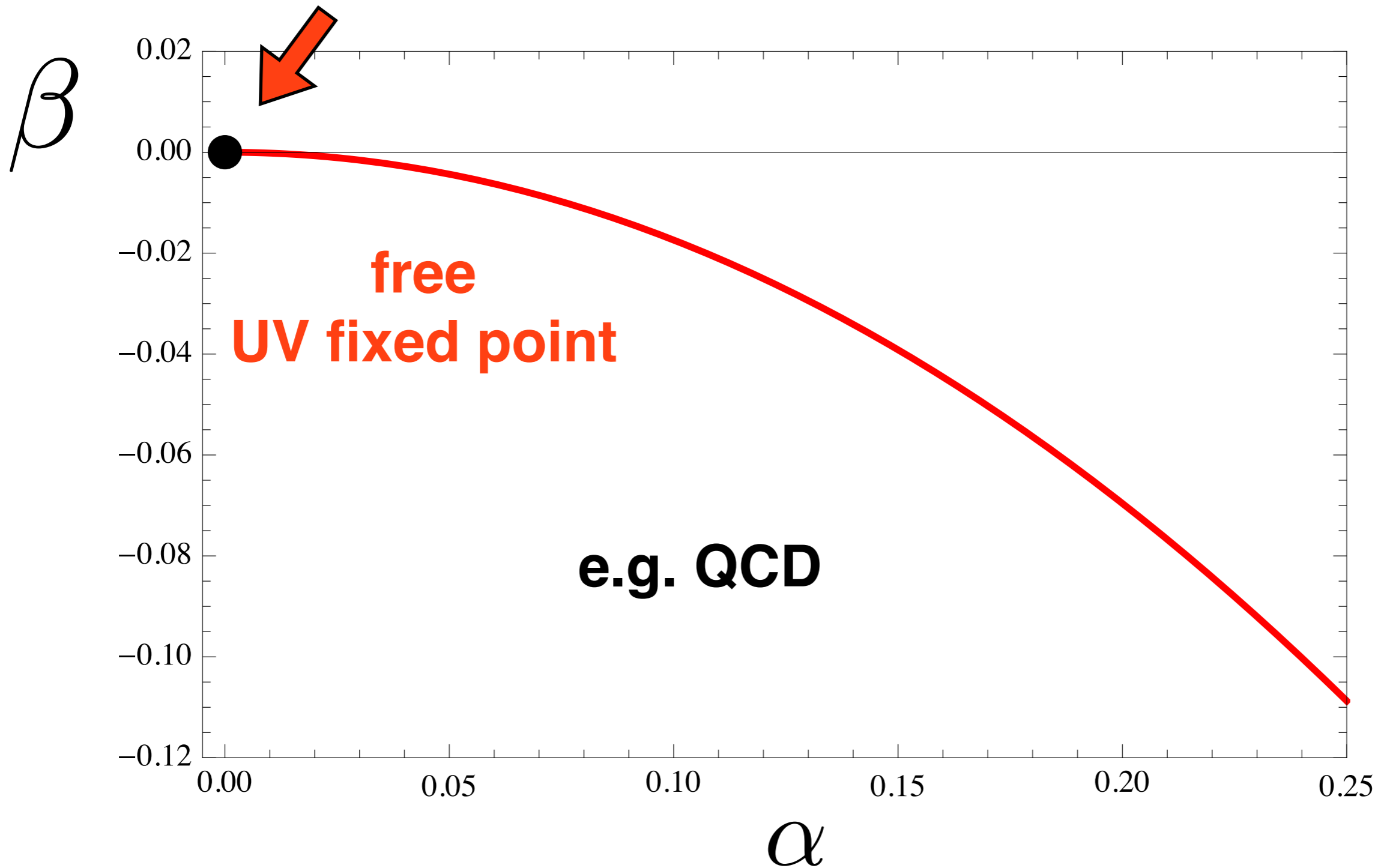
how does **gravity** fit in?

asymptotic freedom

triumph
of QFT



asymptotic freedom



conditions for asymptotic freedom

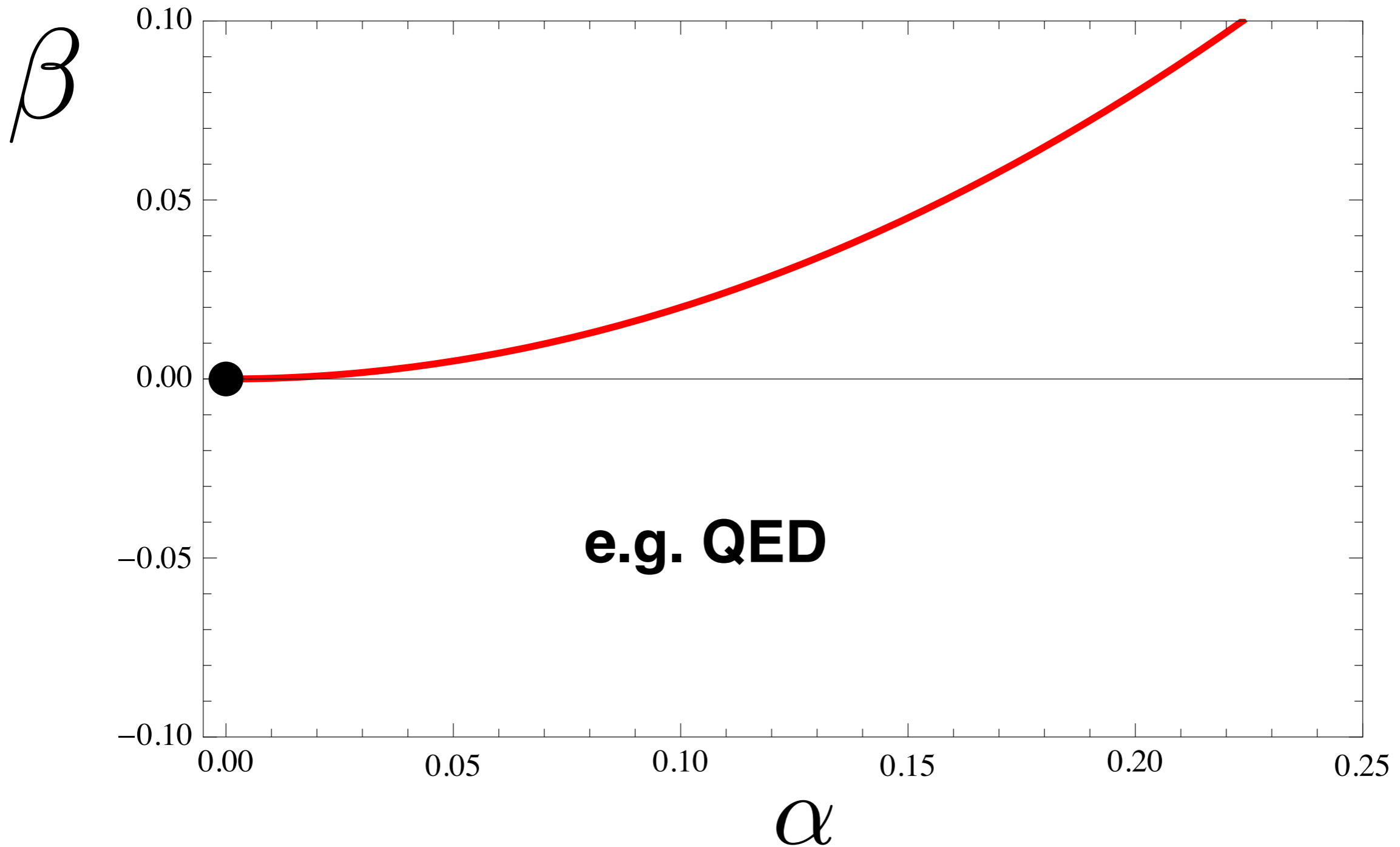
complete asymptotic freedom

couplings achieve **non-interacting** UV fixed point

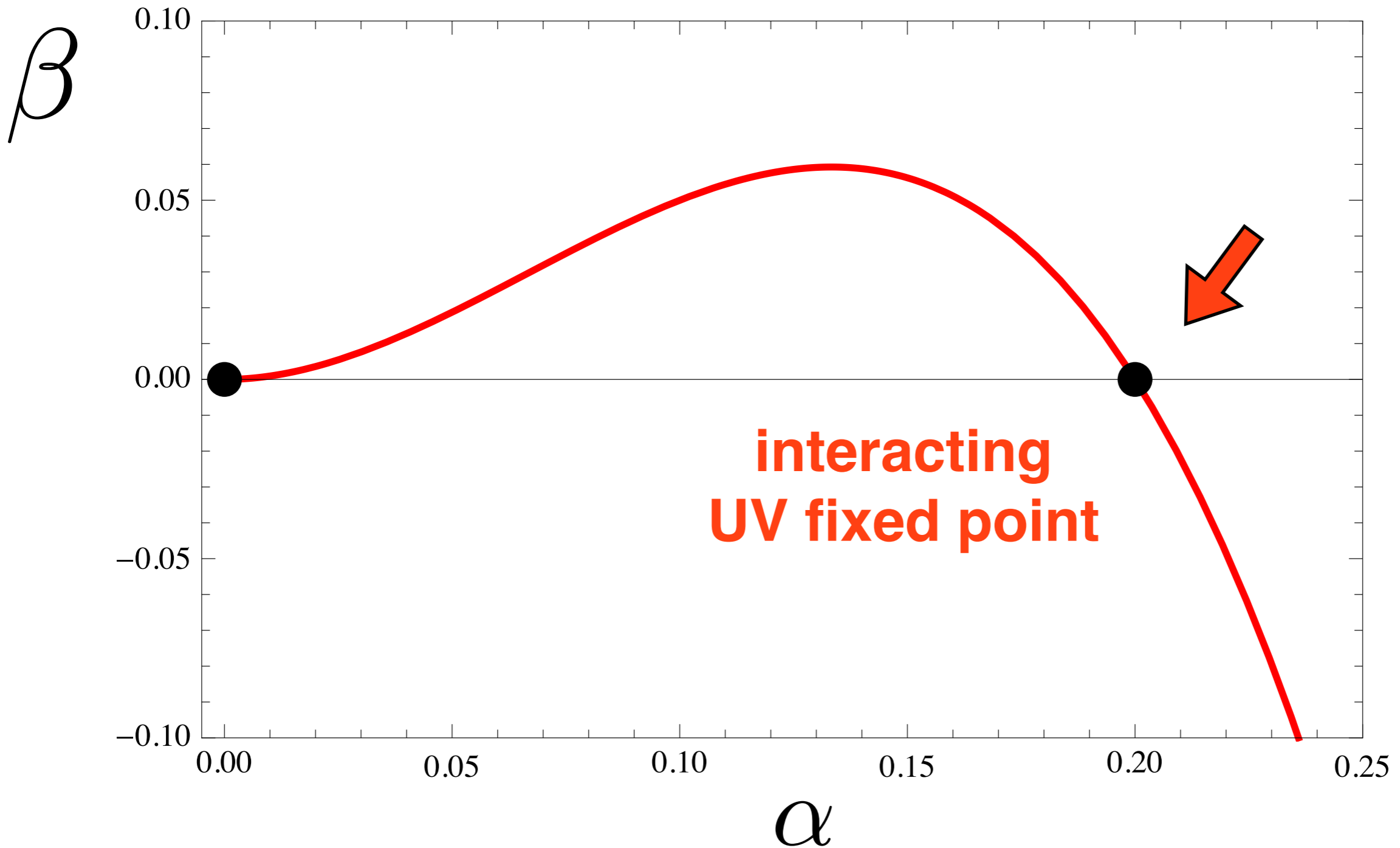
fields	cAF
scalars	no
scalars with fermions	no
U(1) w/o scalars or fermions	no
non-Abelian gauge fields	yes
non-Abelian fields with fermions	yes *)
non-Abelian fields, fermions, scalars	yes *)

*) provided certain auxiliary conditions hold true

infrared freedom



asymptotic safety



asymptotic safety

idea:

some or all couplings achieve

interacting UV fixed point

Wilson '71
Weinberg '79

if so, **new directions** for

BSM physics &, possibly, quantum gravity

proof of existence:

4D gauge-Yukawa theory with

exact asymptotic safety

Litim, Sannino, 1406.2337

Bond, Litim @ERG2016

asymptotic safety

today:

1. **theorems** for asymptotic safety

Bond, Litim 1608.00519

2. weakly **interacting UV completions**
of the Standard Model

3. **constraints** from data (colliders)

AD Bond, G Hiller, K Kowalska, DF Litim, 1702.01727

asymptotic safety

today:

1. **theorems** for asymptotic safety

Bond, Litim 1608.00519

conditions for asymptotic safety

results

Bond, Litim | 608.00519

case	gauge group	matter	Yukawa	asymptotic safety
a)	simple	fermions in irreps	No	No
b)	simple or abelian	fermions, any rep	No	No
		scalars, any rep	No	No
		fermions and scalars, any rep	No	No
c)	semi-simple, with or without abelian factors	fermions, any rep	No	No
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basics of asymptotic safety

gauge theory

$$\partial_t \alpha_g = -B \alpha_g^2 + C \alpha_g^3$$

$$t = \ln \mu / \Lambda$$

$$0 < \alpha^* = B/C \ll 1$$

$$\alpha_* \ll 1$$

loop coefficients

$$B > 0 \quad \text{asymptotic freedom} \quad C < 0 \text{ or } C > 0$$

in the latter case:

$$\alpha_g^* = \frac{B}{C}$$

Banks-Zaks IR FP

basics of asymptotic safety

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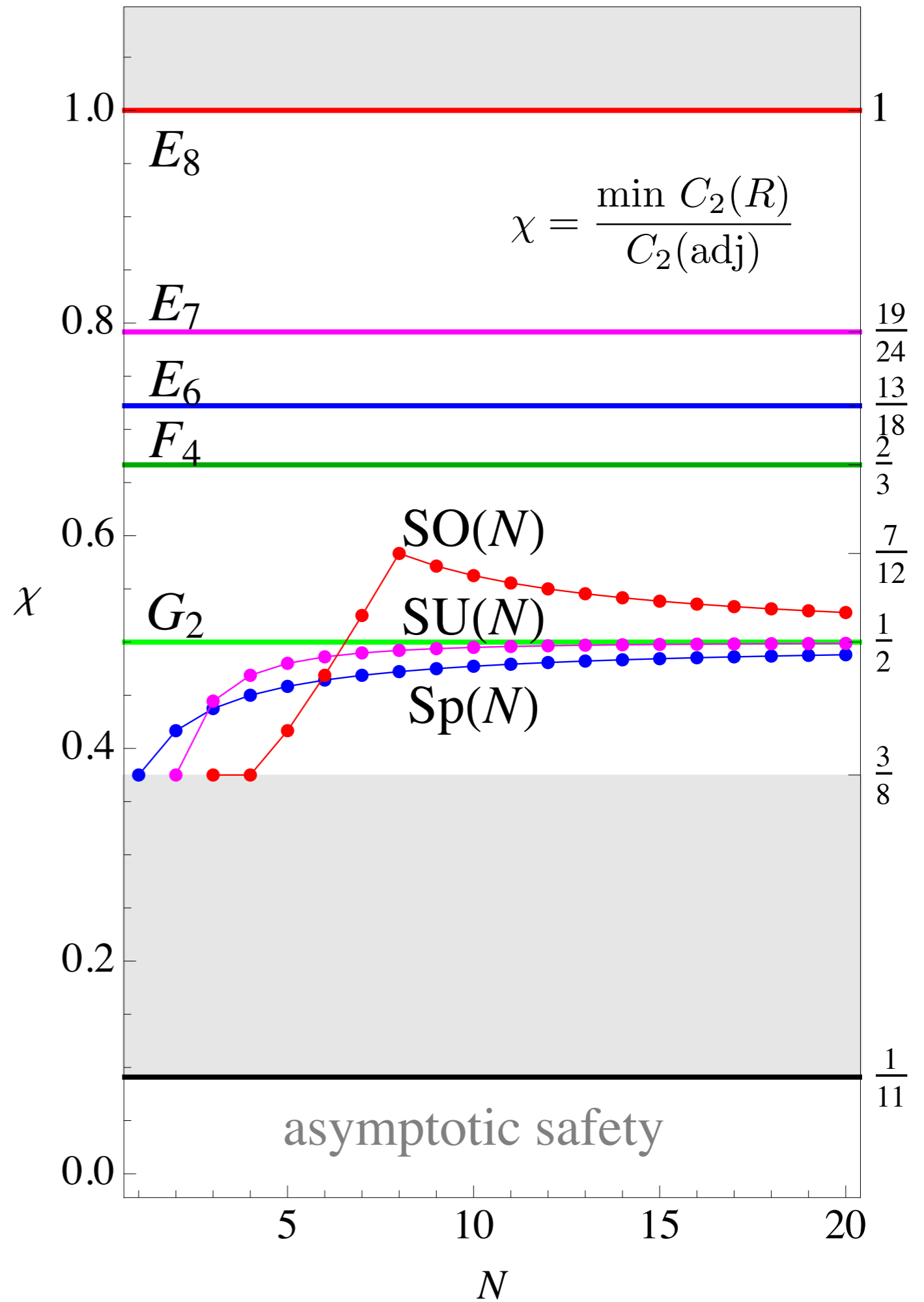
loop coefficients

$$B < 0 \quad \text{infrared freedom}$$

for $C < 0$ we must have

$$C_2^S < \frac{1}{11} C_2^G$$

result:

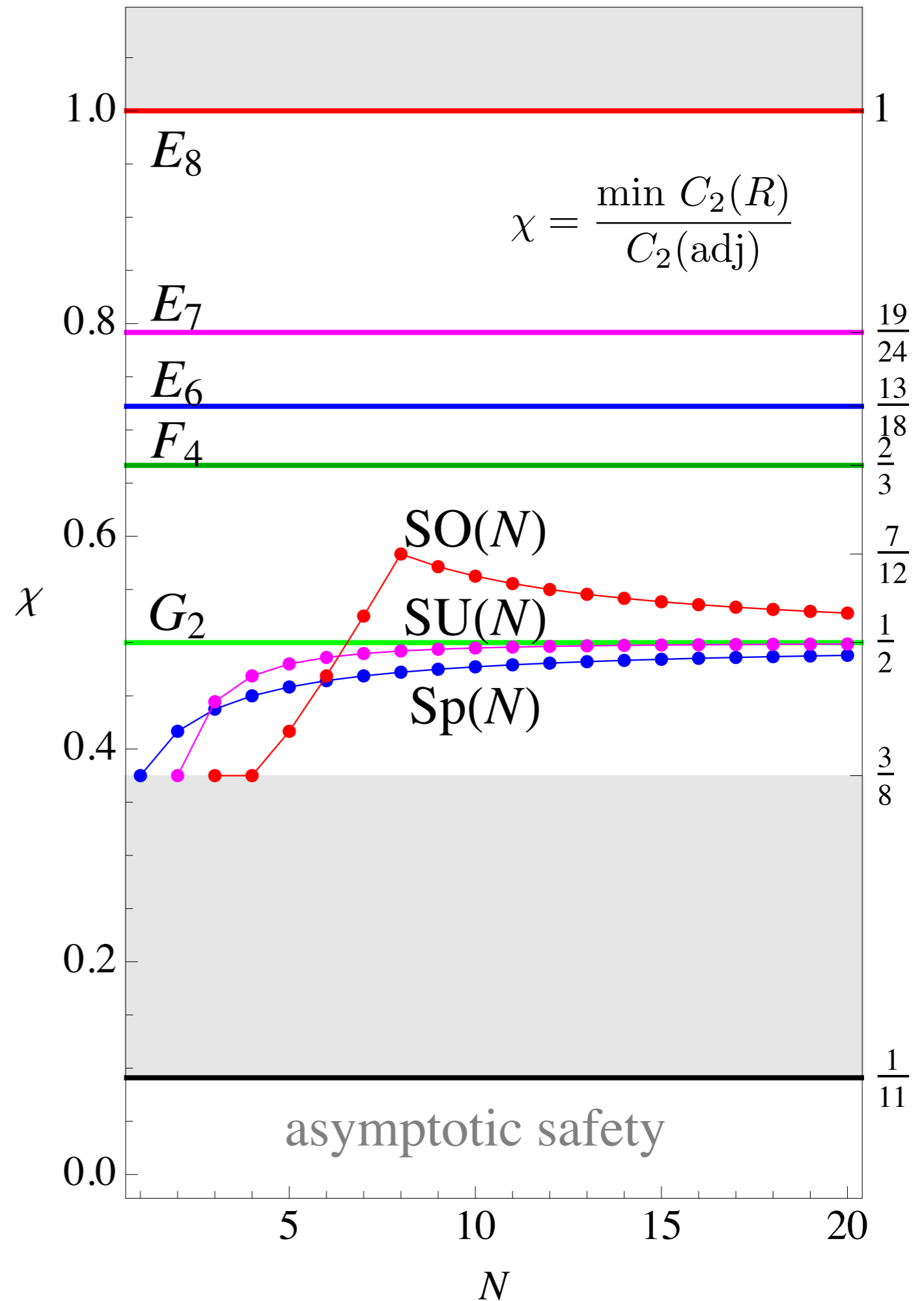


result:

implication:

$$B \leq 0 \quad \Rightarrow \quad C > 0$$

no go theorem



basics of asymptotic safety

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Bond, Litim 1608.00519

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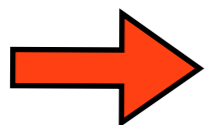
Bond, Litim 1608.00519

can other couplings help?

more gauge: **useless**

scalar quartics: **useless**

Yukawas: **unique viable option**



basics of asymptotic safety

gauge Yukawa theory

$$\partial_t \alpha_g = -B \alpha_g^2 + C \alpha_g^3 - D \alpha_g^2 \alpha_y$$

$$\partial_t \alpha_y = E \alpha_y^2 - F \alpha_g \alpha_y$$

$$t = \ln \mu / \Lambda$$

$$\alpha_* \ll 1$$

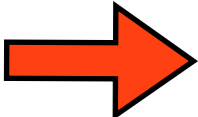
loop coefficients $D, E, F > 0$ in any QFT

basics of asymptotic safety

gauge Yukawa theory

$$\partial_t \alpha_g = -B \alpha_g^2 + C \alpha_g^3 - D \alpha_g^2 \alpha_y$$

$$t = \ln \mu / \Lambda$$


$$\partial_t \alpha_y = E \alpha_y^2 - F \alpha_g \alpha_y$$

$$\alpha_* \ll 1$$

Yukawa nullcline

$$\alpha_y^* = \frac{F}{E} \alpha_g^*$$

basics of asymptotic safety

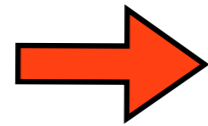
gauge Yukawa theory

$$\partial_t \alpha_g = -B \alpha_g^2 + C \alpha_g^3 - D \alpha_g^2 \alpha_y \quad t = \ln \mu / \Lambda$$

$$\partial_t \alpha_y = E \alpha_y^2 - F \alpha_g \alpha_y \quad \alpha_* \ll 1$$

Yukawa nullcline

$$\alpha_y^* = \frac{F}{E} \alpha_g^*$$



$$\beta_g| = (-B + C' \alpha_g) \alpha_g^2$$

shifted two-loop

$$C \rightarrow C' = C - D \frac{F}{E}$$

interacting UV fixed point iff

$$D F - C E > 0$$

basics of asymptotic safety

gauge Yukawa theory

$$\partial_t \alpha_g = -B \alpha_g^2 + C \alpha_g^3 - D \alpha_g^2 \alpha_y$$

$$t = \ln \mu / \Lambda$$

$$\partial_t \alpha_y = E \alpha_y^2 - F \alpha_g \alpha_y$$

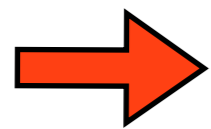
$$\alpha_* \ll 1$$

Yukawa nullcline

$$\alpha_y^* = \frac{F}{E} \alpha_g^*$$

$$\beta_g| = (-B + C' \alpha_g) \alpha_g^2$$

gauge-Yukawa fixed point



$$(\alpha_g^*, \alpha_y^*) = \left(\frac{B}{C'}, \frac{B}{C'} \frac{F}{E} \right)$$

UV or **IR**

basics of asymptotic safety

gauge Yukawa theory

$$\partial_t \alpha_g = -B \alpha_g^2 + C \alpha_g^3 - D \alpha_g^2 \alpha_y$$

$$t = \ln \mu / \Lambda$$

$$\partial_t \alpha_y = E \alpha_y^2 - F \alpha_g \alpha_y$$

$$\alpha_* \ll 1$$

summary of fixed points

$$(\alpha_g^*, \alpha_y^*) = (0, 0)$$

Gaussian

UV or IR

$$(\alpha_g^*, \alpha_y^*) = \left(\frac{B}{C}, 0 \right)$$

Banks-Zaks

IR

$$(\alpha_g^*, \alpha_y^*) = \left(\frac{B}{C'}, \frac{B}{C'} \frac{F}{E} \right)$$

gauge-Yukawa

UV or IR

basics of asymptotic safety

gauge Yukawa theory

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$$\alpha_* \ll 1$$

exact proofs of existence (Veneziano limit)

SU(N) + scalars + fermions

DF Litim, F Sannino, 1406.2337

SU(N) x SU(M) + scalars + fermions

AD Bond, DF Litim, @ERG2016 (to appear)

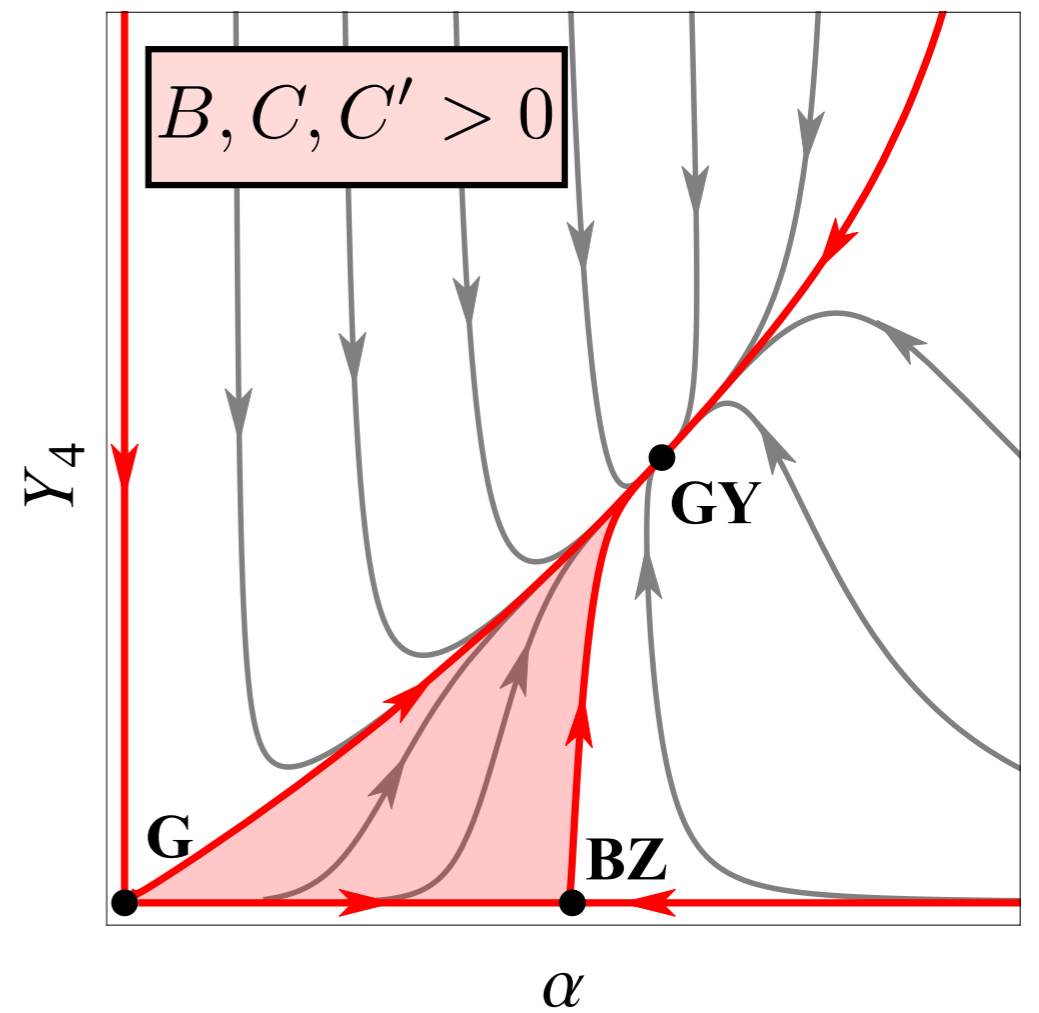
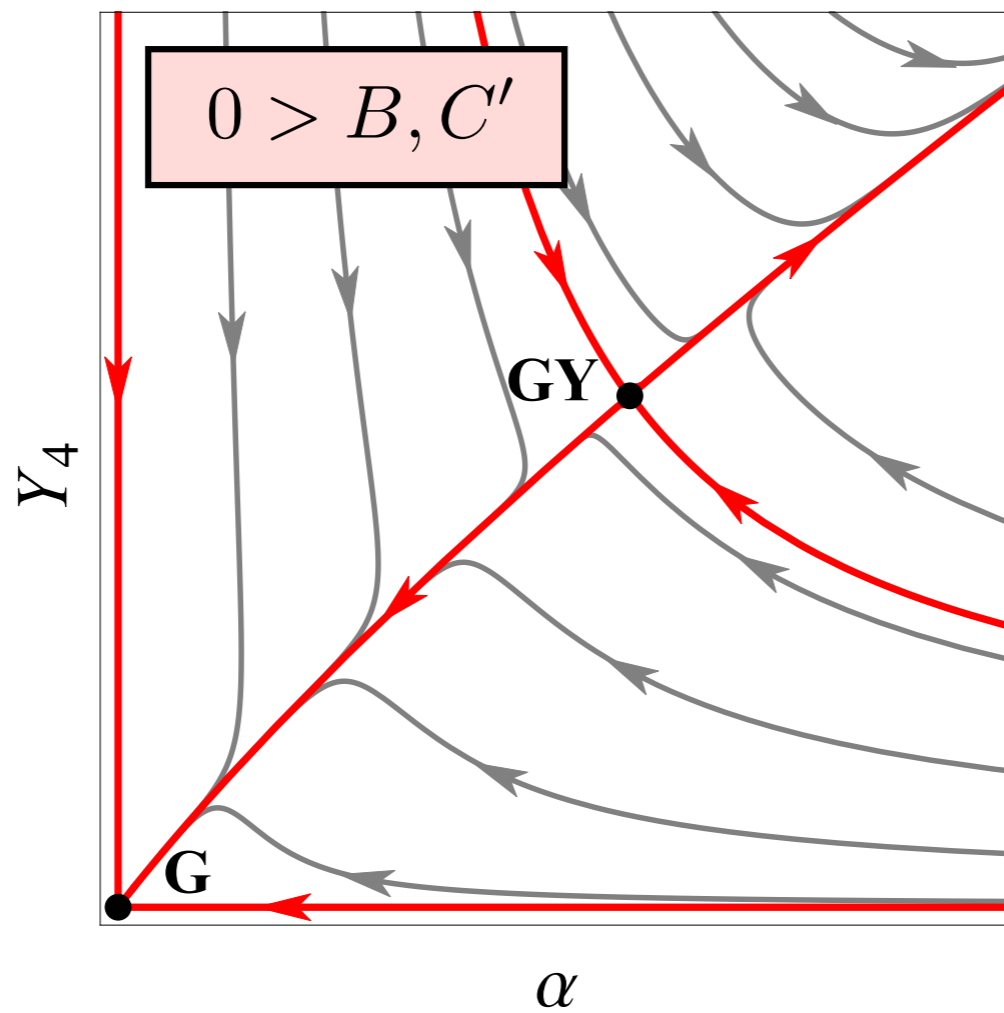
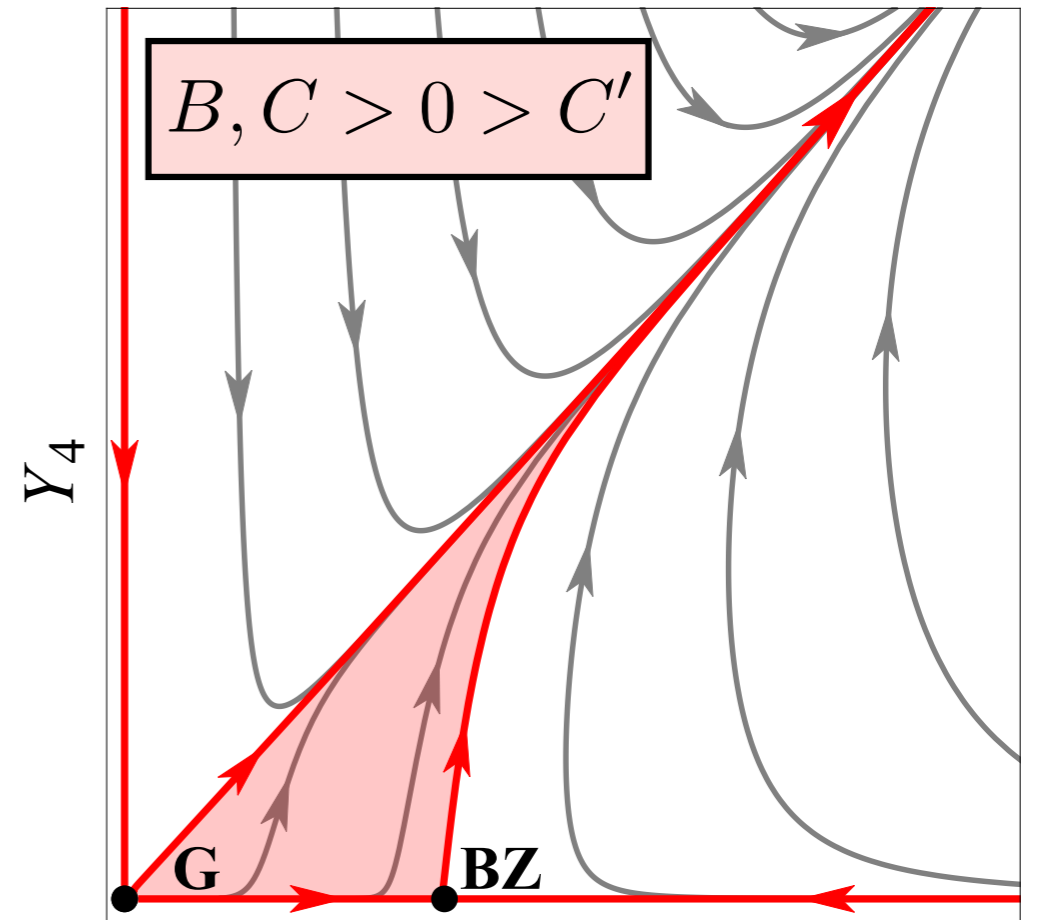
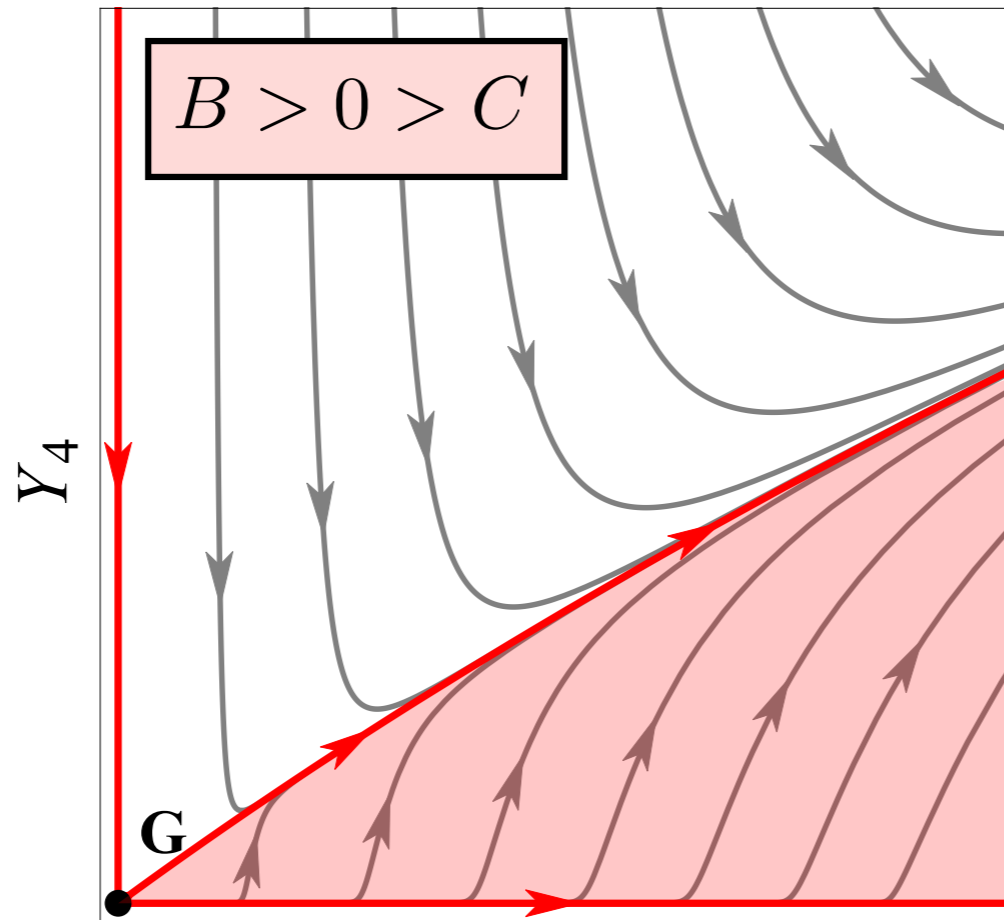
conditions for asymptotic safety

results

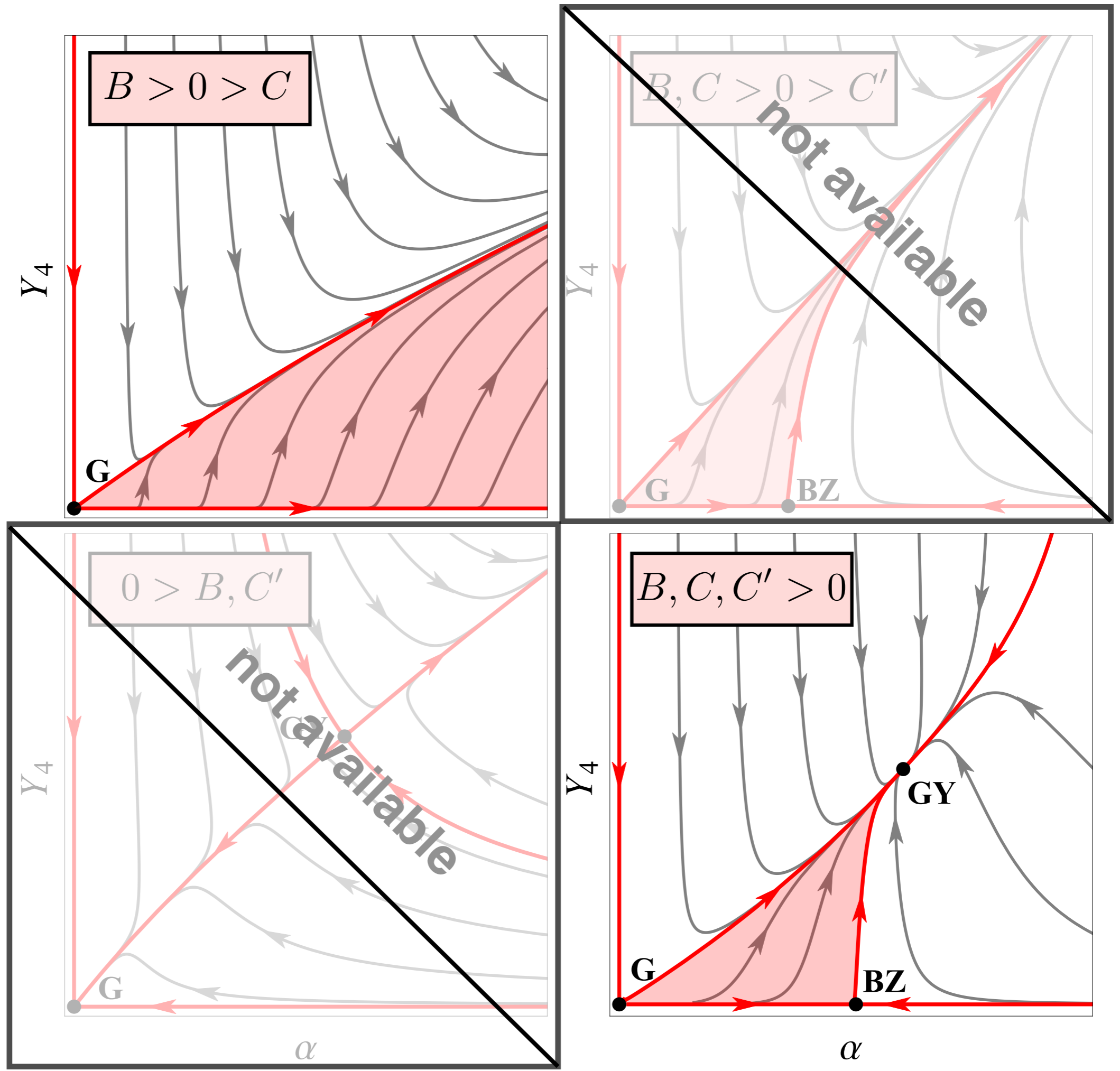
Bond, Litim | 608.00519

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^{*)} provided certain auxiliary conditions hold true



**N=1
SUSY**



asymptotic safety

2. weakly **interacting UV completions** of the Standard Model

AD Bond, G Hiller, K Kowalska, DF Litim, 1702.01727

asymptotic safety beyond the SM

Bond, Hiller, Kowalska, Litim, 1702.01727

N_F **flavors of BSM fermions** $\psi_i(R_3, R_2, Y)$

BSM singlet scalars S_{ij}

global flavor symmetry $U(N_F) \times U(N_F)$

$$L_{\text{BSM, Yukawa}} = -y \text{Tr}(\bar{\psi}_L S \psi_R + \bar{\psi}_R S^\dagger \psi_L)$$

BSM Lagrangean

$$L = L_{\text{SM}} + L_{\text{BSM, kin.}} + L_{\text{BSM, pot.}} + L_{\text{BSM, Yukawa}}$$

UV fixed points

#	gauge couplings		BSM Yukawa	type & info	
FP₁	$\alpha_3^* = 0$	$\alpha_2^* = 0$	$\alpha_y^* = 0$	G · G	non-interacting
FP₂	$\alpha_3^* = 0$	$\alpha_2^* > 0$	$\alpha_y^* > 0$	G · GY	partially interacting
FP₃	$\alpha_3^* > 0$	$\alpha_2^* = 0$	$\alpha_y^* > 0$	GY · G	partially interacting
FP₄	$\alpha_3^* > 0$	$\alpha_2^* > 0$	$\alpha_y^* > 0$	GY · GY	fully interacting

BSM fixed points

FP₂ $\alpha_2^* > 0$ **weak** becomes **strong**
 $\alpha_3^* = 0$ **strong** becomes **weak**

UV critical surface $\delta\alpha_2(\Lambda), \delta\alpha_3(\Lambda)$

FP₃ $\alpha_3^* > 0$ **strong** remains **strong**
 $\alpha_2^* = 0$ **weak** remains **weak**

UV critical surface $\delta\alpha_2(\Lambda), \delta\alpha_3(\Lambda)$

FP₄ $\frac{\alpha_2^*}{\alpha_3^*} \rightarrow \frac{3}{2}$ **weak** becomes the
new strong

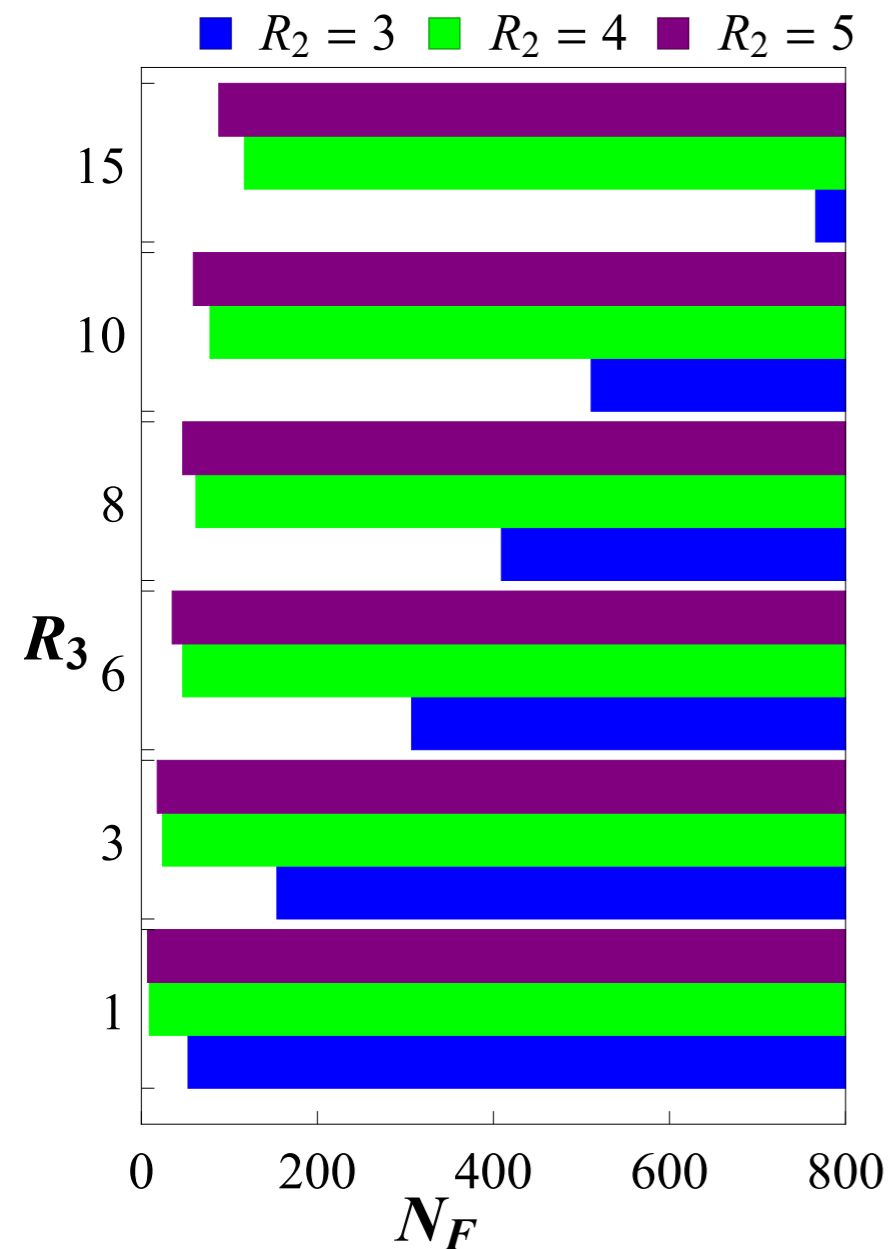
UV critical surface $\delta\alpha_3(\Lambda)$

BSM fixed points

FP₂

$$\alpha_2^* > 0$$

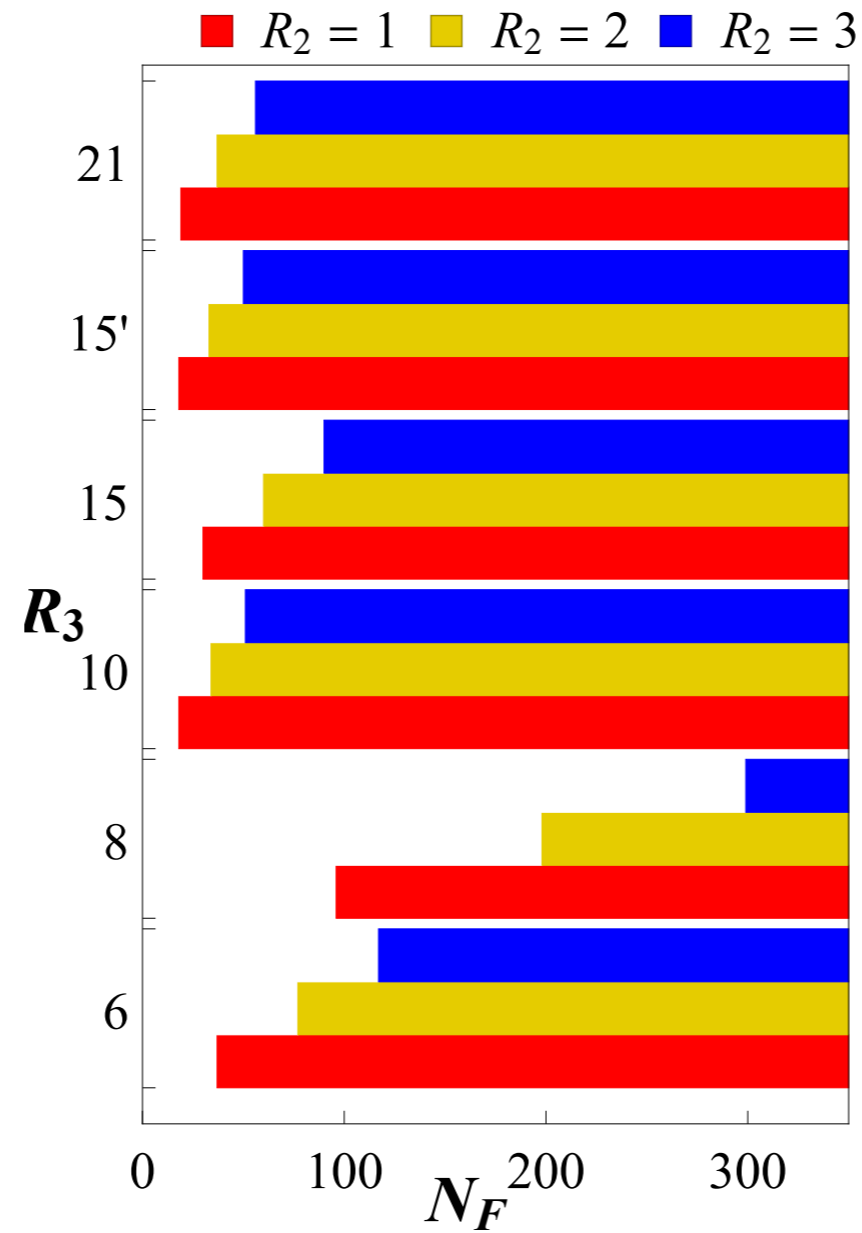
$$\alpha_3^* = 0$$



FP₃

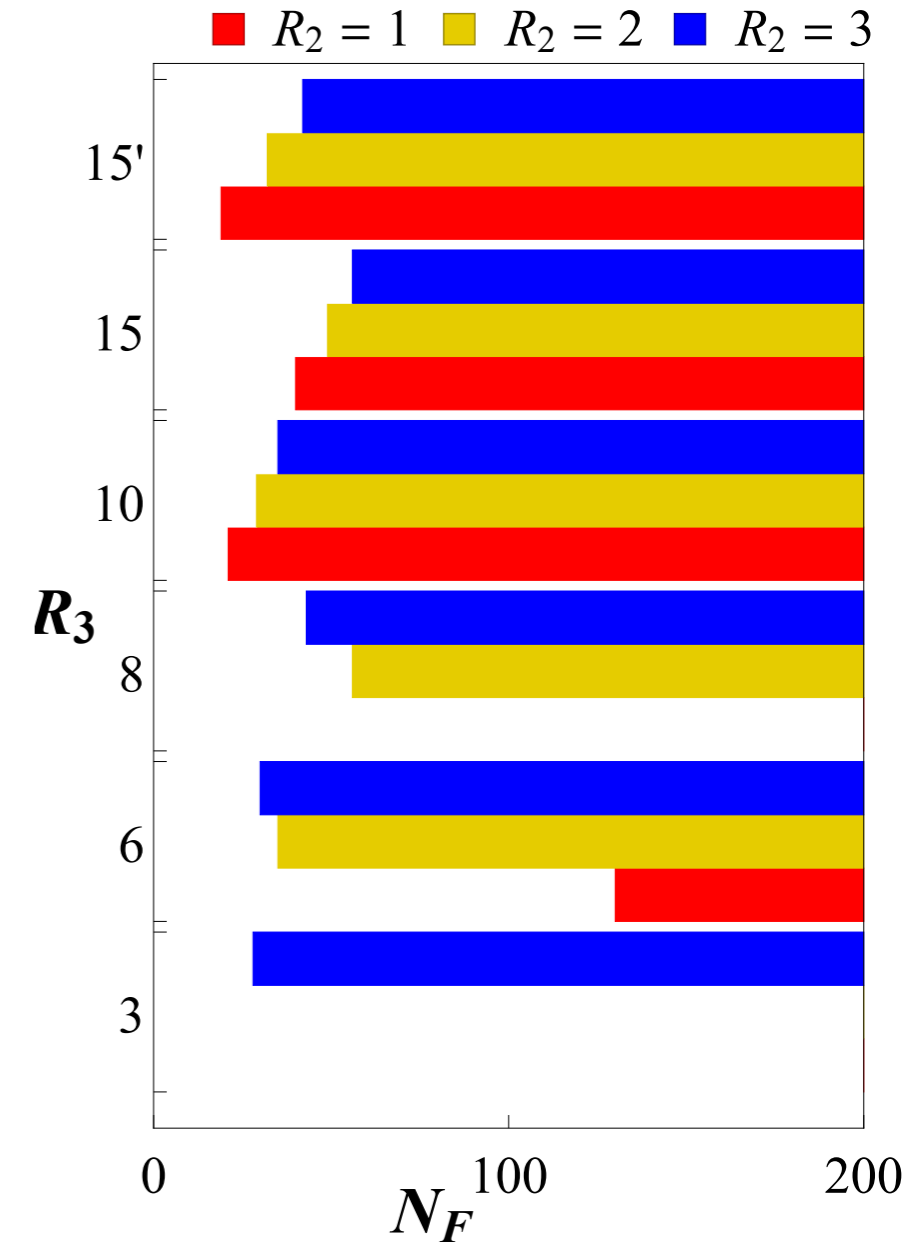
$$\alpha_3^* > 0$$

$$\alpha_2^* = 0$$

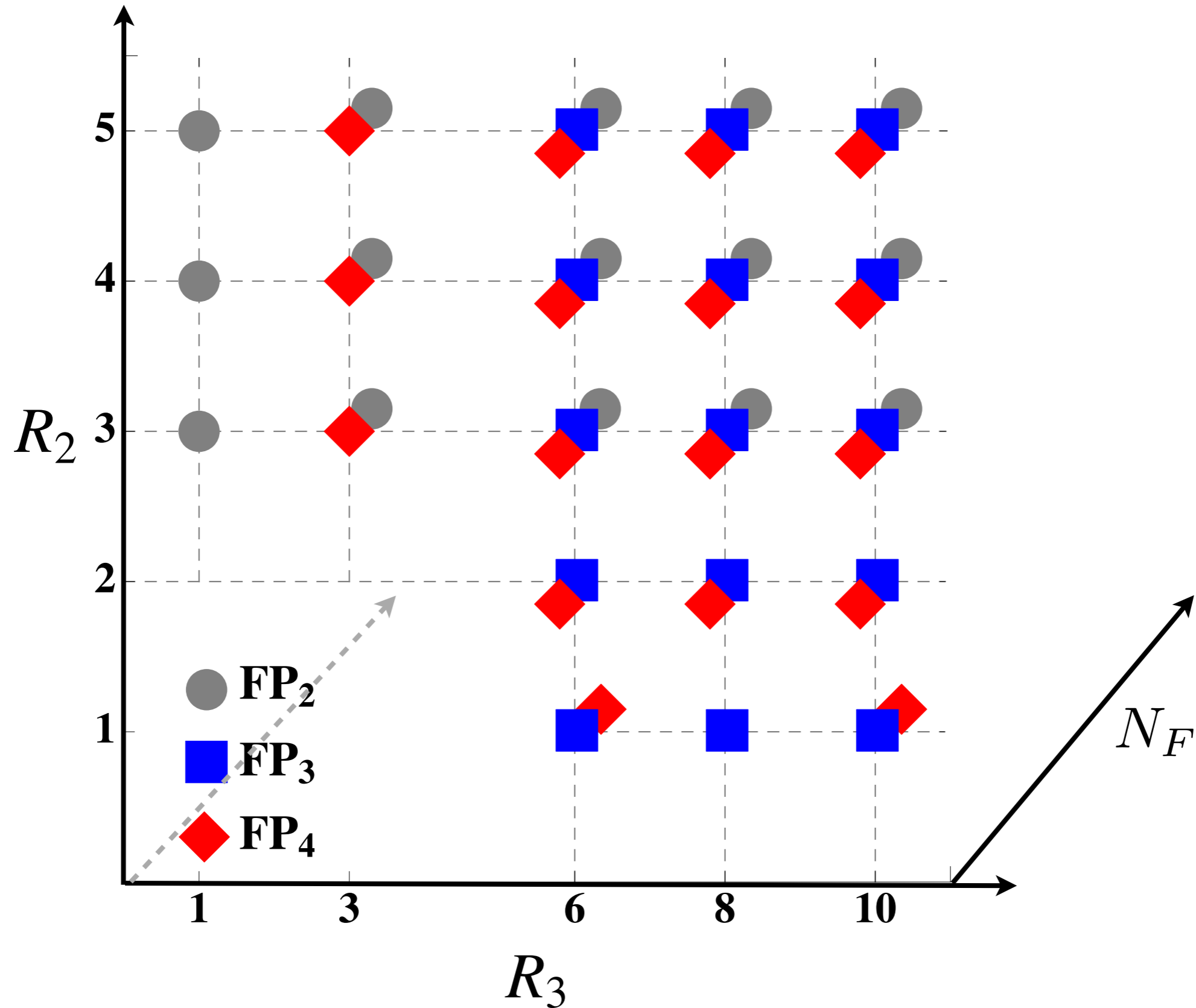


FP₄

$$\alpha_2^*, \alpha_3^* > 0$$



summary of fixed points

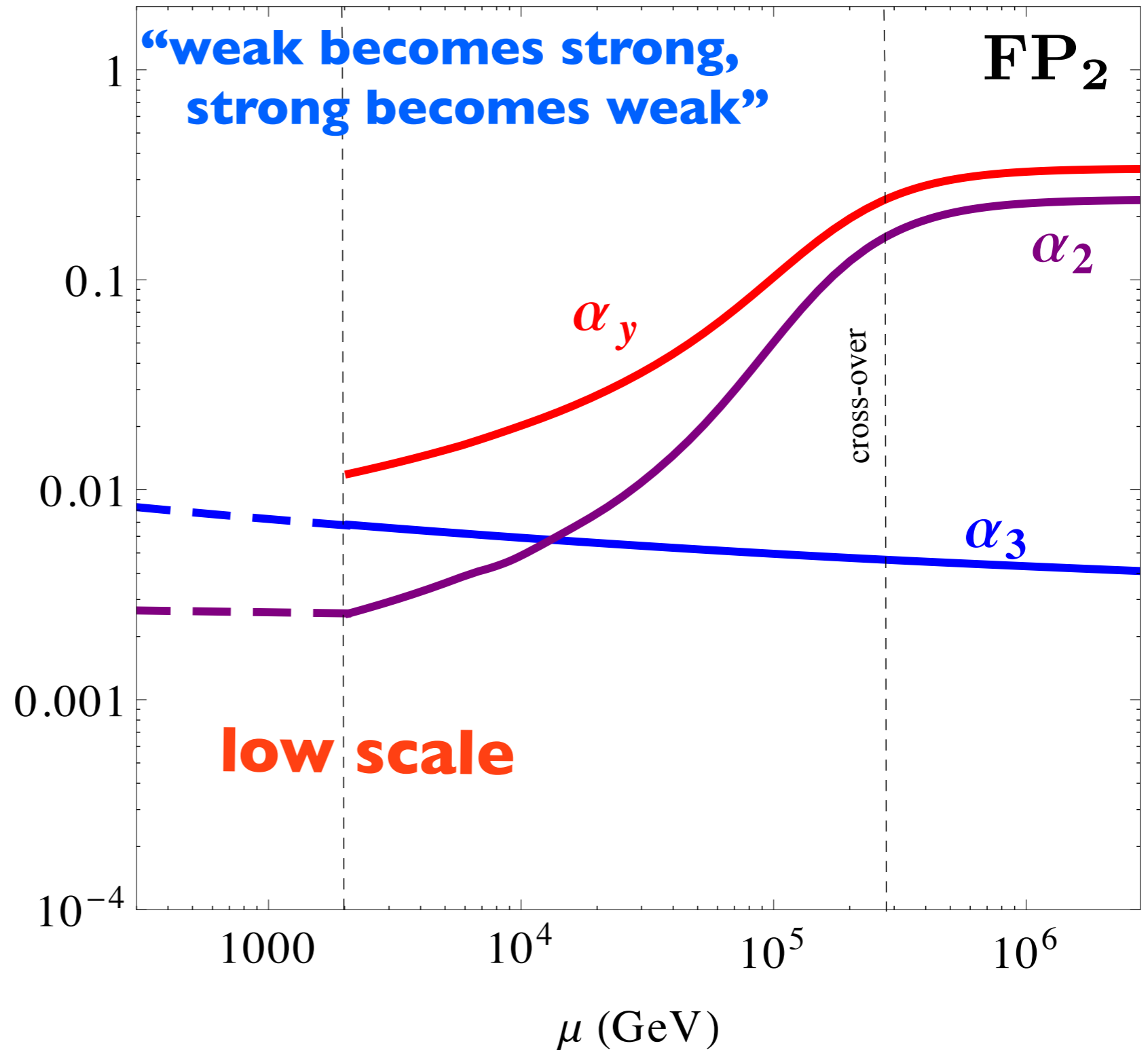


benchmark models

model	parameter (R_3, R_2, N_F)	UV fixed points			type
		α_3^*	α_2^*	α_y^*	
A	(1, 4, 12)	0	0.2407	0.3385	FP ₂ ●
B	(10, 1, 30)	0.1287	0	0.1158	FP ₃ ■
		0.1292	0.2769	0.1163	FP ₄ ◆
C	(10, 4, 80)	0.3317	0	0.0995	FP ₃ ■
		0.0503	0.0752	0.0292	FP ₄ ◆
D	(3, 4, 290)	0	0.8002	0.1500	FP ₂ ●
		0.0416	0.0895	0.0066	FP ₂ ●
E	(3, 3, 72)	0.0615	0.0056		FP ₄ ◆
		0.1499	0.2181	0.0471	FP ₄ ◆

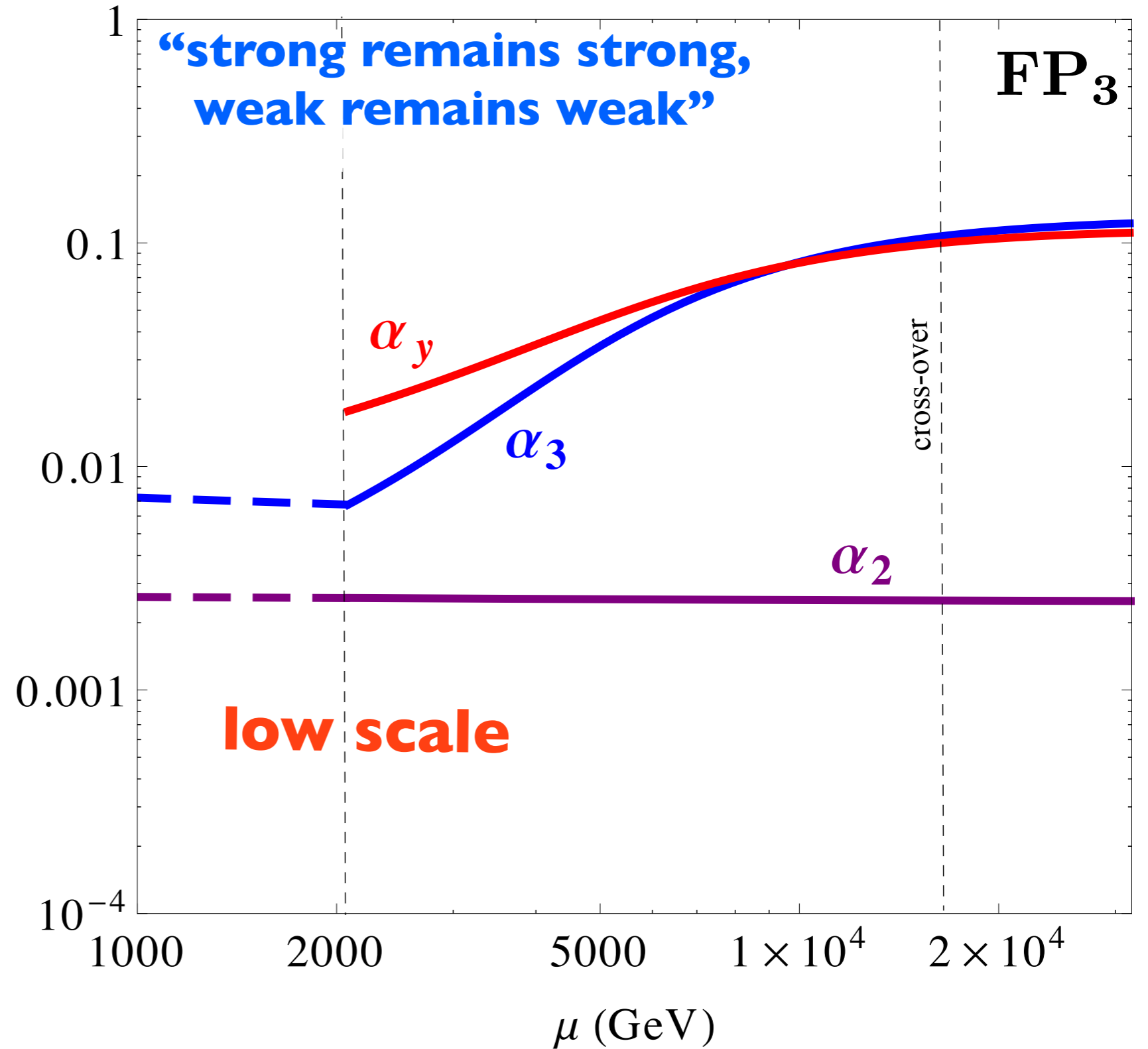
model A

$$(R_3, R_2, N_F) = (1, 4, 12)$$



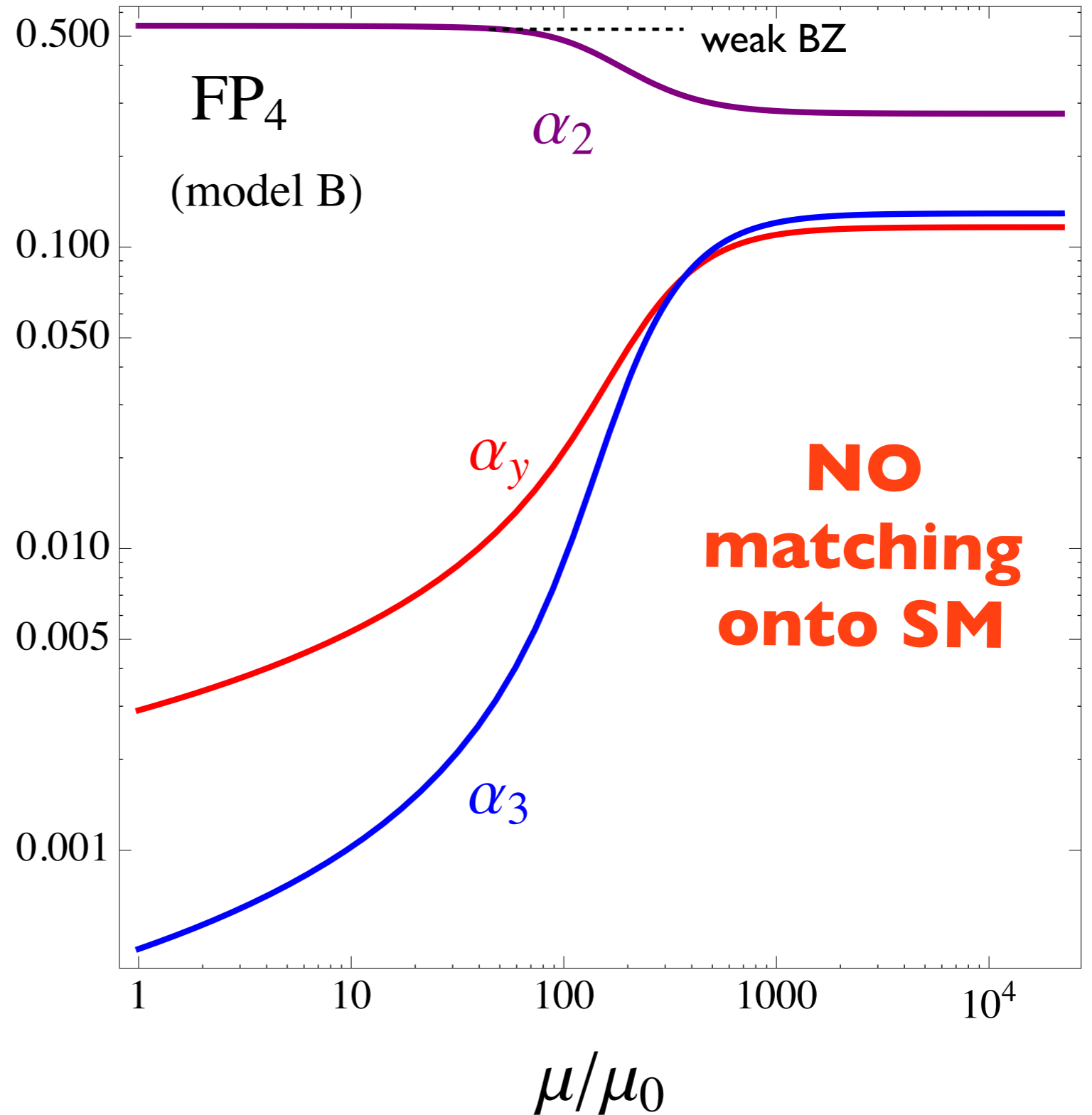
model B

$$(R_3, R_2, N_F) = (10, 1, 30)$$



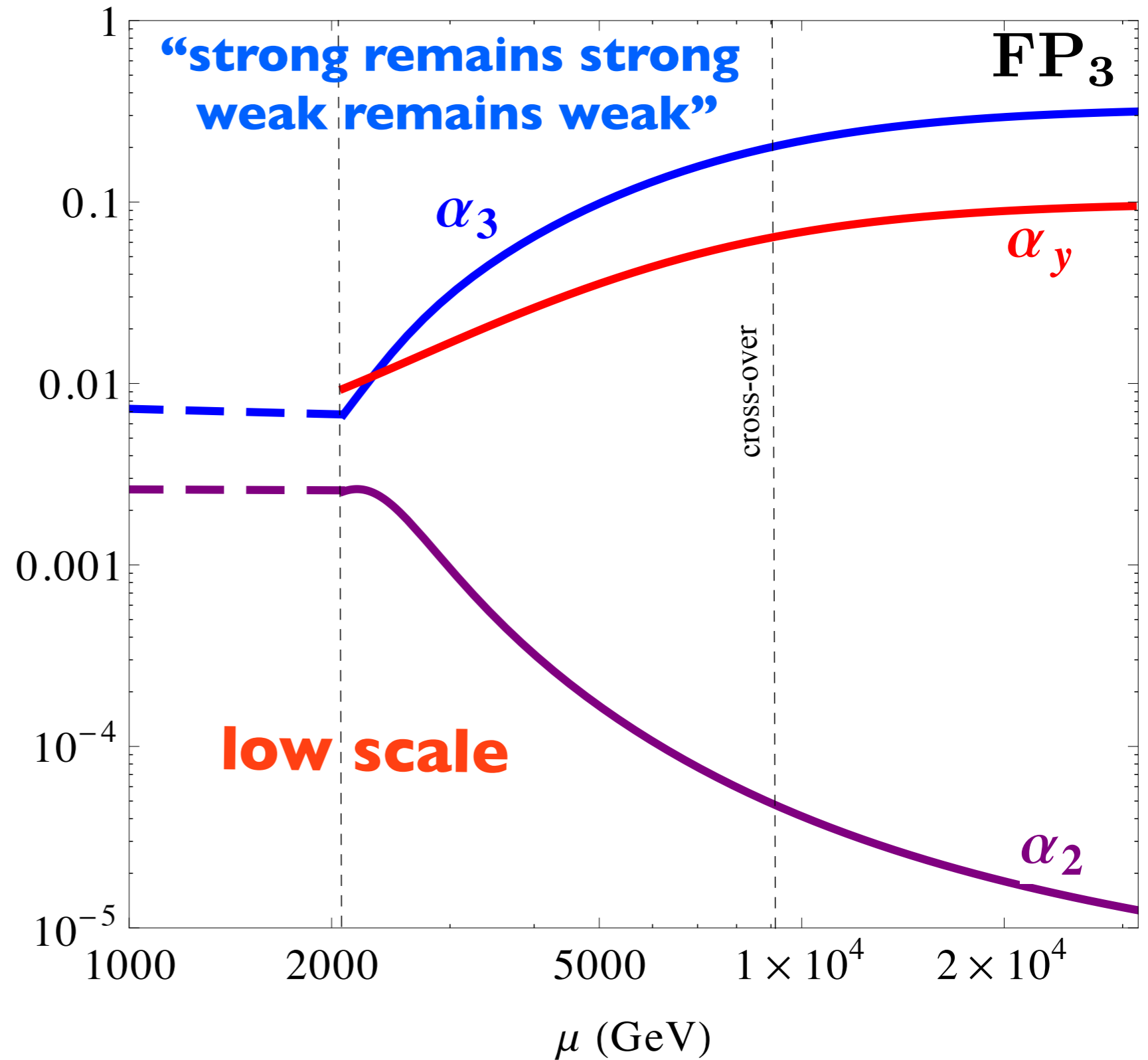
model B

$$(R_3, R_2, N_F) = (10, 1, 30)$$



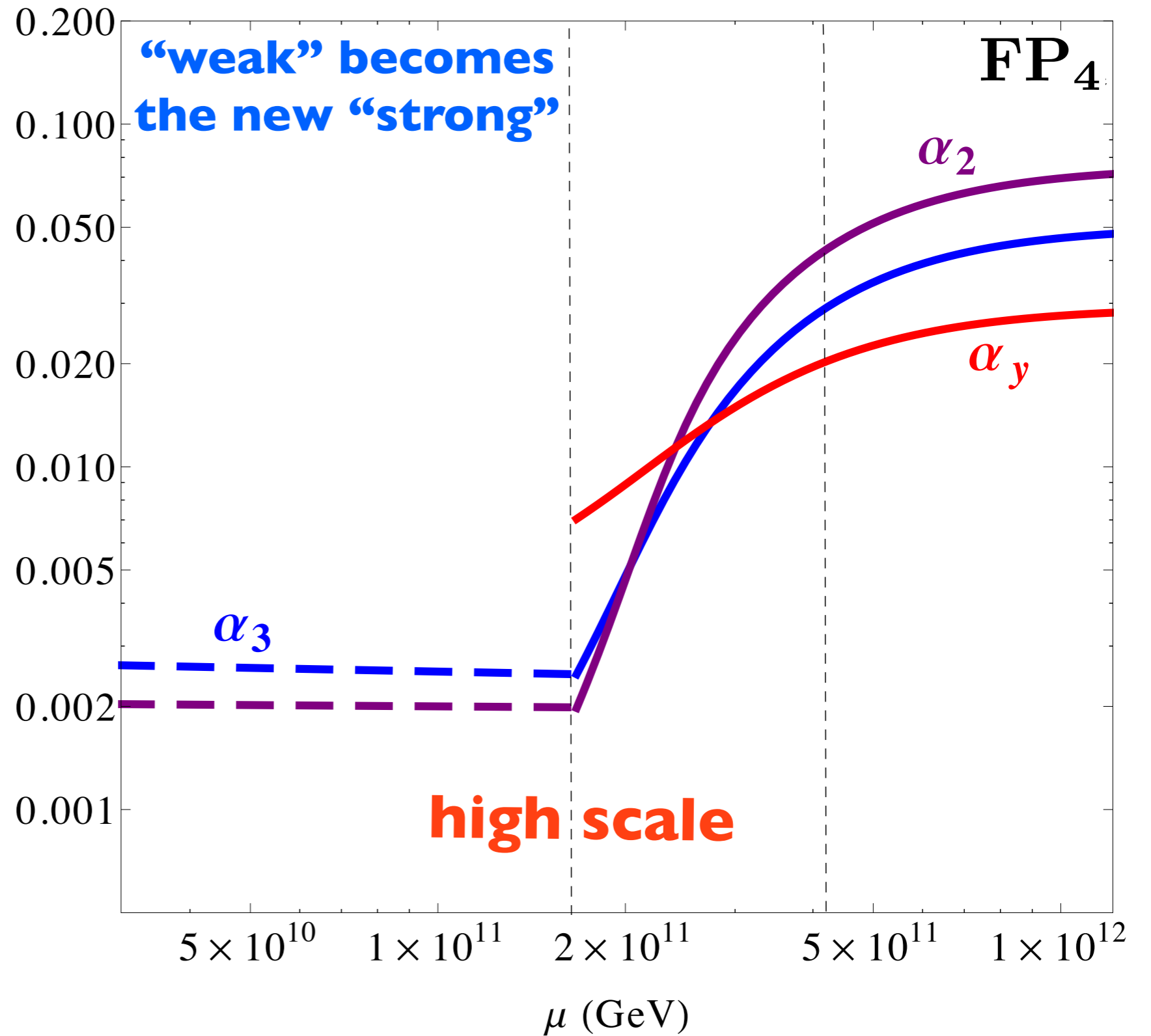
model C

$$(R_3, R_2, N_F) = (10, 4, 80)$$



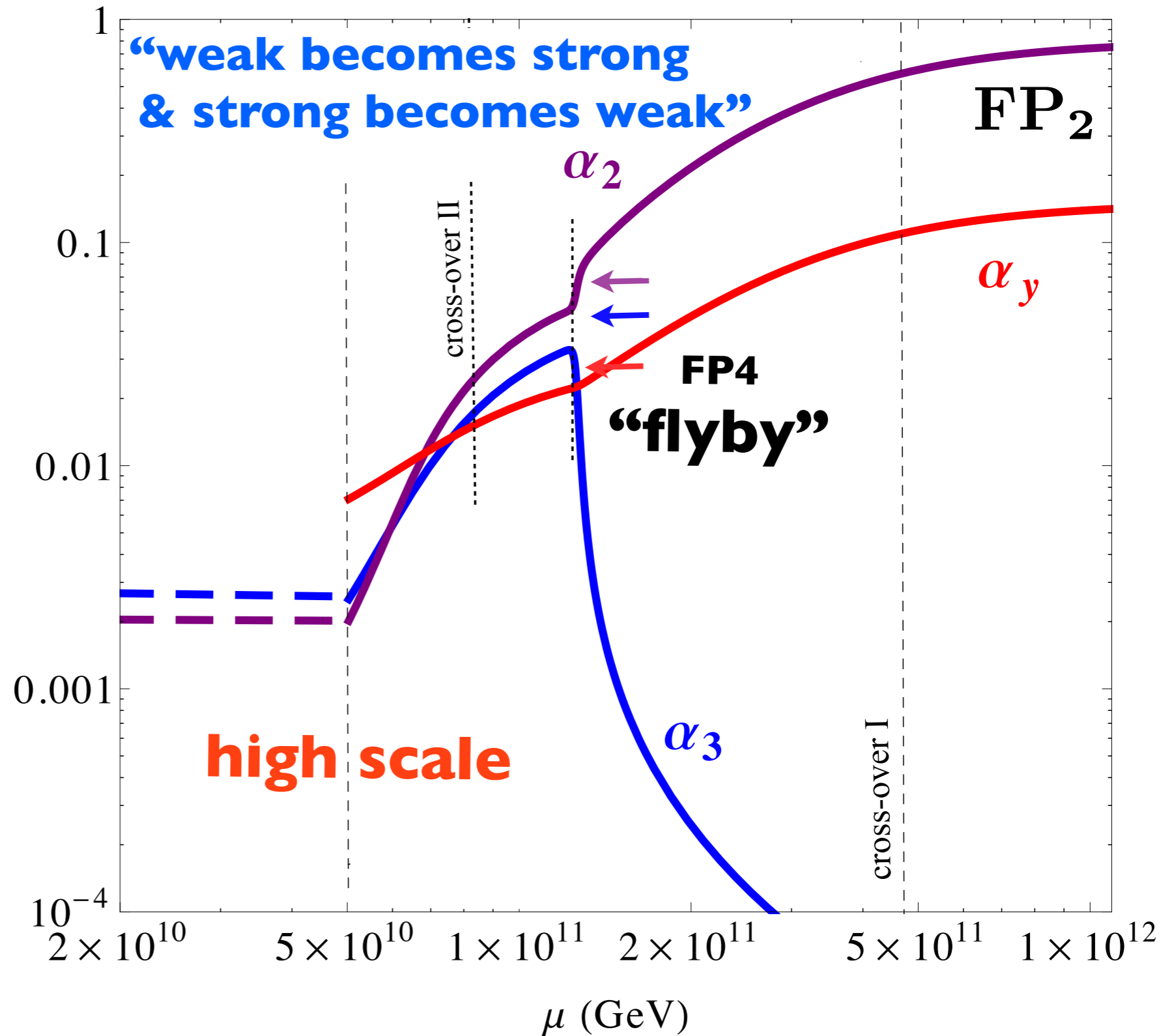
model C

$$(R_3, R_2, N_F) = (10, 4, 80)$$



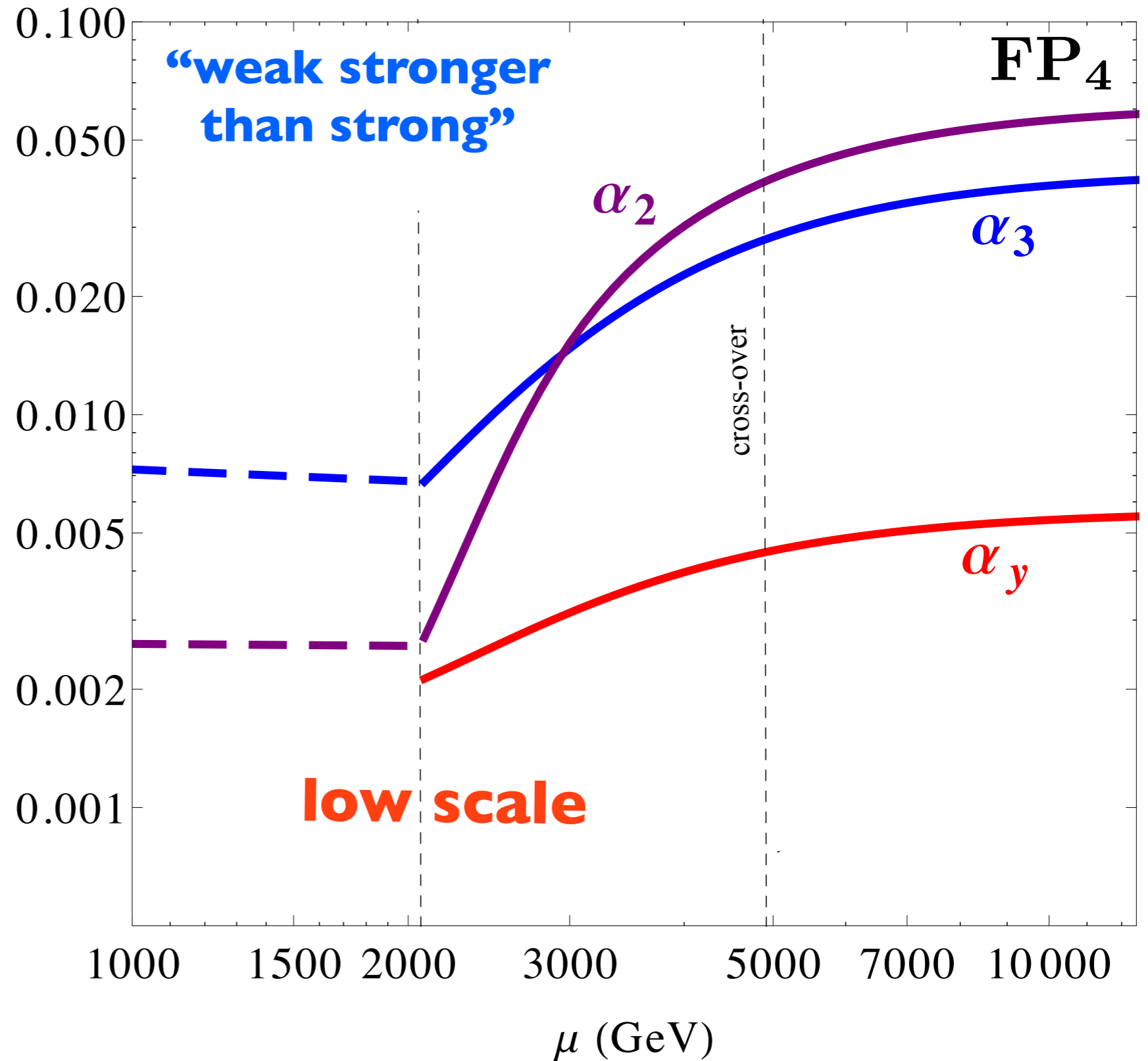
model C

$$(R_3, R_2, N_F) = (10, 4, 80)$$



model D

$$(R_3, R_2, N_F) = (3, 4, 290)$$



summary of SM matching: when it works

FP₂

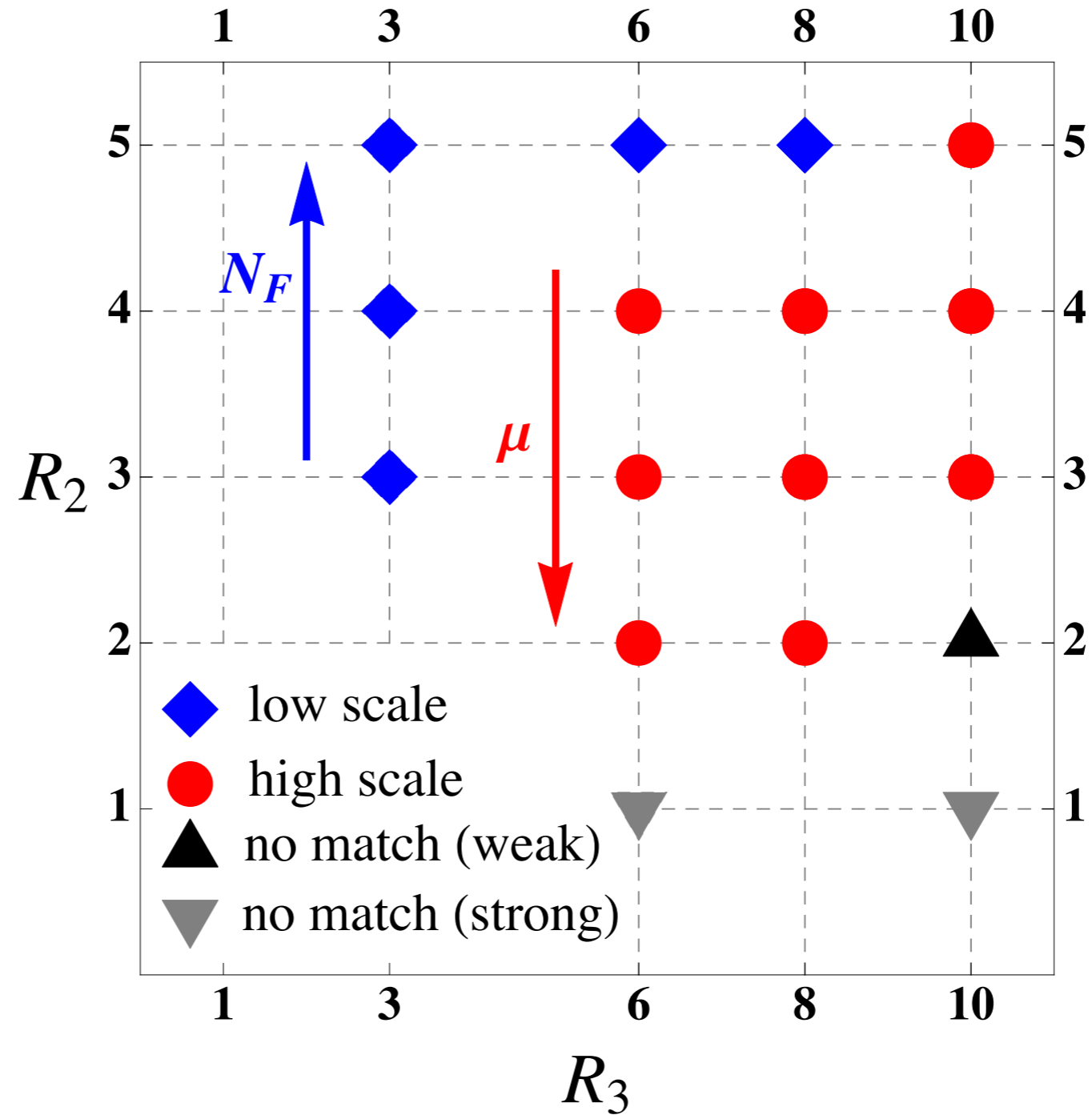
genuinely, except in special circumstances
(competition with other nearby FPs)

FP₃

genuinely, except in special circumstances
(competition with other nearby FPs)

summary of SM matching: when it works

FP₄



asymptotic safety

3. **constraints** from data (colliders)

AD Bond, G Hiller, K Kowalska, DF Litim, 1702.01727

assume **low scale** matching

some BSM masses within **TeV** energy range

assume $R_3 \neq 1$ for LHC

($R_3 = 1$ can be tested at future e^+e^- colliders)

flavor symmetry: stable BSM fermions

broken flavor symmetry: **lightest BSM fermion stable**

constraints from

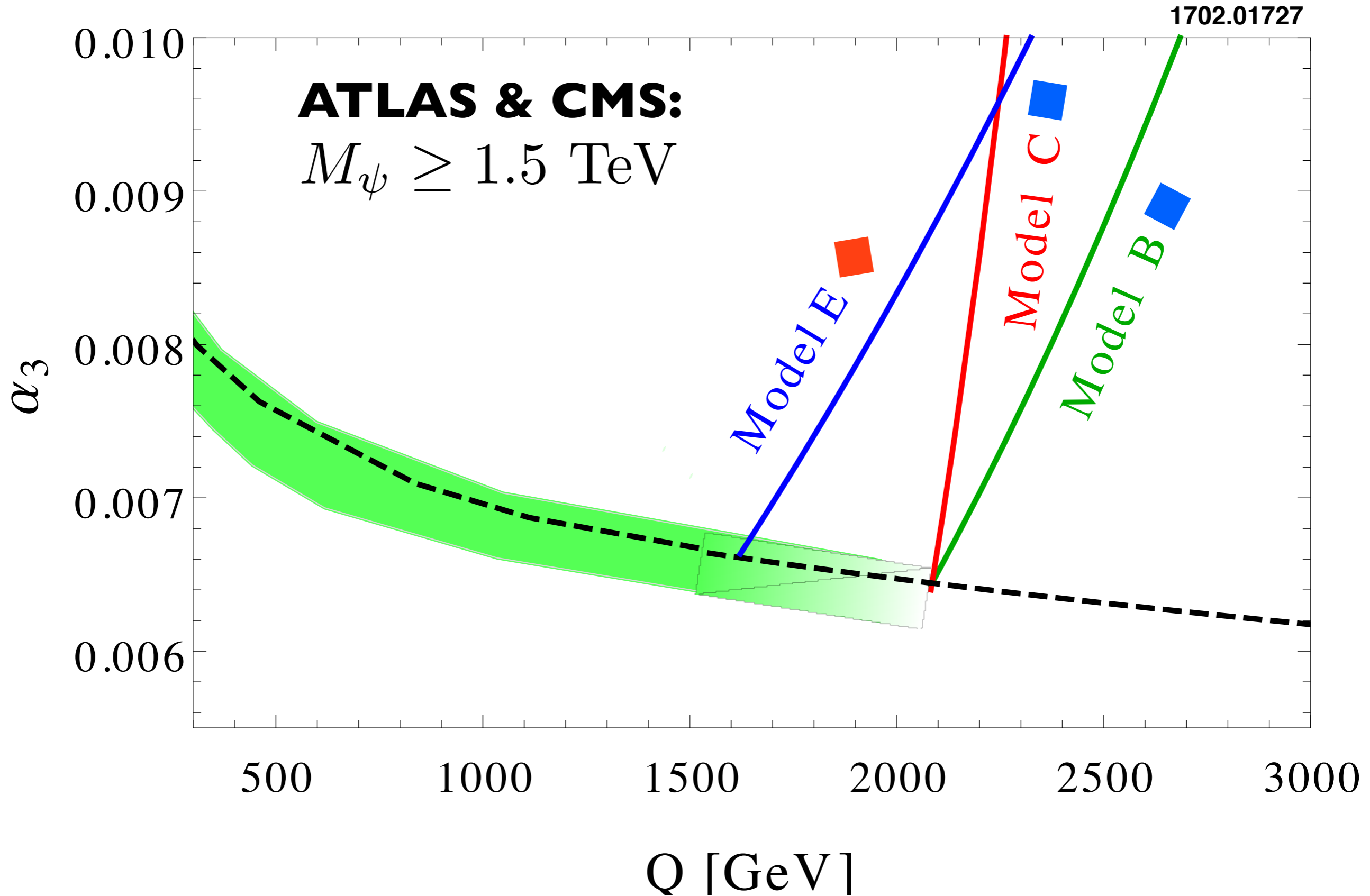
running couplings

the weak sector

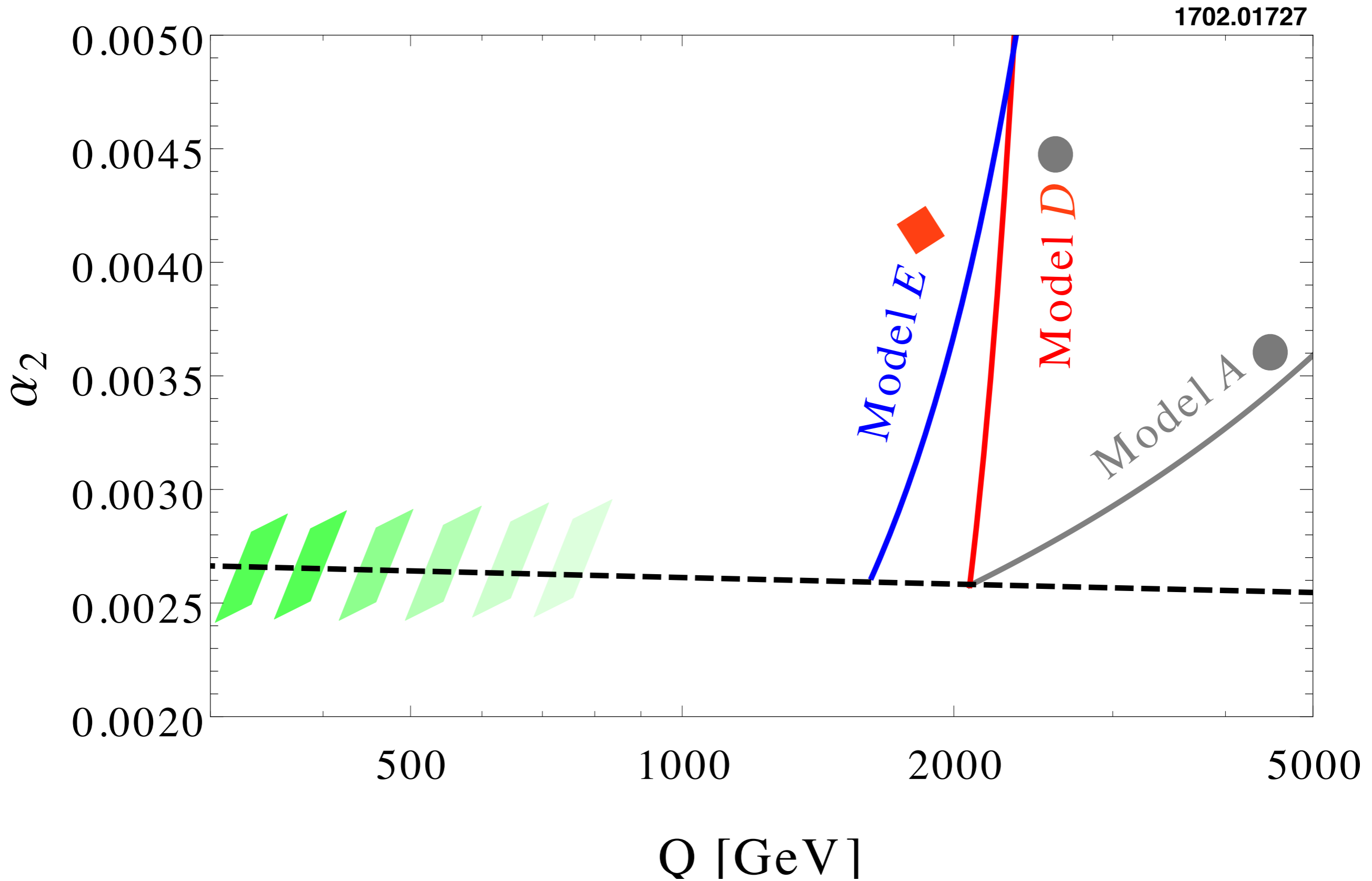
long-lived QCD bound states

di-boson searches

SU(3) BSM running



SU(2) BSM running



di-boson spectra and resonances

assume **resonant production** of BSM scalars

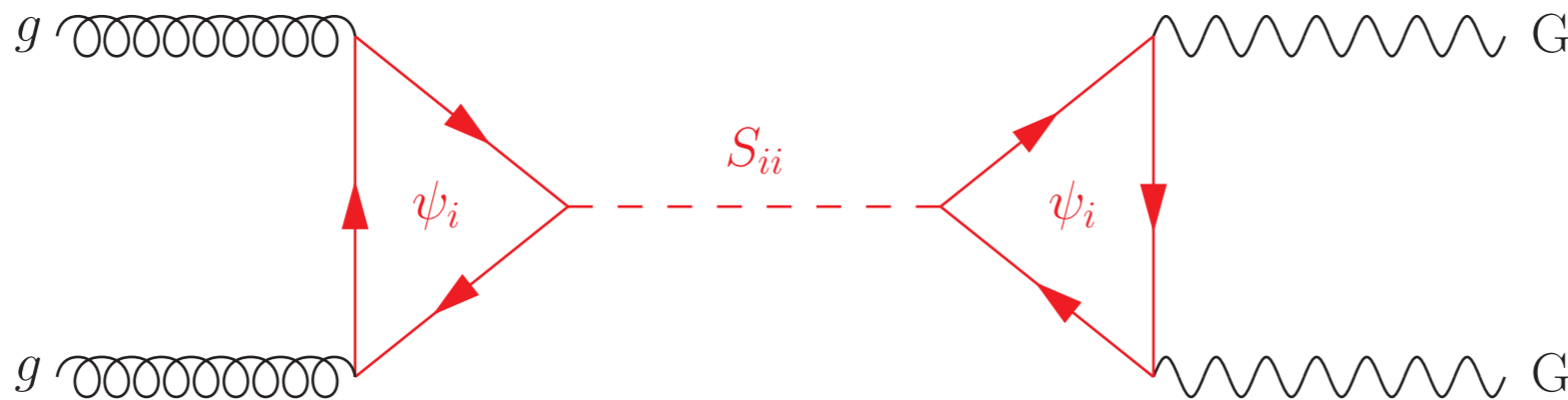
$$M_S < \sqrt{s}$$

$$M_S < 2M_\psi$$

“**low Ms**” $M_S \lesssim M_\psi$

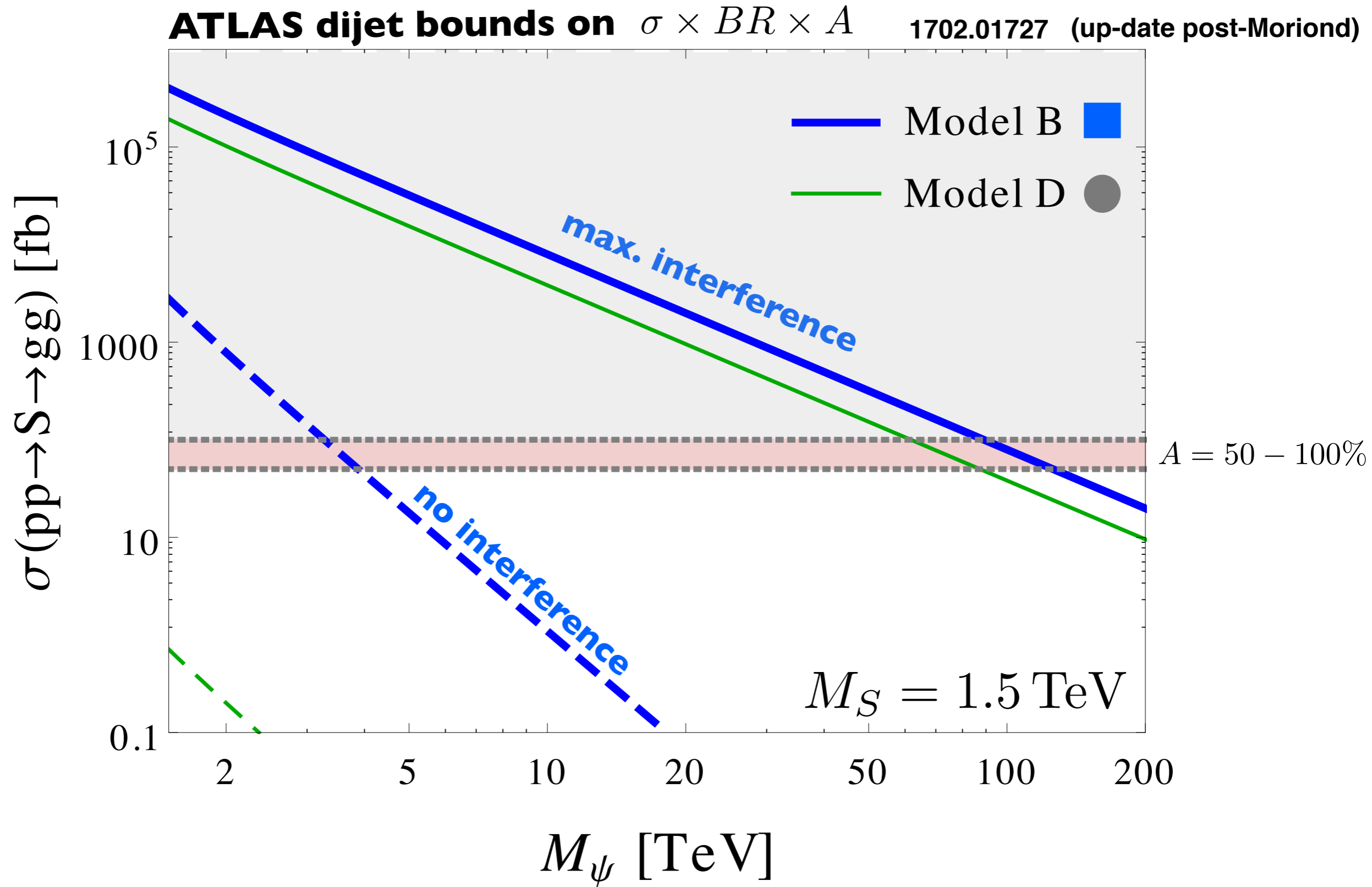
“**high Ms**” $M_\psi \lesssim M_S < 2M_\psi$

loop-mediated decay into $GG = gg, \gamma\gamma, ZZ, Z\gamma,$ or WW

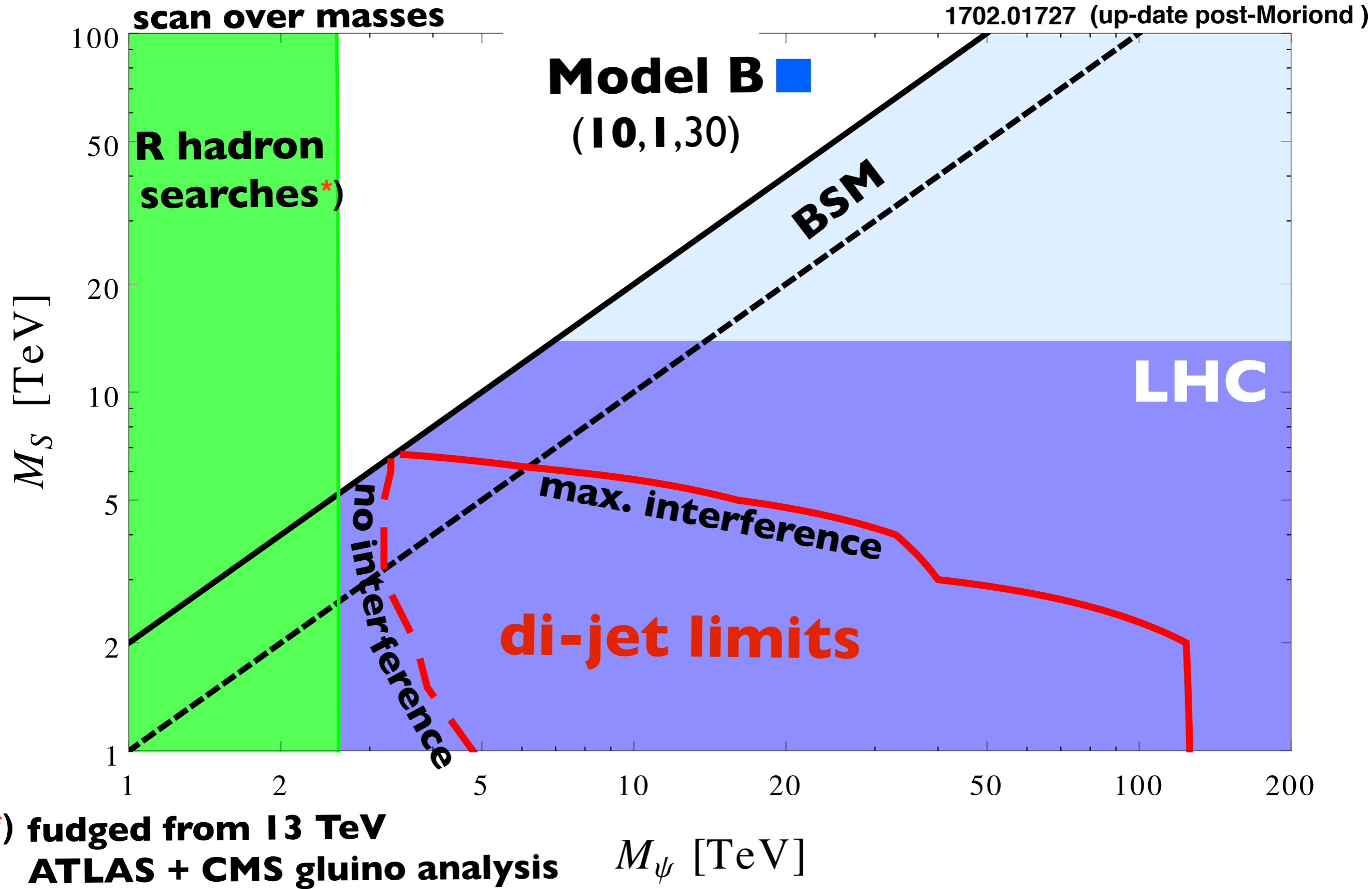


interference effects

dijet cross section



mass exclusion limits



theorems for fixed points and asymptotic safety
new directions beyond asymptotic freedom

weakly interacting **UV completions** of the SM
UV FPs can be **partially or fully** interacting
matching to SM explained, works in many cases

window of opportunities for BSM
new physics can be probed at LHC
constraints from colliders