

# Temperature dependence of shear viscosity of SU(3) gluodynamics within lattice simulation

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## Outline:

- Introduction
- Details of the calculation
- Shear viscosity
  - Fitting of the data
  - Backus-Gilbert method
- Bulk viscosity (preliminary results)
- Conclusion

## Relativistic Hydrodynamics

- $T^{\mu\nu} = (e + p)u^\mu u^\nu + pg^{\mu\nu} + (\eta\nabla^\langle\mu u^{\nu\rangle} + \zeta\Delta^{\mu\nu}\nabla_\alpha u^\alpha) + \dots$   
 $\nabla^\alpha = \Delta^{\alpha\nu}\partial_\nu, \quad \Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$   
 $\nabla^\langle\mu u^{\nu\rangle} = \nabla^\mu u^\nu + \nabla^\nu u^\mu - \frac{2}{3}\Delta^{\mu\nu}\nabla_\alpha u^\alpha$
- EOM  $\partial_\mu T^{\mu\nu} = 0$
- Non-relativistic limit ( $u^\mu = (1, \vec{v})$ )
  - *Continuity equation:*  $\partial_t \rho + \rho(\vec{\partial}\vec{v}) + \vec{v}\vec{\partial}\rho = 0$
  - *Navier–Stokes equation:*  $\frac{\partial v^i}{\partial t} + v^k \frac{\partial v^i}{\partial x^k} = -\frac{1}{\rho} \frac{\partial p}{\partial x^i} - \frac{1}{\rho} \frac{\partial \Pi^{ki}}{\partial x^k}$
  - *Viscous stress tensor:*  $\Pi^{ik} = -\eta\left(\frac{\partial v^i}{\partial x^k} + \frac{\partial v^k}{\partial x^i} - \frac{2}{3}\delta^{ik}\frac{\partial v^l}{\partial x^l}\right) - \zeta\delta^{ik}\frac{\partial v^l}{\partial x^l}$
- $\eta$ –shear viscosity,  $\zeta$ –bulk viscosity

## Calculation of transport coefficients:

- Shear viscosity

$$\eta = \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int_0^{\infty} dt \int d^3x e^{i\omega t} \langle [T_{12}(x), T_{12}(0)] \rangle$$

$$\rho_{12,12}(\omega) = -\frac{1}{\pi} \text{Im} G_R^{12,12}(\omega, \vec{k} = 0)$$

$$\eta = \pi \lim_{\omega \rightarrow 0} \frac{1}{\omega} \rho_{12,12}(\omega)$$

- Bulk viscosity

$$\zeta = \frac{1}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int_0^{\infty} dt \int d^3x e^{i\omega t} \langle [T_{\mu}^{\mu}(x), T_{\nu}^{\nu}(0)] \rangle$$

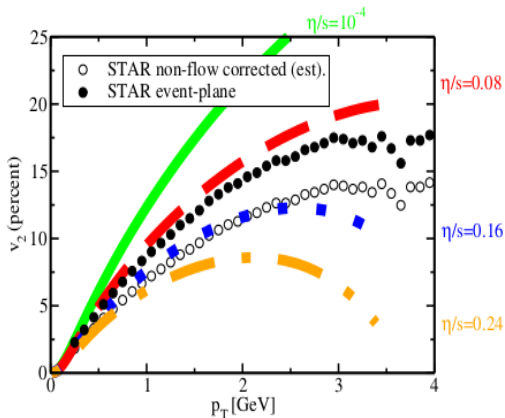
$$\rho_{\mu\mu,\nu\nu}(\omega) = -\frac{1}{\pi} \text{Im} G_R^{\mu\mu,\nu\nu}(\omega, \vec{k} = 0)$$

$$\zeta = \frac{\pi}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \rho_{\mu\mu,\nu\nu}(\omega)$$

- Analytic continuation

$$G_E(\omega, \vec{p}) = -G_R(i\omega, \vec{p}), \quad \omega > 0$$

$$C_E(t) = \int_0^{\infty} d\omega \rho(\omega) \frac{\text{ch}\left(\frac{\omega}{2T} - \omega t\right)}{\text{sh}\left(\frac{\omega}{2T}\right)}$$

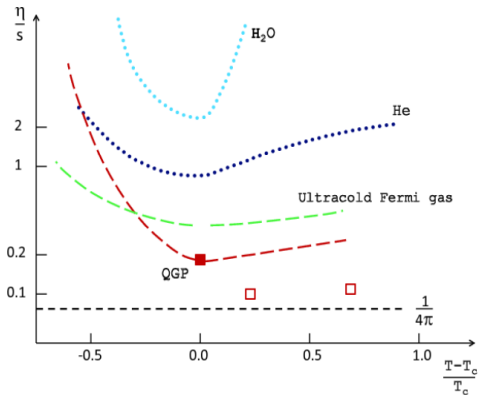


Elliptic flow from STAR experiment (Nucl. Phys. A 757, 102 (2005))

$$\frac{dN}{d\phi} \sim (1 + 2v_1 \cos(\phi) + 2v_2 \cos^2(\phi))$$

Quark-gluon plasma is close to ideal liquid ( $\frac{\eta}{s} = (1 - 3) \frac{1}{4\pi}$ )

M. Luzum and P. Romatschke, Phys. Rev. C 78, 034915 (2008)

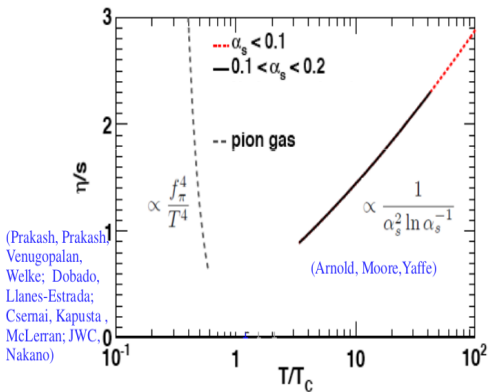


S.Cremonini, U.Gursoy, P.Szepietowski, JHEP 1208 (2012) 167

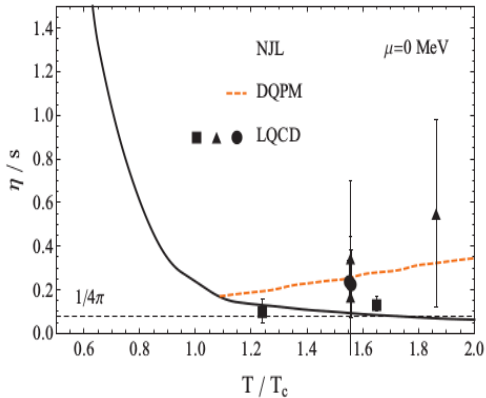
Comparison of different liquids

**Quark-gluon plasma is the most ideal liquid**

## Study of shear viscosity (perturbative calculation + pion gas)



## Study of shear viscosity (effective models)



R. Marty, E. Bratkovskaya, W. Cassing, J. Aichelin and H. Berrehrh, Phys. Rev. C 88, 045204 (2013)



## Other works (SU(3) gluodynamics):

- Karsch, F. et al. Phys.Rev. D35 (1987)
- A. Nakamura, S. Sakai Phys. Rev. Lett. 94, 072305 (2005)
- H. B. Meyer, Phys.Rev. D76 (2007) 101701
- H. B. Meyer, Nucl.Phys. A830 (2009) 641C-648C

## Results:

- $\frac{\eta}{s} = 0.134 \pm 0.033$  ( $T/T_c = 1.65, 8 \times 28^3$ )
- $\frac{\eta}{s} = 0.102 \pm 0.056$  ( $T/T_c = 1.24, 8 \times 28^3$ )
- $\frac{\eta}{s} = 0.20 \pm 0.03$  ( $T/T_c = 1.58, 16 \times 48^3$ )
- $\frac{\eta}{s} = 0.26 \pm 0.03$  ( $T/T_c = 2.32, 16 \times 48^3$ )

## SU(2) gluodynamics:

- $\frac{\eta}{s} = 0.134 \pm 0.057$  ( $T/T_c = 1.2, 16 \times 32^3$ )

N.Yu. Astrakhantsev, V.V. Braguta, A.Yu. Kotov, JHEP 1509 (2015) 082

# Lattice calculation of shear viscosity

The first step:

Measurement of the correlation function:

$$C_E(t) = \langle T_{12}(t) T_{12}(0) \rangle$$

The second step:

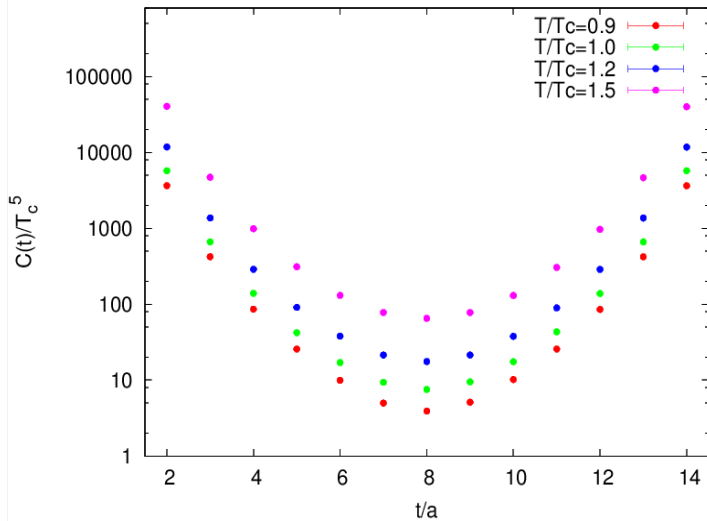
Calculation of the spectral function  $\rho(\omega)$ :

$$C_E(t) = \int_0^\infty d\omega \rho(\omega) \frac{ch\left(\frac{\omega}{2T} - \omega t\right)}{sh\left(\frac{\omega}{2T}\right)}$$
$$\eta = \pi \lim_{\omega \rightarrow 0} \frac{\rho(\omega)}{\omega}$$

## Details of the calculation

- SU(3) gluodynamics
- Two-level algorithm
- Lattice size  $32^3 \times 16$
- Temperatures  $T/T_c = 0.9, 0.925, 0.95, 1.0, 1.1, 1.2, 1.35, 1.5$
- Accuracy  $\sim 2 - 3\%$  at  $t = \frac{1}{2T}$
- $\langle T_{12}(x) T_{12}(y) \rangle \sim (\langle T_{11}(x) T_{11}(y) \rangle - \langle T_{11}(x) T_{22}(y) \rangle)$
- Clover discretization for the  $\hat{F}_{\mu\nu}$
- Renormalization of EMT: F. Karsch, Nucl.Phys. B205 (1982) 285-300
- ...

## Correlation functions



## Spectral function

$$C(t) = \int_0^\infty d\omega \rho(\omega) \frac{\text{ch}\left(\frac{\omega}{2T} - \omega t\right)}{\text{sh}\left(\frac{\omega}{2T}\right)}$$

Properties of the spectral function:

- $\rho(\omega) \geq 0$ ,  $\rho(-\omega) = -\rho(\omega)$
- Asymptotic freedom:  $\rho(\omega)|_{\omega \rightarrow \infty}^{NLO} = \frac{1}{10} \frac{d_A}{(4\pi)^2} \omega^4 \left(1 - \frac{5N_c \alpha_s}{9\pi}\right)$   
~ 90% of the total contribution  $t = 1/2/T$
- Hydrodynamics:  $\rho(\omega)|_{\omega \rightarrow 0} = \frac{\eta}{\pi} \omega$

## Spectral function

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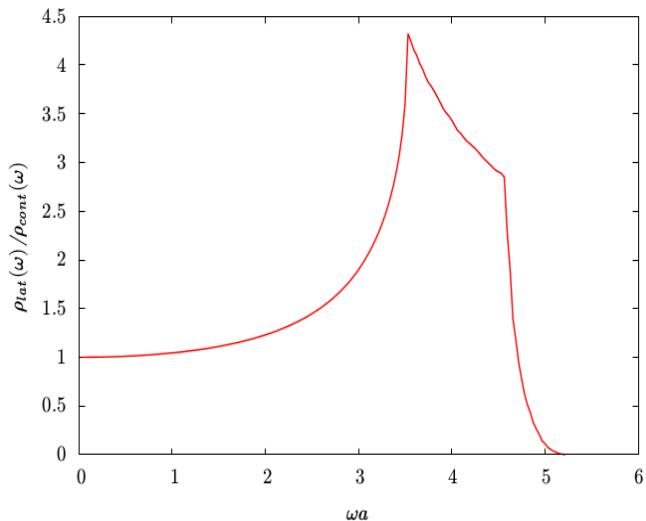
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Ansatz for the spectral function (QCD sum rules motivation)

$$\rho(\omega) = \frac{\eta}{\pi} \omega \theta(\omega_0 - \omega) + A \rho_{lat}(\omega) \theta(\omega - \omega_0)$$

## Lattice spectral function $\rho_{lat}$

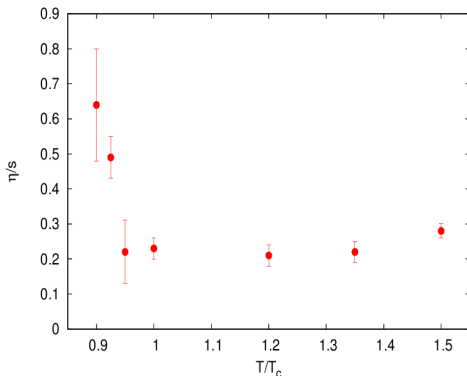
Takes into account discretization errors in temporal direction



## Spectral function

$$\rho_1(\omega) = \frac{\eta}{\pi} \omega \theta(\omega_0 - \omega) + A \rho_{lat}(\omega) \theta(\omega - \omega_0)$$

$\chi^2/dof \sim 1, A \sim 1, \omega_0/T \sim 7 - 8$



## Two additional ansatzs:

- $\rho_2(\omega) = \frac{1}{2} B \omega (1 + \tanh[\gamma(\omega_0 - \omega)]) + \frac{1}{2} A \rho_{lat}(\omega) (1 + \tanh[\gamma(\omega - \omega_0)])$
- $\rho_3(\omega) = B \omega (1 + C\omega^2) \theta(\omega_0 - \omega) + A \rho_{lat}(\omega) \theta(\omega - \omega_0)$



## Properties of the spectral function

- Hydrodynamical approximation works well up to  $\omega < \pi T \sim 1 \text{ GeV}$  (H.B. Meyer, arXiv:0809.5202)
- Asymptotic freedom works well from  $\omega > 3 \text{ GeV}$
- Poor knowledge of the spectral function in the region  $\omega \in (1, 3) \text{ GeV}$   
 $\Rightarrow$  Main source of uncertainty in the fitting procedure

## Backus-Gilbert method for the spectral function

- Problem: find  $f(\omega)$  from the integral equation

$$C(x_i) = \int_0^\infty d\omega f(\omega) K(x_i, \omega), \quad K(x_i, \omega) = \frac{\text{ch}\left(\frac{\omega}{2T} - \omega x_i\right)}{\text{sh}\left(\frac{\omega}{2T}\right)}$$

- Define an estimator  $\tilde{f}(\bar{\omega})$  ( $\delta(\bar{\omega}, \omega)$  - resolution function):

$$\tilde{f}(\bar{\omega}) = \int_0^\infty d\omega \hat{\delta}(\bar{\omega}, \omega) f(\omega)$$

- Let us expand  $\delta(\bar{\omega}, \omega)$  as

$$\delta(\bar{\omega}, \omega) = \sum_i b_i(\bar{\omega}) K(x_i, \omega) \quad \tilde{f}(\bar{\omega}) = \sum_i b_i(\bar{\omega}) C(x_i)$$

- Goal: minimize the width of the resolution function

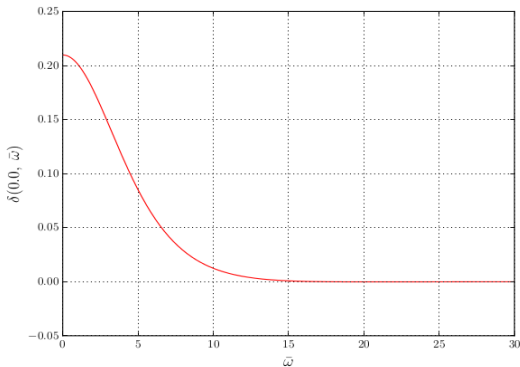
$$b_i(\bar{\omega}) = \frac{\sum_j W_{ij}^{-1} R_j}{\sum_{ij} R_i W_{ij}^{-1} R_j},$$

$$W_{ij} = \int d\omega K(x_i, \omega) (\omega - \bar{\omega})^2 K(x_j, \omega), \quad R_i = \int d\omega K(x_i, \omega)$$

- Regularization by the covariance matrix  $S_{ij}$ :

$$W_{ij} \rightarrow \lambda W_{ij} + (1 - \lambda) S_{ij}, \quad 0 < \lambda < 1$$

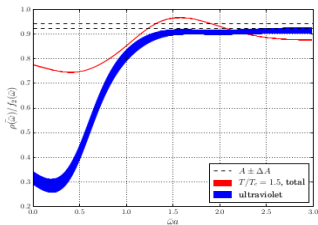
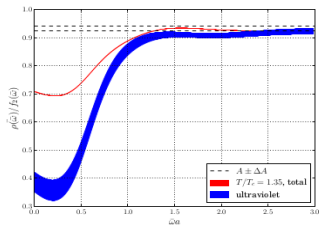
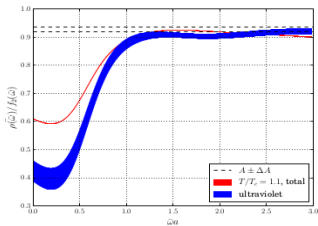
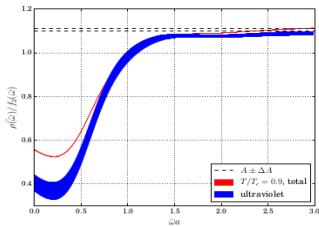
## Resolution function $\delta(0, \omega)$ ( $T/T_c = 1.2$ , $\lambda = 0.001$ )



- Width of the resolution function  $\omega/T \sim 4$
- Hydrodynamical approximation works up to  $\omega/T < \pi$
- Problem: large contribution from ultraviolet tail (more than  $\sim 100\%$ )

## Model for the ultraviolet contribution

$$\rho_{ultr} = A\rho_{lat}(\omega)\theta(\omega - \omega_0)$$

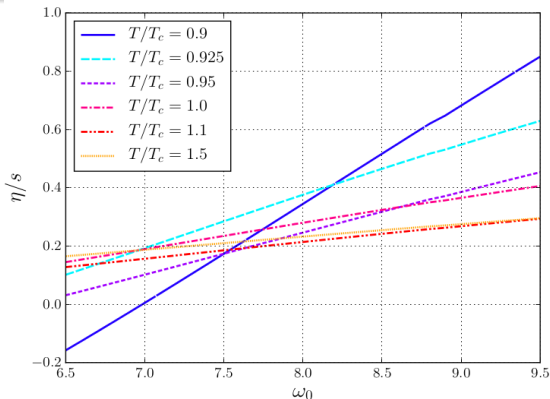


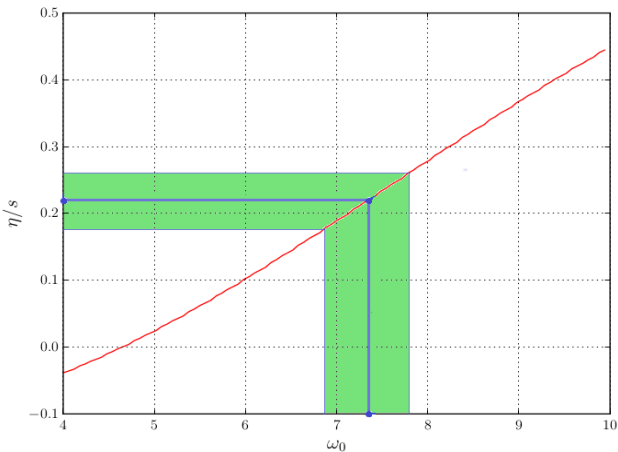
## Solution:

- Take ultraviolet contribution in the form:

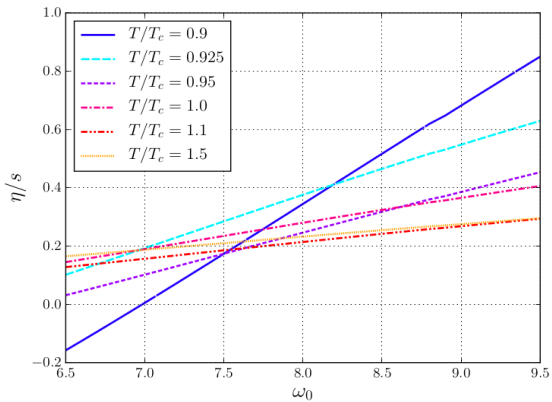
$$\rho_{ultr} = A\rho_{lat}(\omega)\theta(\omega - \omega_0)$$

- Determine the value of the constant A from the BG method
- Subtract ultraviolet contribution and obtain  $\eta/s$  as a function of  $\omega_0$

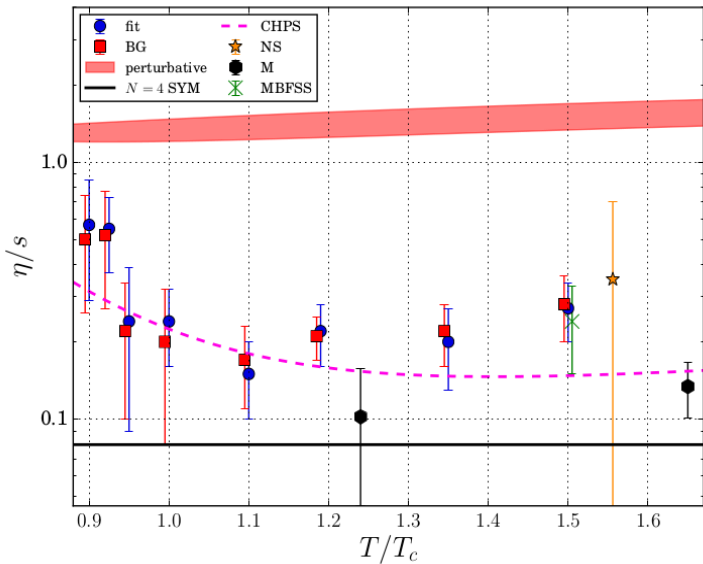




- For  $T/T_c = 1$ .  $\omega_0/T = 7.33 \pm 0.47$
- $\eta/s = 0.22 \pm 0.04$



# Our results

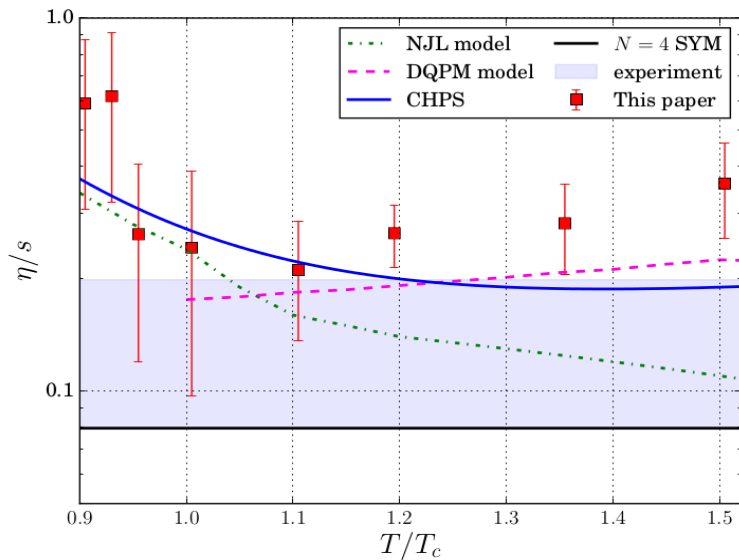




Estimation of the  $\eta/s$  for QCD with dynamical quarks ( $N_f = 3$ ):

$$\left(\frac{\eta}{s}\right)_{QCD} = \left(\frac{(\eta/s)_{QCD}}{(\eta/s)_{Gluedynamics}}\right)_{pert} \times \left(\frac{\eta}{s}\right)_{Gluedynamics}$$

## Our results



## Bulk viscosity

- $\zeta = \frac{1}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int_0^\infty dt \int d^3x e^{i\omega t} \langle [T_\mu^\mu(x), T_\nu^\nu(0)] \rangle$
- $C_E(t) = \int_0^\infty d\omega \rho_{\mu\mu,\nu\nu}(\omega) \frac{ch\left(\frac{\omega}{2T} - \omega t\right)}{sh\left(\frac{\omega}{2T}\right)}$
- $\zeta = \frac{\pi}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \rho_{\mu\mu,\nu\nu}(\omega)$

## Violation of conformal symmetry in QCD

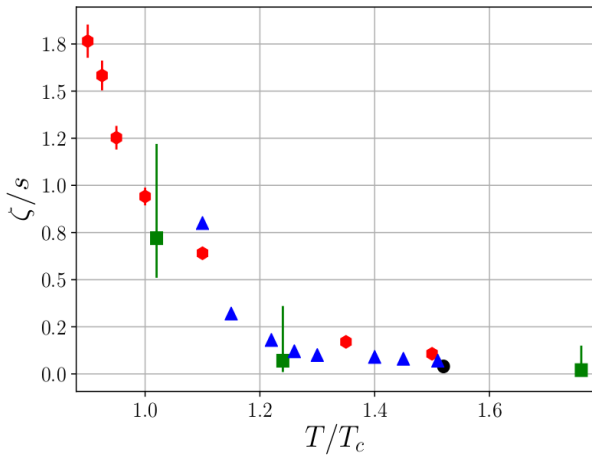
- In conformal theory  $T_\mu^\mu = 0$
- Conformal symmetry is violated by quantum corrections
- Trace anomaly:  $T_\mu^\mu = \frac{\beta(g)}{2g} \text{Tr}[G^{\mu\nu} G_{\mu\nu}] + \sum_i m_i \bar{q}q$

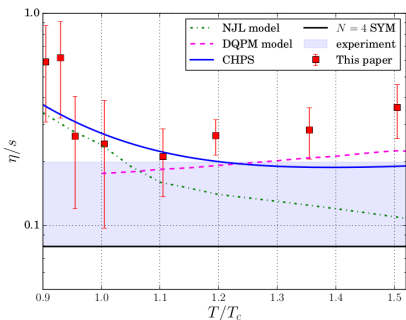
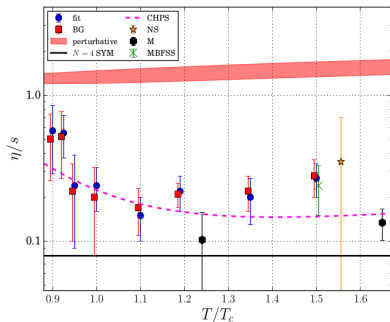
## Properties of spectral density:

- $\rho_{\mu\mu,\nu\nu}(\omega) \geq 0$ ,  $\rho_{\mu\mu,\nu\nu}(-\omega) = -\rho_{\mu\mu,\nu\nu}(\omega)$
- Asymptotic freedom:  $\rho(\omega)|_{\omega \rightarrow \infty}^{NLO} = d_A \left( \frac{11\alpha_s}{(4\pi)^2} \right)^2 \omega^4$
- Hydrodynamics:  $\rho(\omega)|_{\omega \rightarrow 0} = 9 \frac{\zeta}{\pi} \omega$
- The contribution of asymptotic freedom is suppressed ( $\sim \alpha_s^2$ )

## Bulk viscosity from midpoint

$$C_E(\beta/2) = \int_0^\infty d\omega \rho_{hyd}(\omega) \frac{1}{sh\left(\frac{\omega}{2T}\right)}$$





## Conclusion:

- We calculated  $\eta/s$  for set of temperatures  $T/T_c \in (0.9, 1.5)$
- Applied fitting procedure and Backus-Gilbert method for the SF
- $\eta/s$  is close to  $N=4$  SYM and in agreement with experiment
- Large deviation from perturbative results