

# Is there a low $p_T$ anomaly in the pion momentum spectrum at LHC?

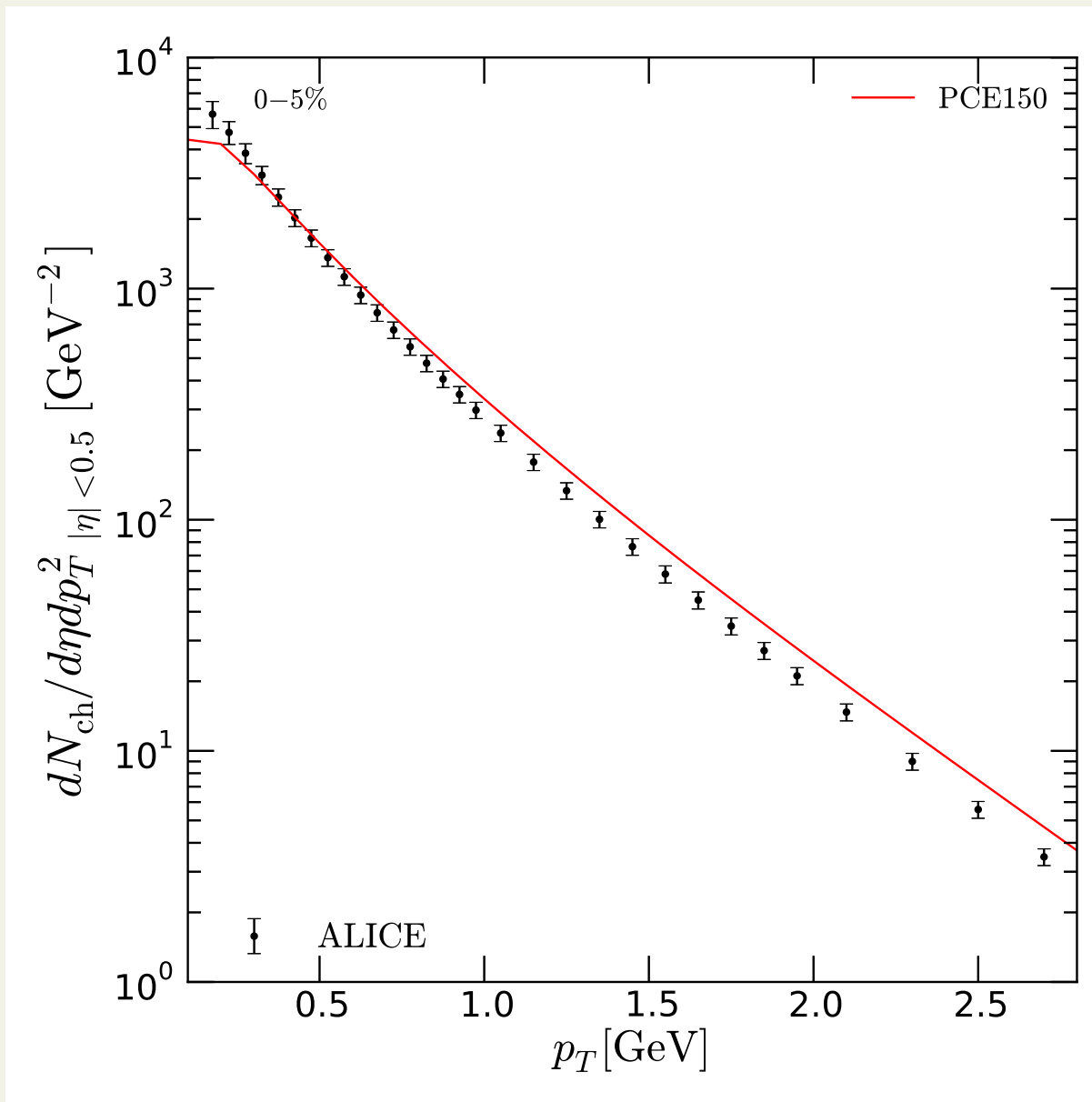
**Pasi Huovinen**  
Uniwersytet Wrocławski

**Transport meeting**

June 16, 2014, **Institut für Theoretische Physik, Frankfurt**

in collaboration with  
**P. M. Lo, M. Marczenko, K. Redlich, and C. Sasaki**

# Charged hadron $p_T$ spectrum at LHC

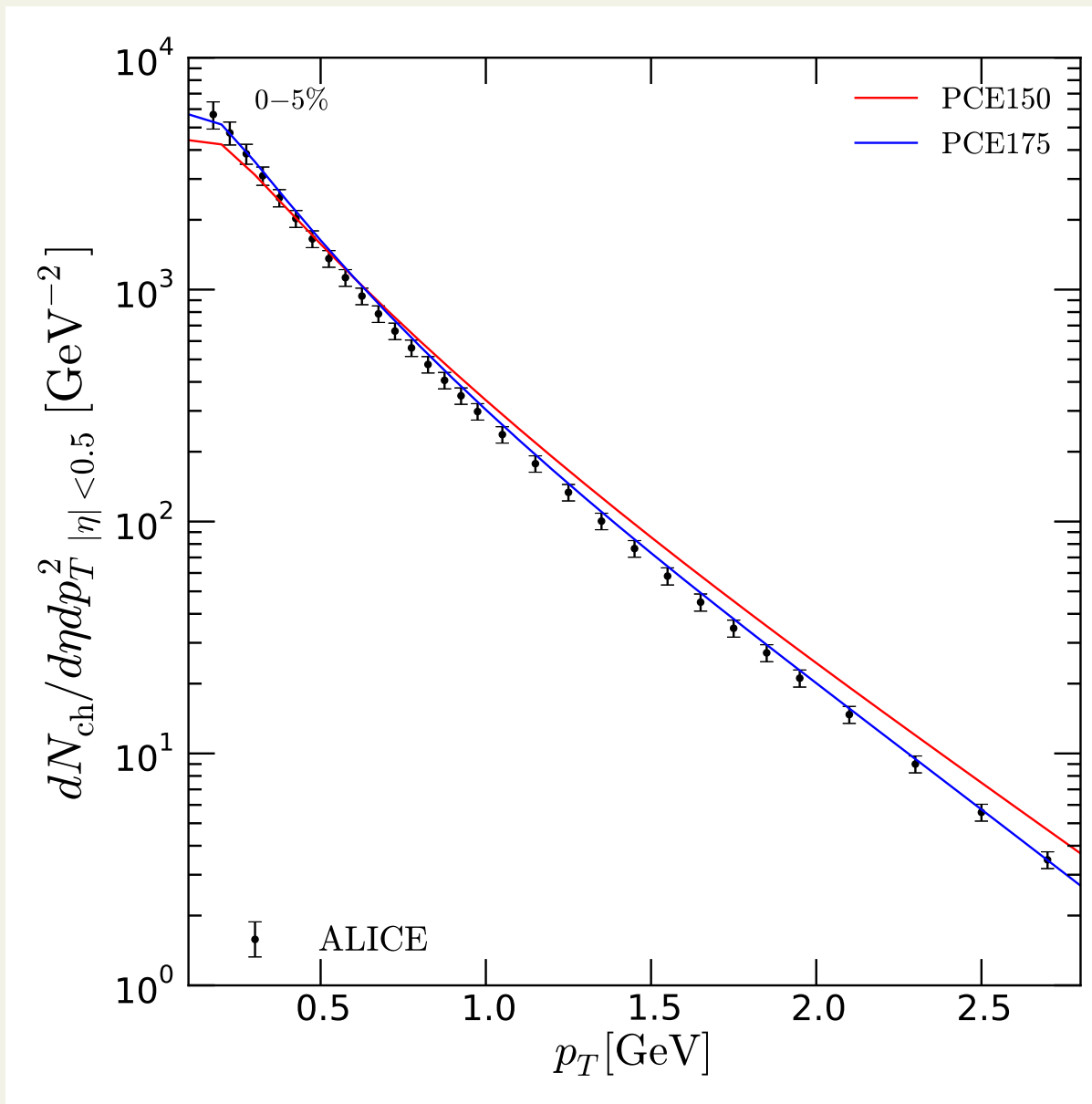


- viscous hydro
- initial state:  
pQCD+saturation
- $\tau_0 \approx 0.2 \text{fm}/c$

**PCE150:**  
fit to  $\pi$ ,  $K$ ,  $p$  yields  
no fit to spectrum

©H. Niemi

# Charged hadron $p_T$ spectrum at LHC



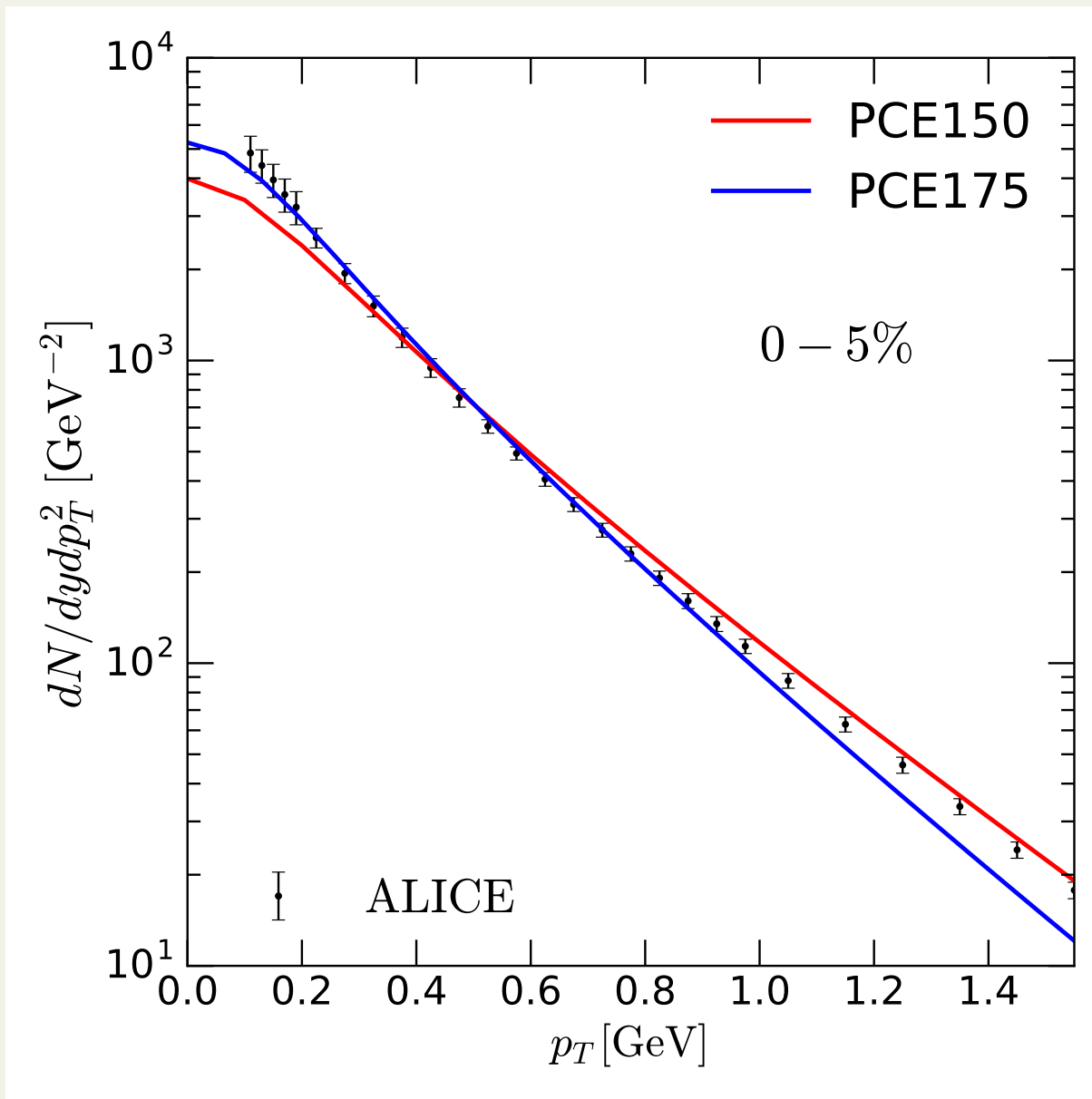
- viscous hydro
- initial state:  
pQCD+saturation
- $\tau_0 \approx 0.2\text{fm}/c$

**PCE150:**  
fit to  $\pi$ ,  $K$ ,  $p$  yields  
no fit to spectrum

**PCE175:**  
no fit to yields  
fits the spectrum

©H. Niemi

# Pion $p_T$ spectrum at LHC



©H. Niemi

- need more resonances
- yield proportional to Boltzmann factor

$$N \propto \exp\left(-\frac{m}{T}\right)$$

- need more resonances
- yield proportional to Boltzmann factor

$$N \propto \exp\left(-\frac{m}{T}\right)$$

- resonance mass?

- **need more resonances**
- **yield proportional to Boltzmann factor**

$$N \propto \exp\left(-\frac{m}{T}\right)$$

- **resonance mass?**
- **usually no width, i.e. resonances have their pole mass**

## effect of Breit-Wigner width on number density:

$$n = \int d^3\mathbf{p} f(p)$$
$$\Rightarrow n = \int d^3\mathbf{p} \int dm^2 \frac{d\rho}{dm^2} f(p, m)$$

where

$$\frac{d\rho}{dm^2} = \frac{1}{N} \frac{m_R \Gamma}{(m^2 - m_R^2)^2 + m_R^2 \Gamma^2},$$

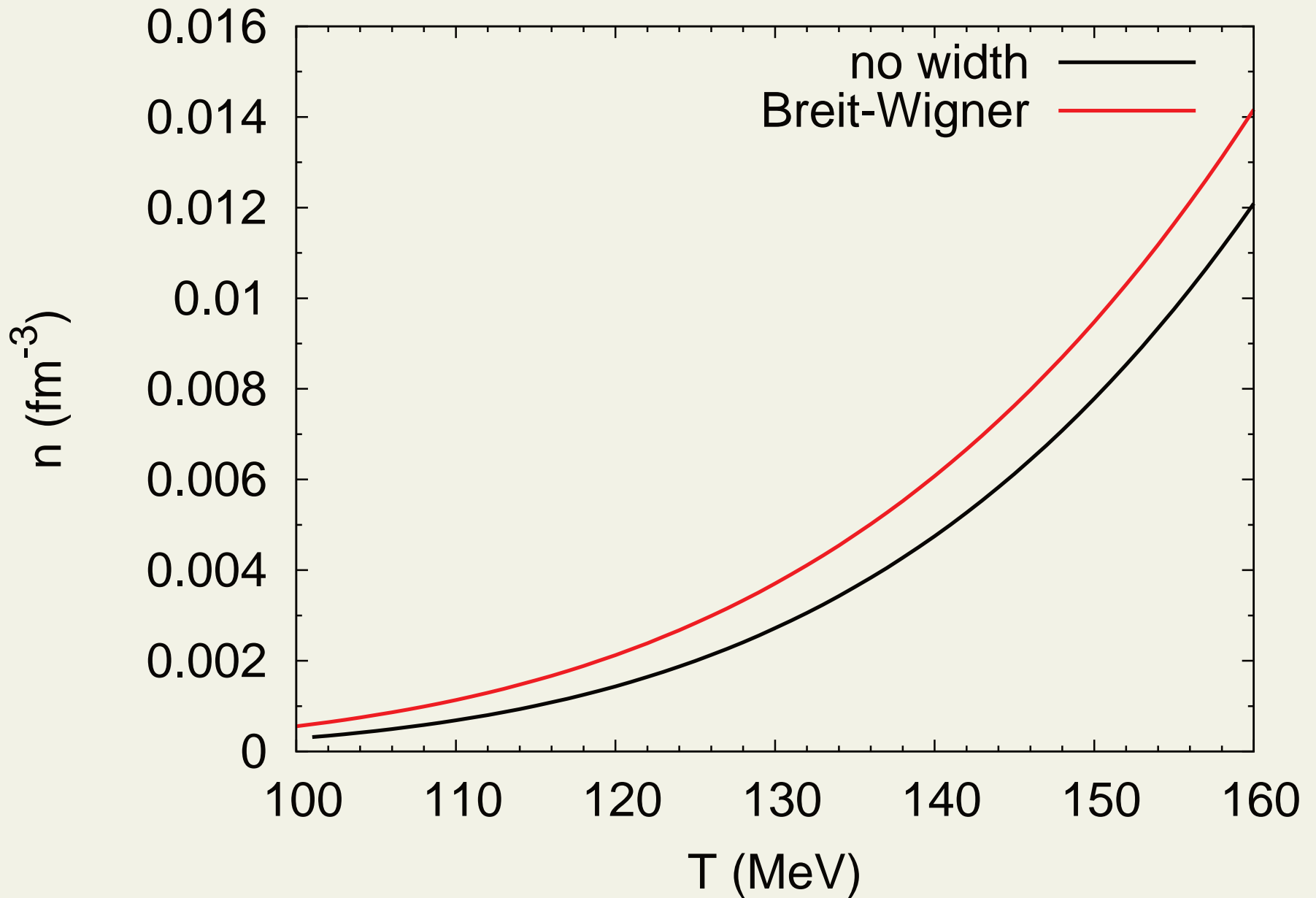
with normalisation

$$N = \int_{m_0}^{\infty} dm^2 \frac{m_R \Gamma}{(m^2 - m_R^2)^2 + m_R^2 \Gamma^2},$$

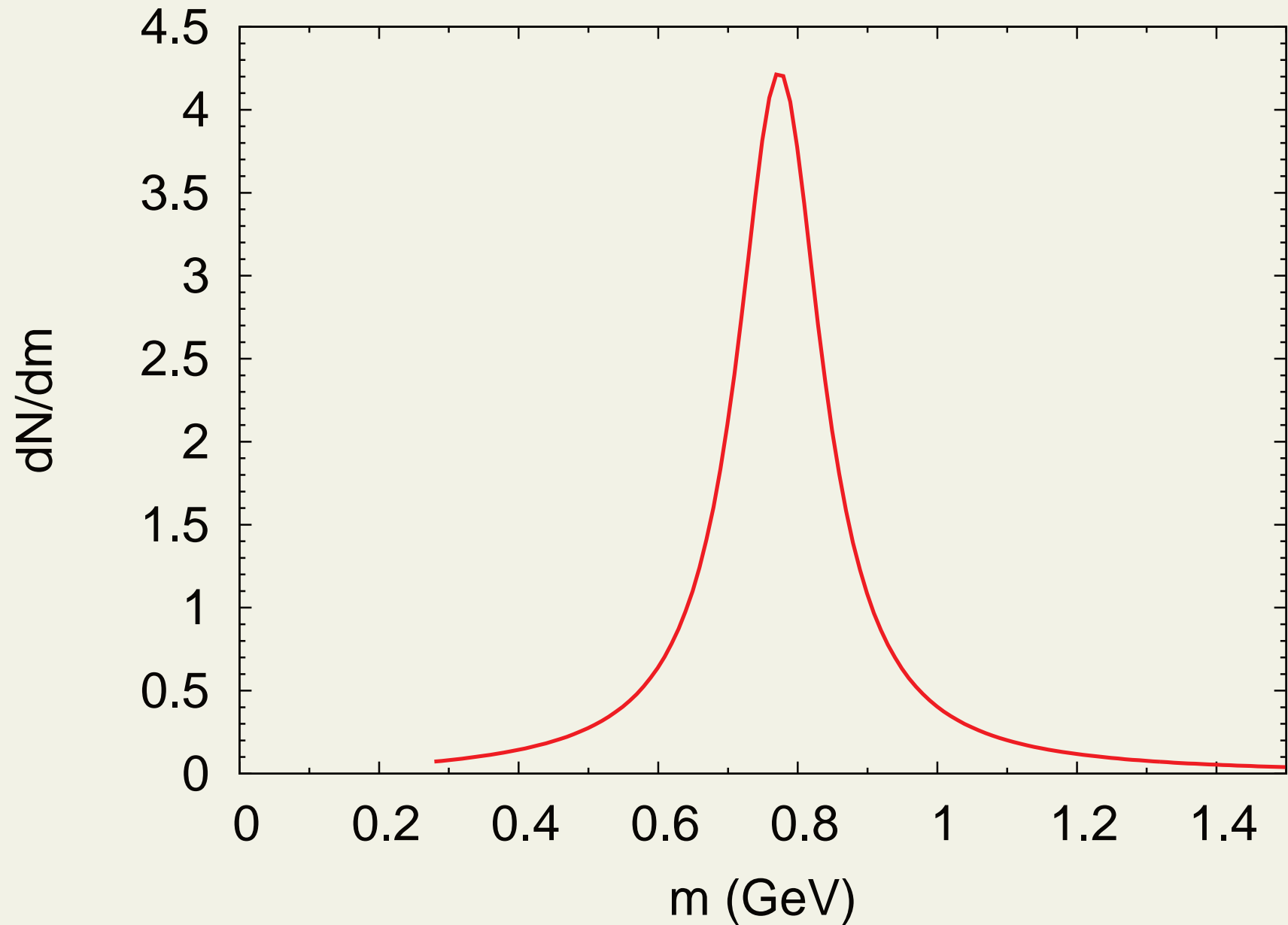
**For  $\rho^0$   $m_R = 775.26 \text{ MeV}$  and  $\Gamma = 147.8 \text{ MeV}$**



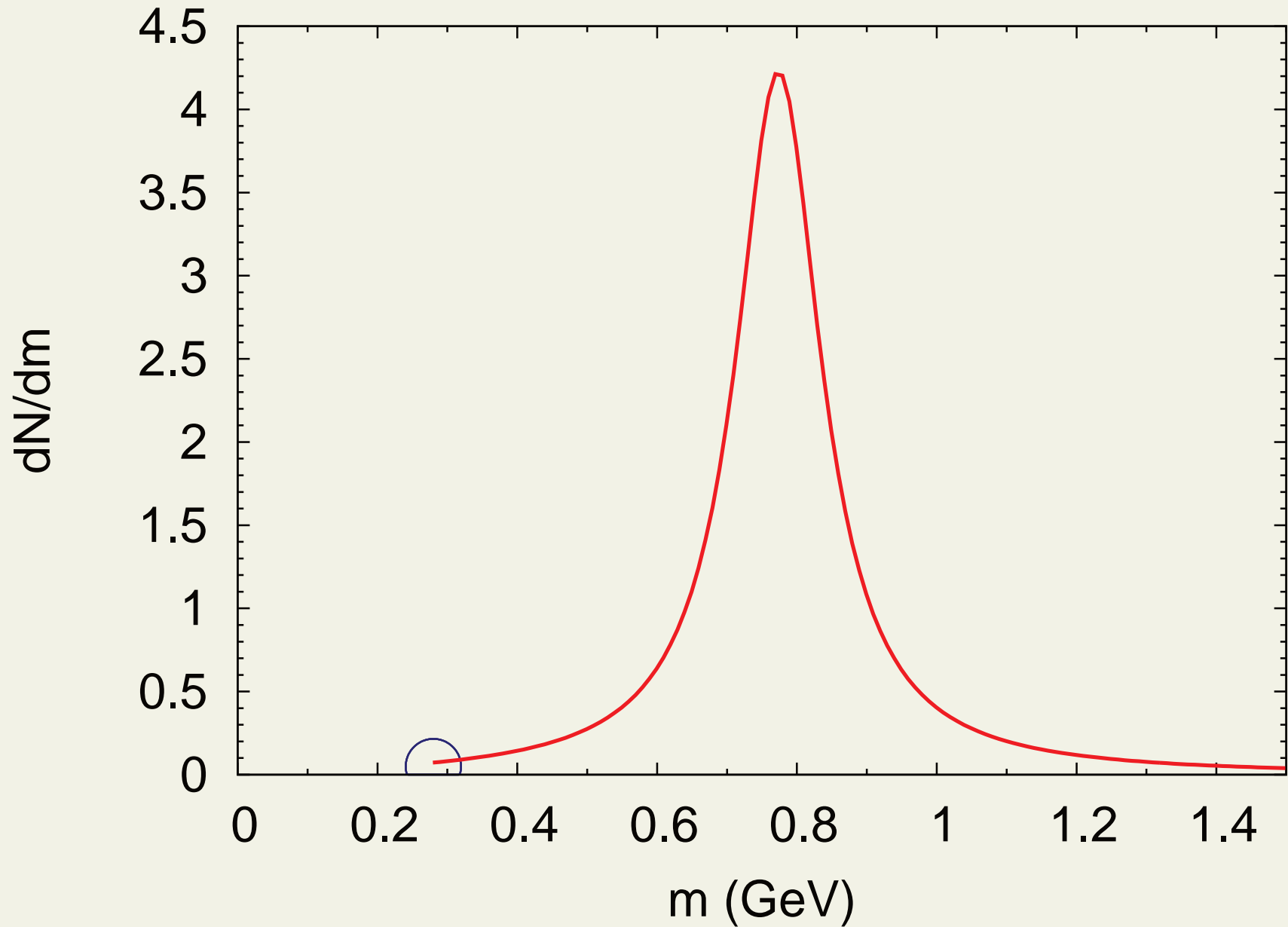
# $\rho$ -density



# Breit-Wigner



# Breit-Wigner



# Mass dependent width

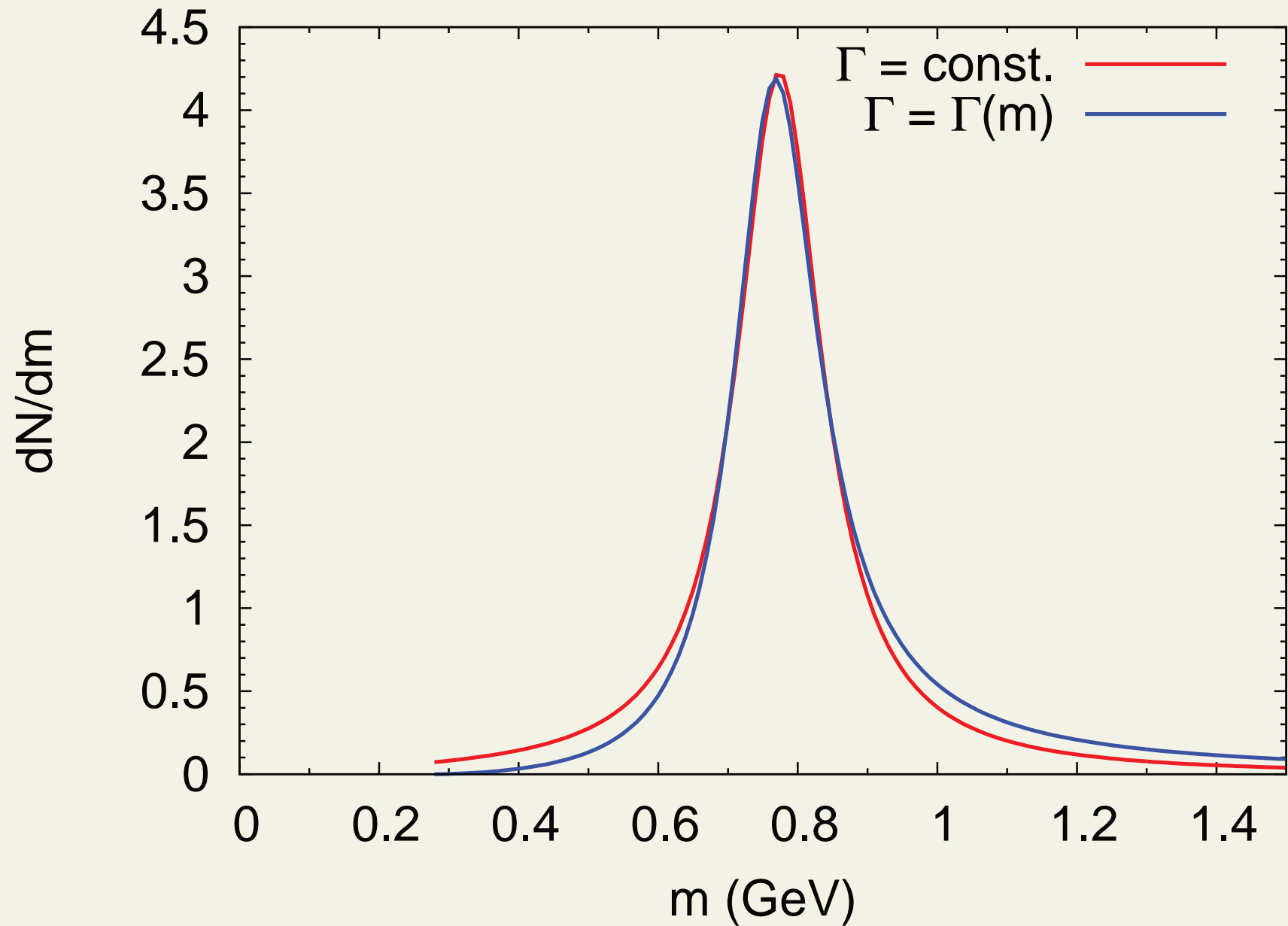
$$\frac{d\rho}{dm^2} = \frac{1}{N} \frac{m_R \Gamma(m)}{(m^2 - m_R^2)^2 + m_R^2 \Gamma(m)^2},$$

with width

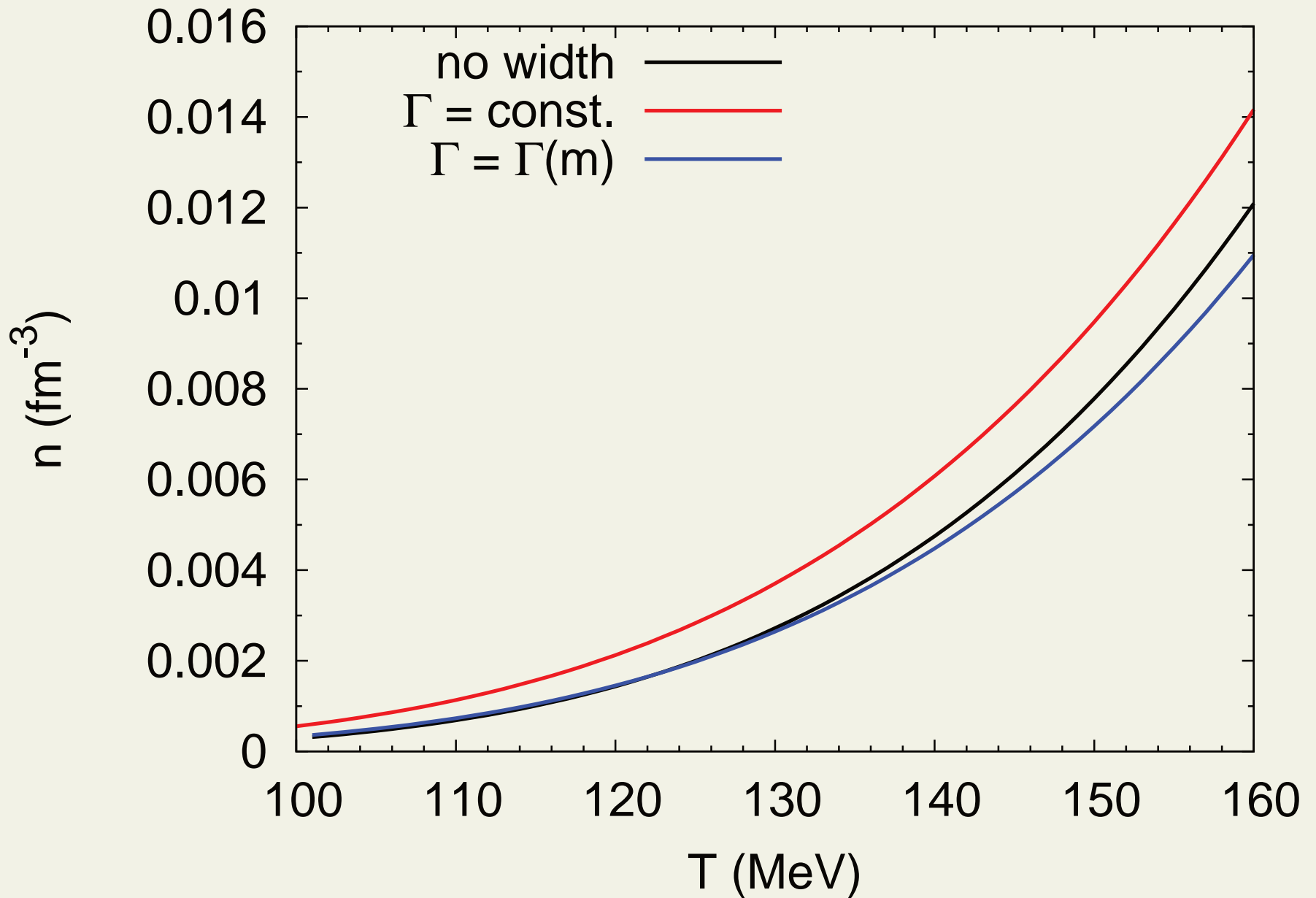
$$\Gamma(m) = \frac{1}{2} \frac{p_{\text{CMS}}^3 r_0^2}{1 + p_{\text{CMS}}^2 r_0^2}$$

where  $r_0 = 6.3 \text{ GeV}^{-1}$

# Breit-Wigner



# $\rho$ -density



# relativistic Breit-Wigner

$$\frac{d\rho}{dm^2} = \frac{1}{N} \frac{m_R \Gamma(m)}{(m^2 - m_R^2)^2 + m_R^2 \Gamma(m)^2}$$

**or:**

$$\frac{d\rho}{dm^2} = \frac{1}{N} \frac{m \Gamma(m)}{(m^2 - m_R^2)^2 + m^2 \Gamma(m)^2}$$

# relativistic Breit-Wigner

$$\frac{d\rho}{dm^2} = \frac{1}{N} \frac{m_R \Gamma(m)}{(m^2 - m_R^2)^2 + m_R^2 \Gamma(m)^2}$$

or:

$$\frac{d\rho}{dm^2} = \frac{1}{N} \frac{m \Gamma(m)}{(m^2 - m_R^2)^2 + m^2 \Gamma(m)^2}$$

**But if  $\Gamma(m) \propto m$  at large  $m$ ,**

$$N = \int_{m_0}^{\infty} dm^2 \frac{m \Gamma(m)}{(m^2 - m_R^2)^2 + m^2 \Gamma(m)^2} = \infty$$



## Particle Data Group about $\rho$ :

...the line shape does not correspond to a relativistic Breit-Wigner function...but requires some additional shape parameter

# Garbage in, garbage out



**Dashen-Ma-Berstein theorem:** If interactions mediated by *narrow* resonances, properties of interacting hadron gas are those of noninteracting hadron-resonance gas  $\Rightarrow$  **Hadron resonance gas model**

**Dashen-Ma-Berstein theorem:** If interactions mediated by *narrow* resonances, properties of interacting hadron gas are those of noninteracting hadron-resonance gas  $\Rightarrow$  **Hadron resonance gas model**

**Dashen-Ma-Berstein:** S-matrix formulation of statistical mechanics:

$\Rightarrow$  Second virial coefficient can be evaluated in terms of scattering phase shift (as far as interaction is manifested in elastic scattering)

$\Rightarrow$  relativistic Beth-Uhlenbeck form

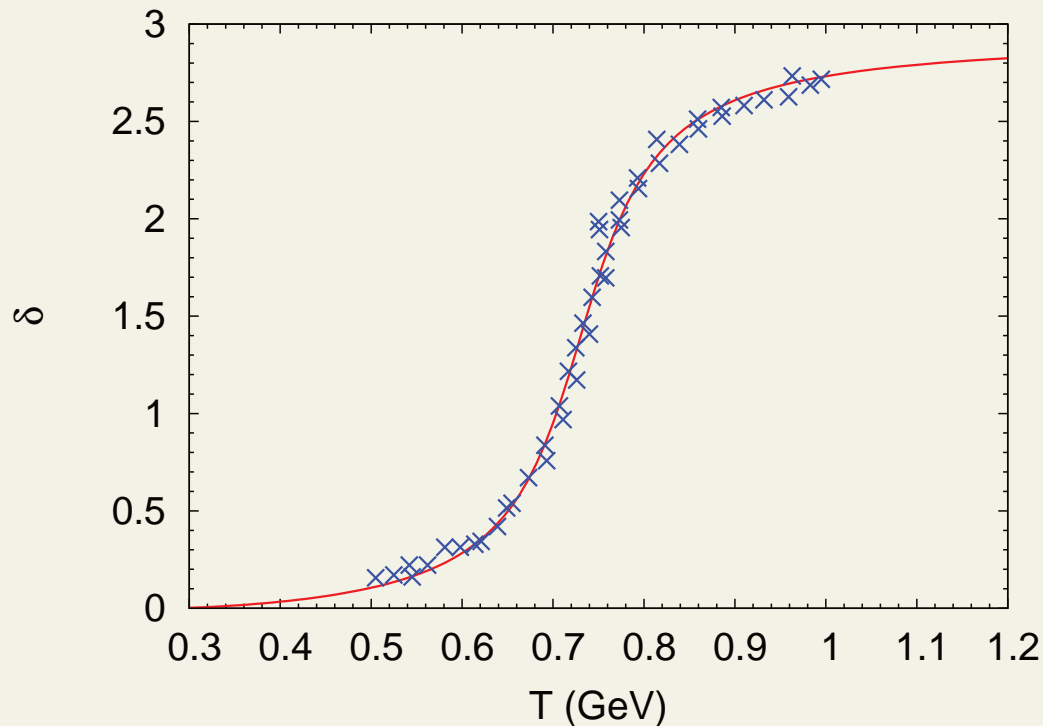
# Beth-Uhlenbeck

- effects of interactions expressed in terms of scattering phase shifts

$$n = \int d^3\mathbf{p} \int dm \frac{d\rho}{dm} f(p, m)$$

with

$$\frac{d\rho}{dm} = \frac{1}{\pi} \frac{d\delta}{dm}$$

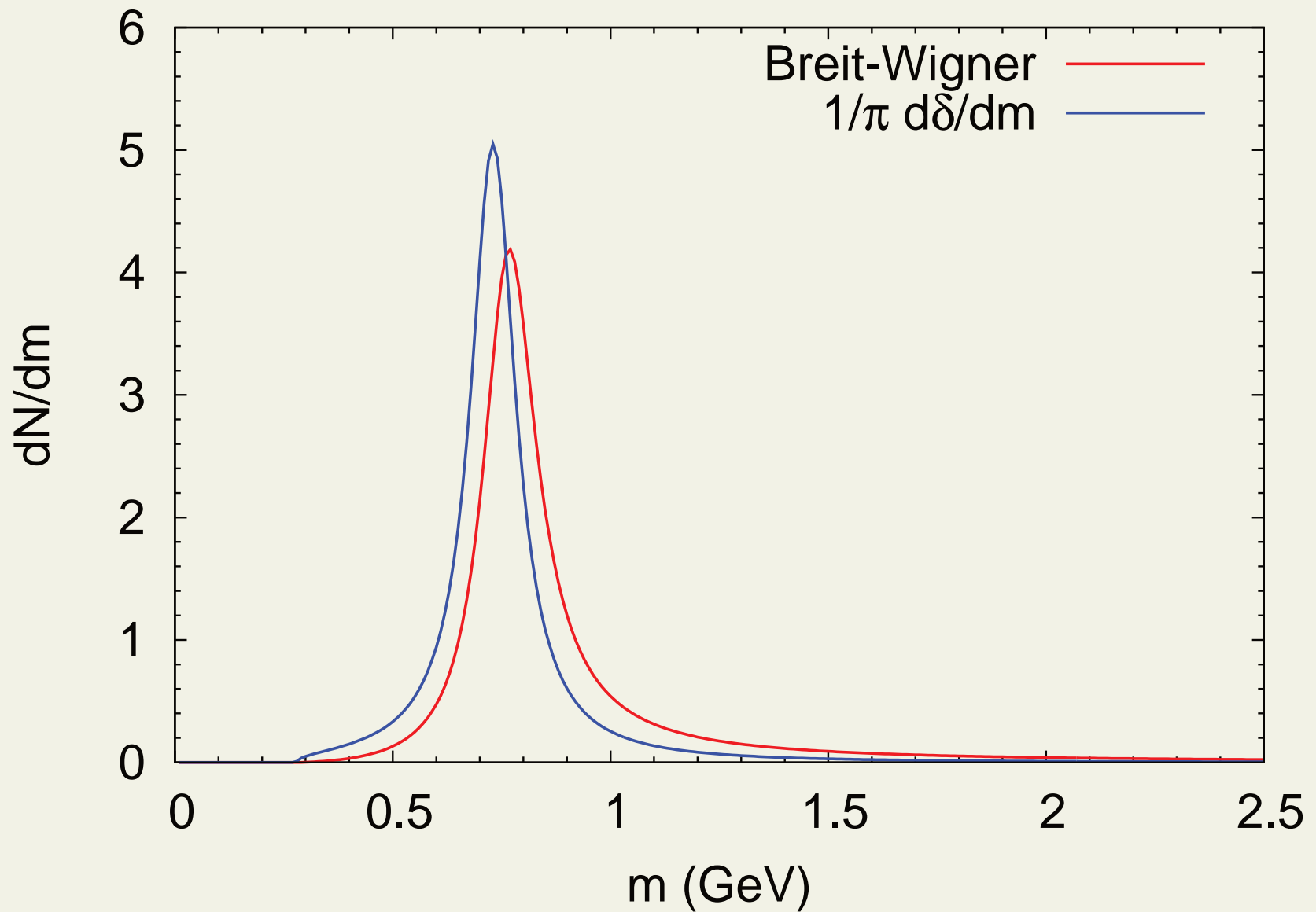


$$\delta(m) = \arctan \frac{-2\alpha}{3} \frac{p_{\text{CMS}}^3}{m(m^2 - \hat{m}_r^2)}$$

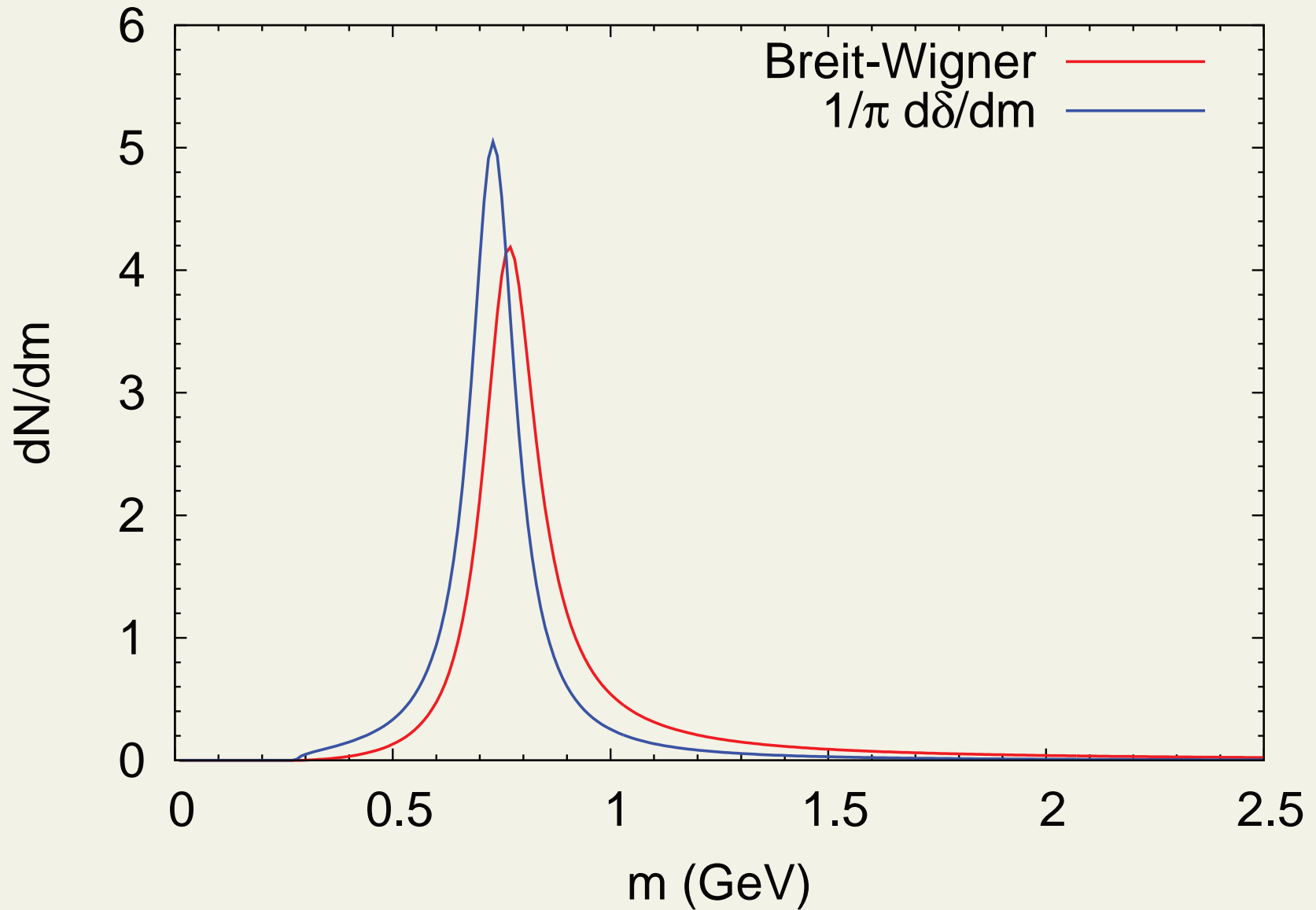
where

$$\alpha = 2.64526$$

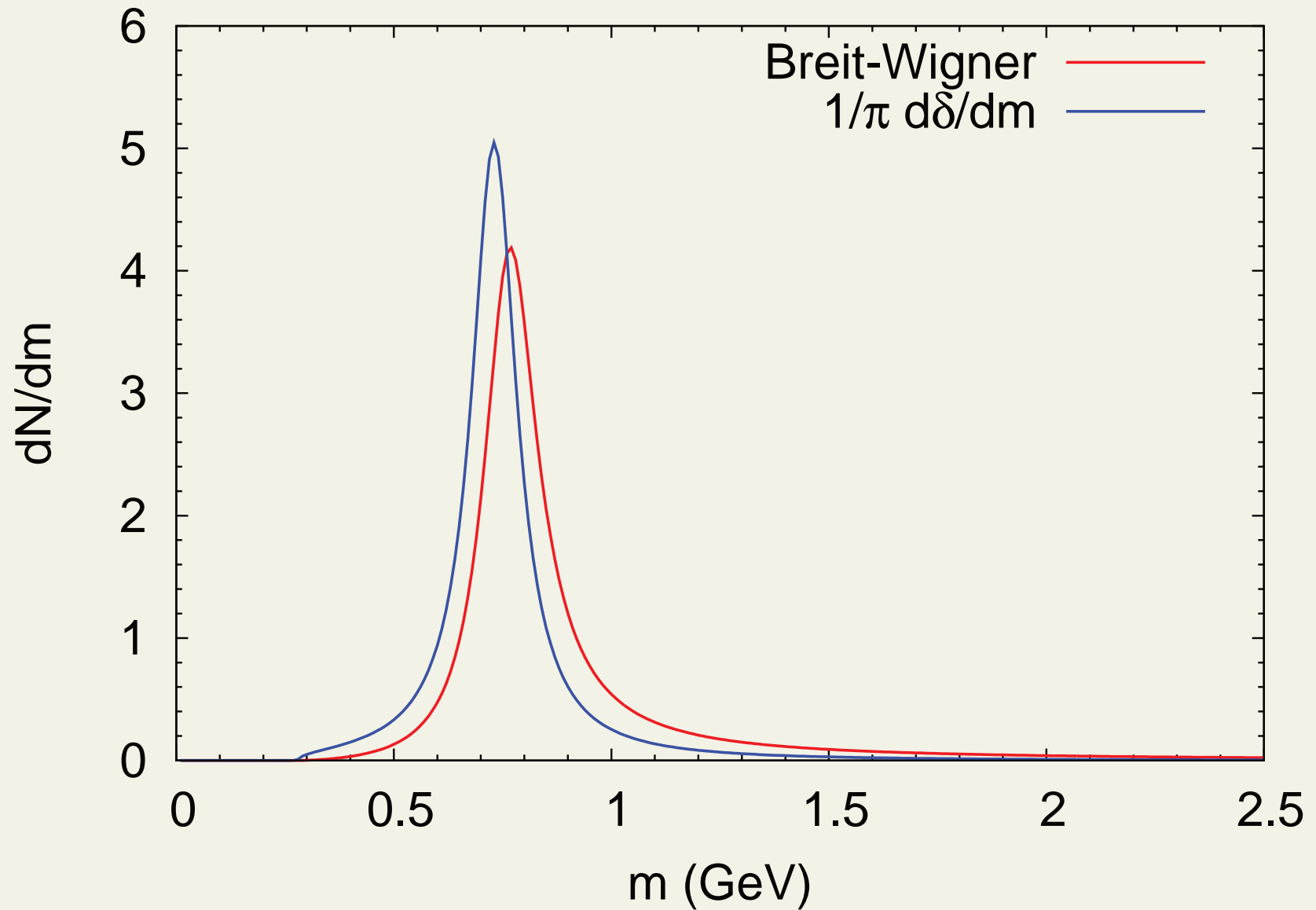
$$\hat{m}_r = 0.741395$$



**Note:** 
$$\frac{dn}{dm} = \frac{1}{\pi} \frac{d\delta}{dm} = \frac{2\hat{m}_r^2 \Gamma_\delta(m)}{(m^2 - \hat{m}_r^2)^2 + m^2 \Gamma_\delta(m)^2} + \dots$$



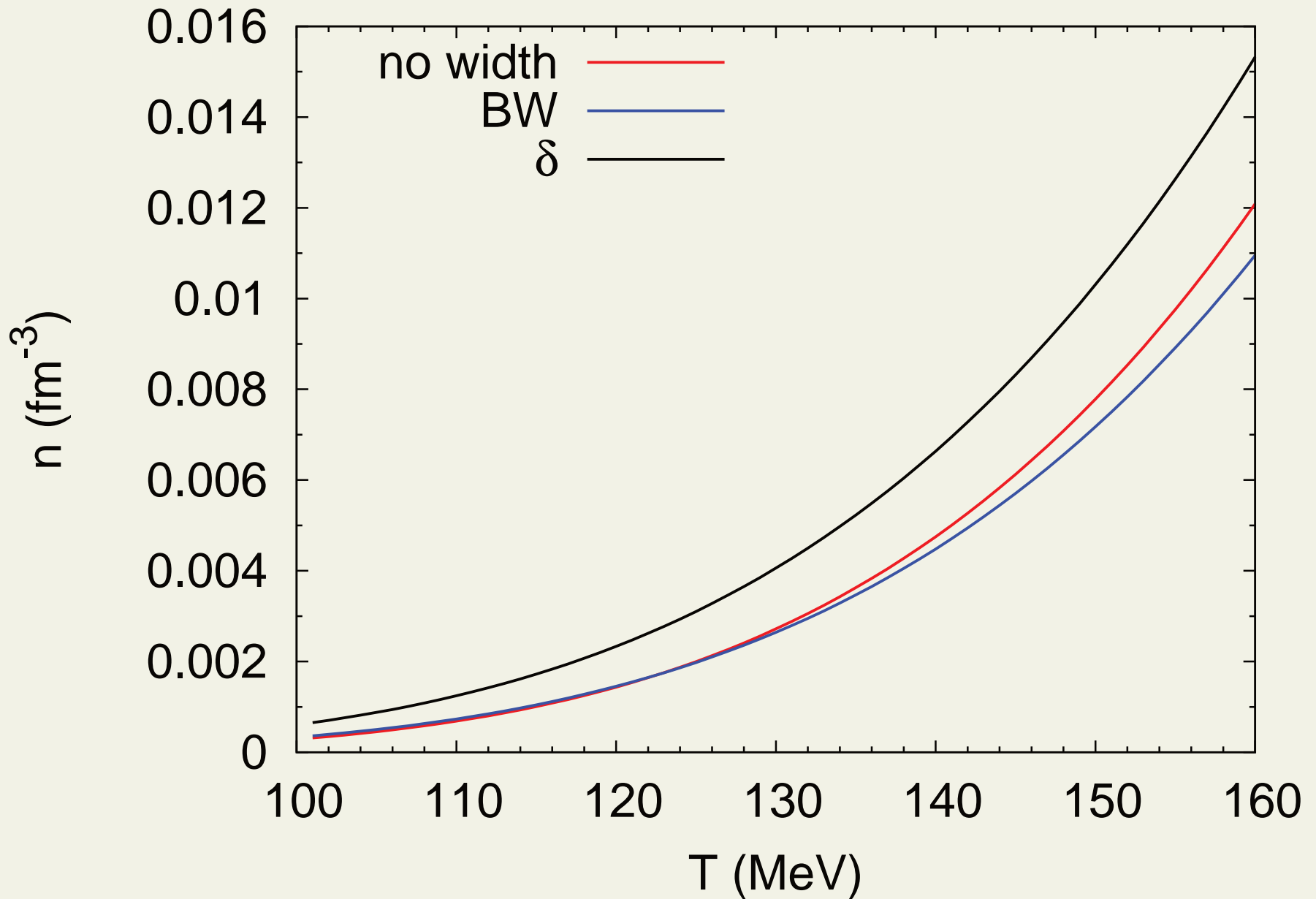
**Note:** 
$$\frac{dn}{dm} = \frac{1}{\pi} \frac{d\delta}{dm} = \frac{2\hat{m}_r^2 \Gamma_\delta(m)}{(m^2 - \hat{m}_r^2)^2 + m^2 \Gamma_\delta(m)^2} + \frac{\Gamma_\delta(m)}{p_{\text{CMS}}^2} \frac{3m_\pi^2 (\hat{m}_r^2 - m^2)}{(m^2 - \hat{m}_r^2)^2 + m^2 \Gamma_\delta(m)^2}$$



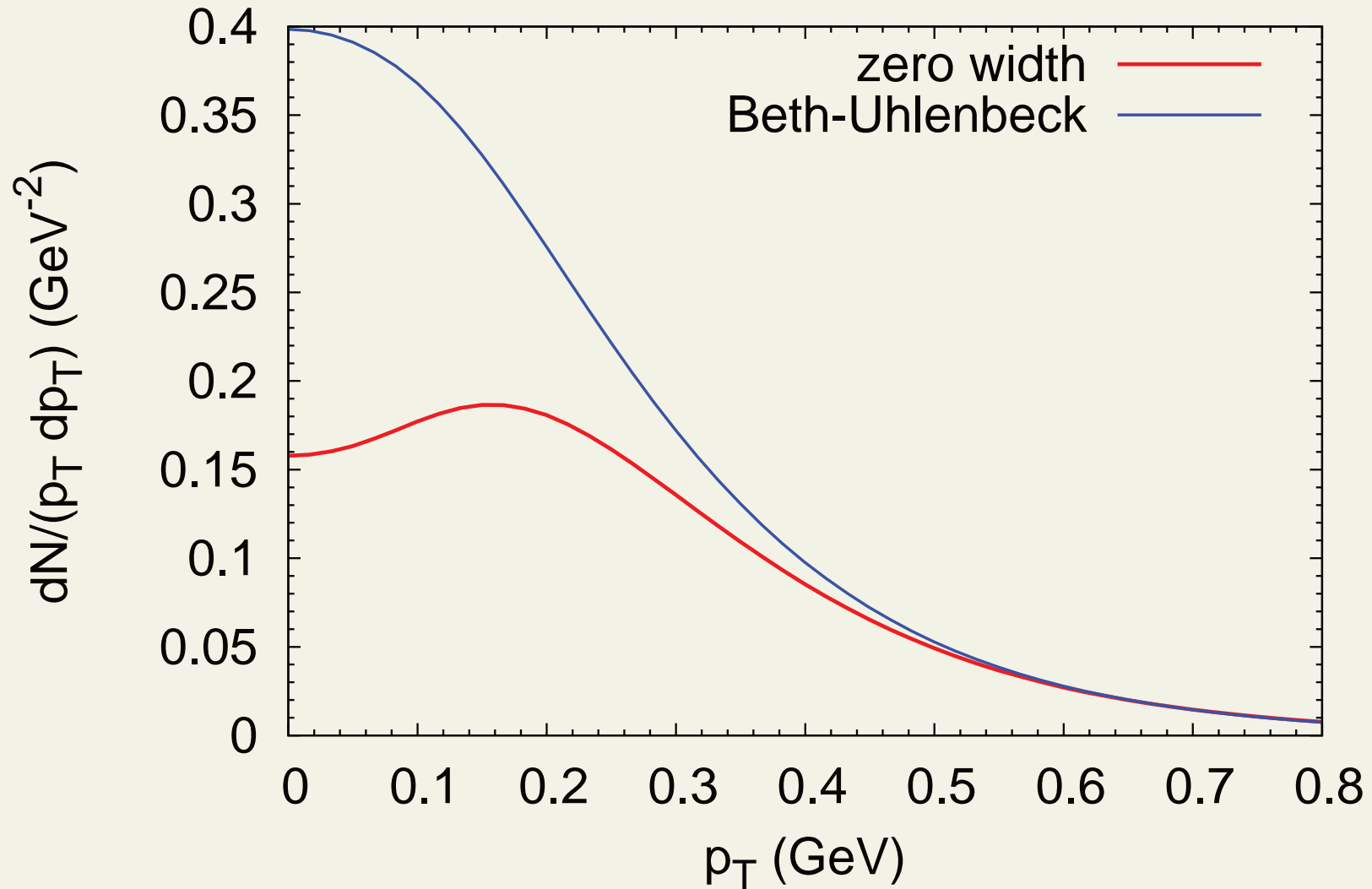
**Note:**  $m_\rho = 775.26 \text{ MeV}$  vs.  $\hat{m}_r = 741.395 \text{ MeV}$



# $\rho$ -density

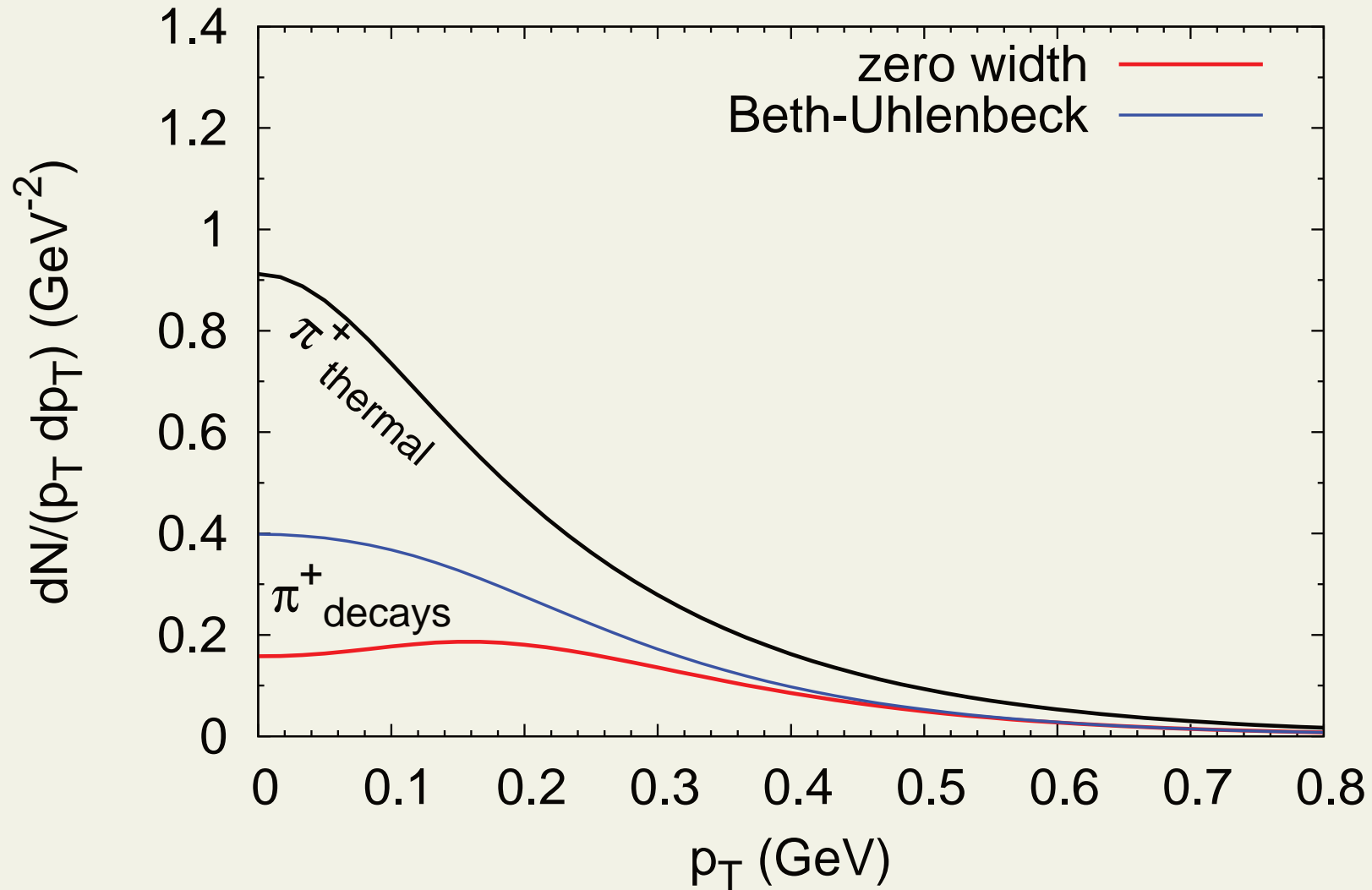


# Pions from $\rho$ decays



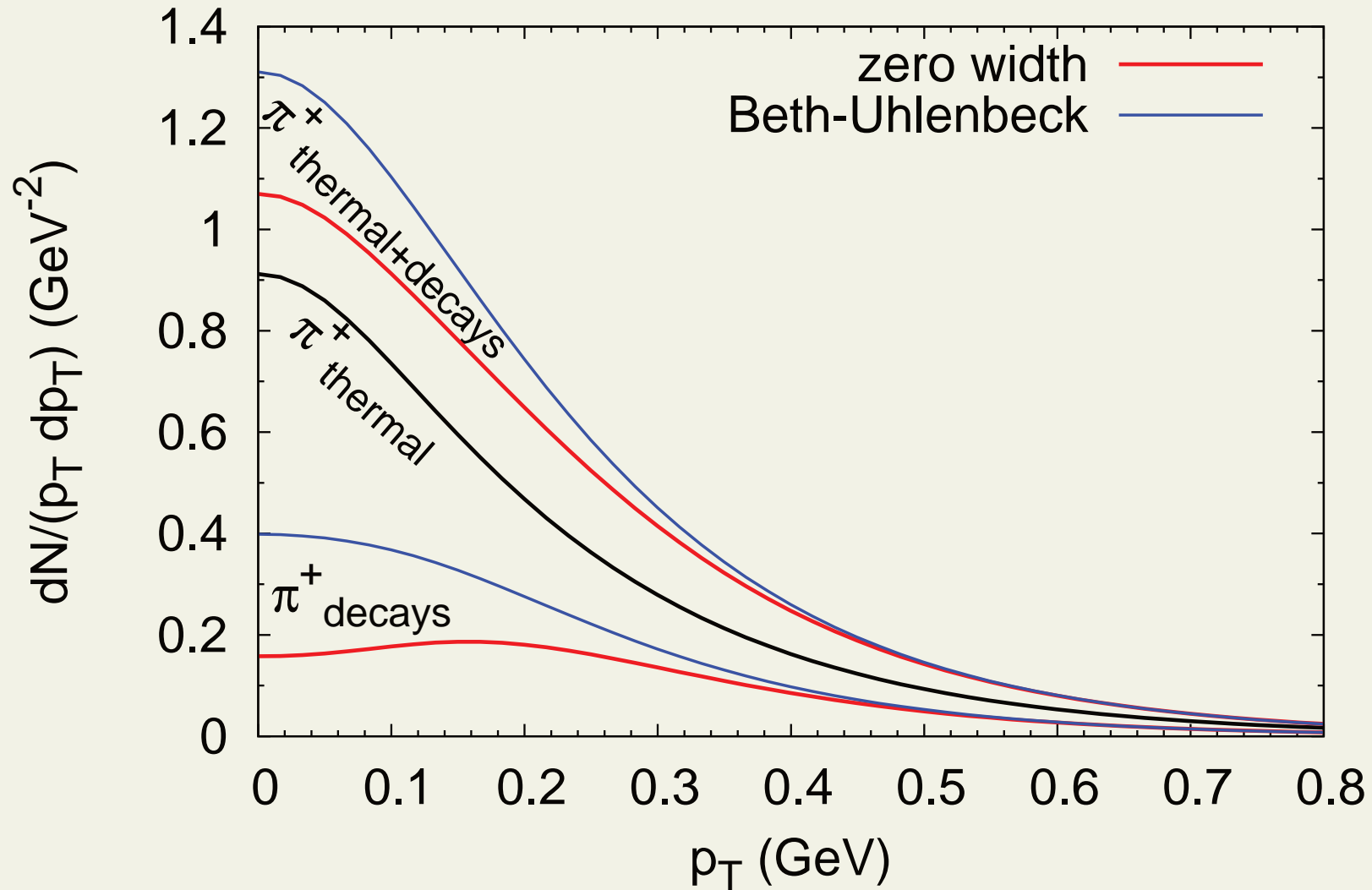
- static source,  $T = 155 \text{ MeV}$

# Thermal pions + pions from $\rho$ decays



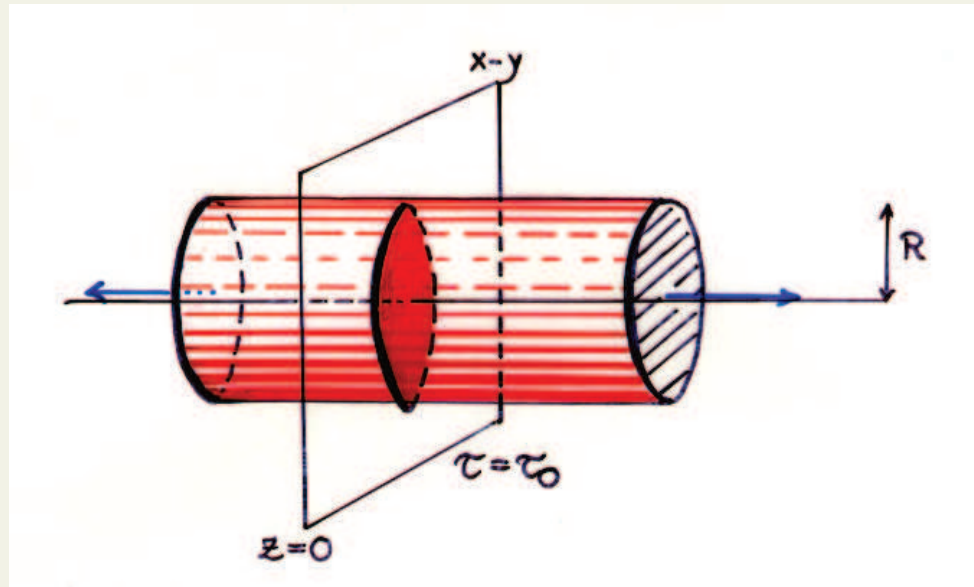
- static source,  $T = 155 \text{ MeV}$

# Thermal pions + pions from $\rho$ decays



- static source,  $T = 155$  MeV

# blast-wave parametrisation

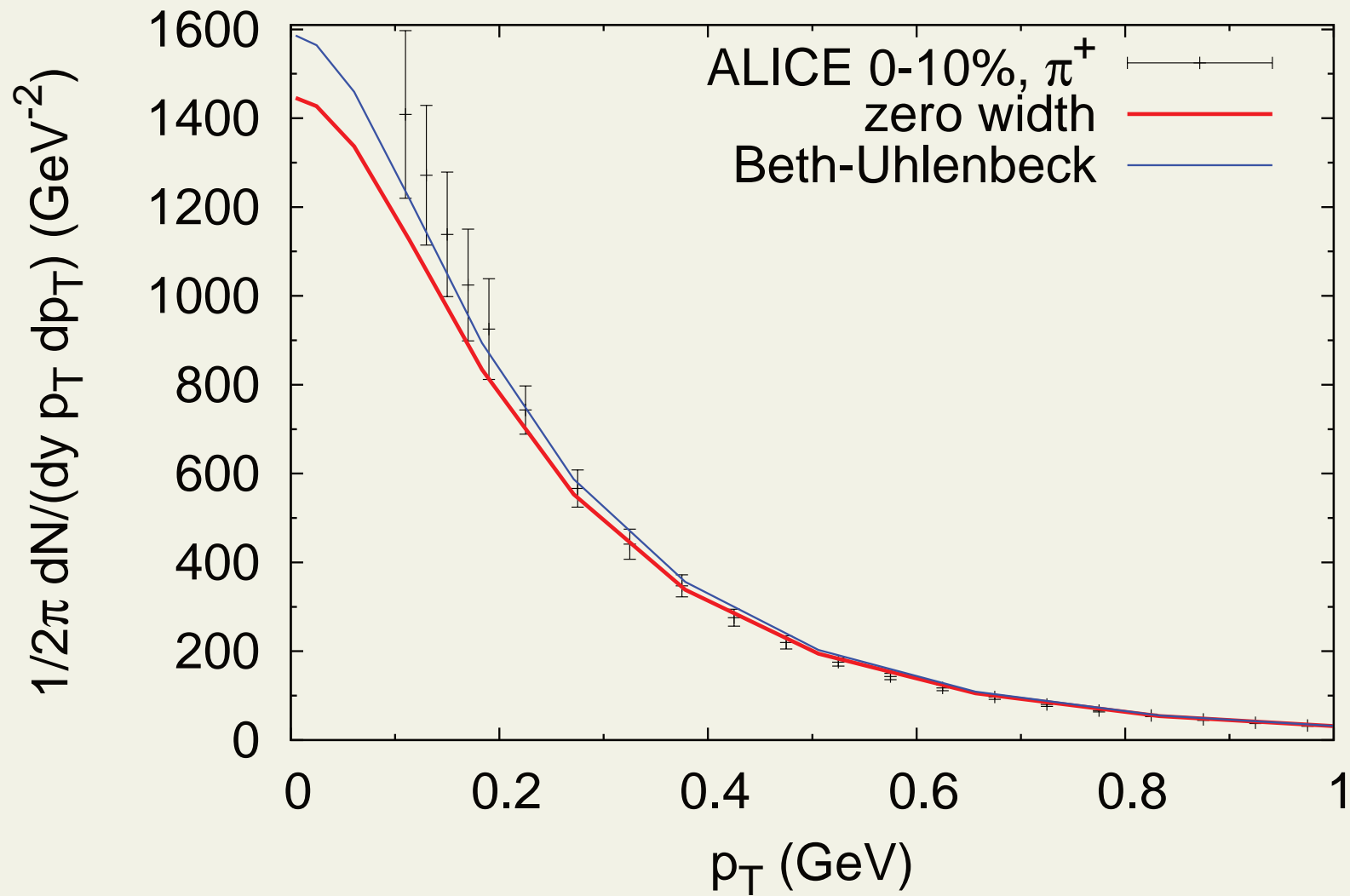


- boost invariant & cylidricly symmetric
- decoupling at constant  $\tau$ , i.e. volume emission
- transverse velocity  $v = v(r)$

$$E \frac{dN}{dp^3} = \frac{g\tau m_T}{2\pi^2} \int_0^R r dr \int_{m_{th}}^{\infty} dm \frac{d\rho}{dm} \sum_{n=1}^{\infty} (\mp 1)^{n+1} I_0 \left( n \frac{p_T \gamma_r(r) v_r(r)}{T} \right) K_1 \left( n \frac{m_T \gamma_r(r)}{T} \right)$$

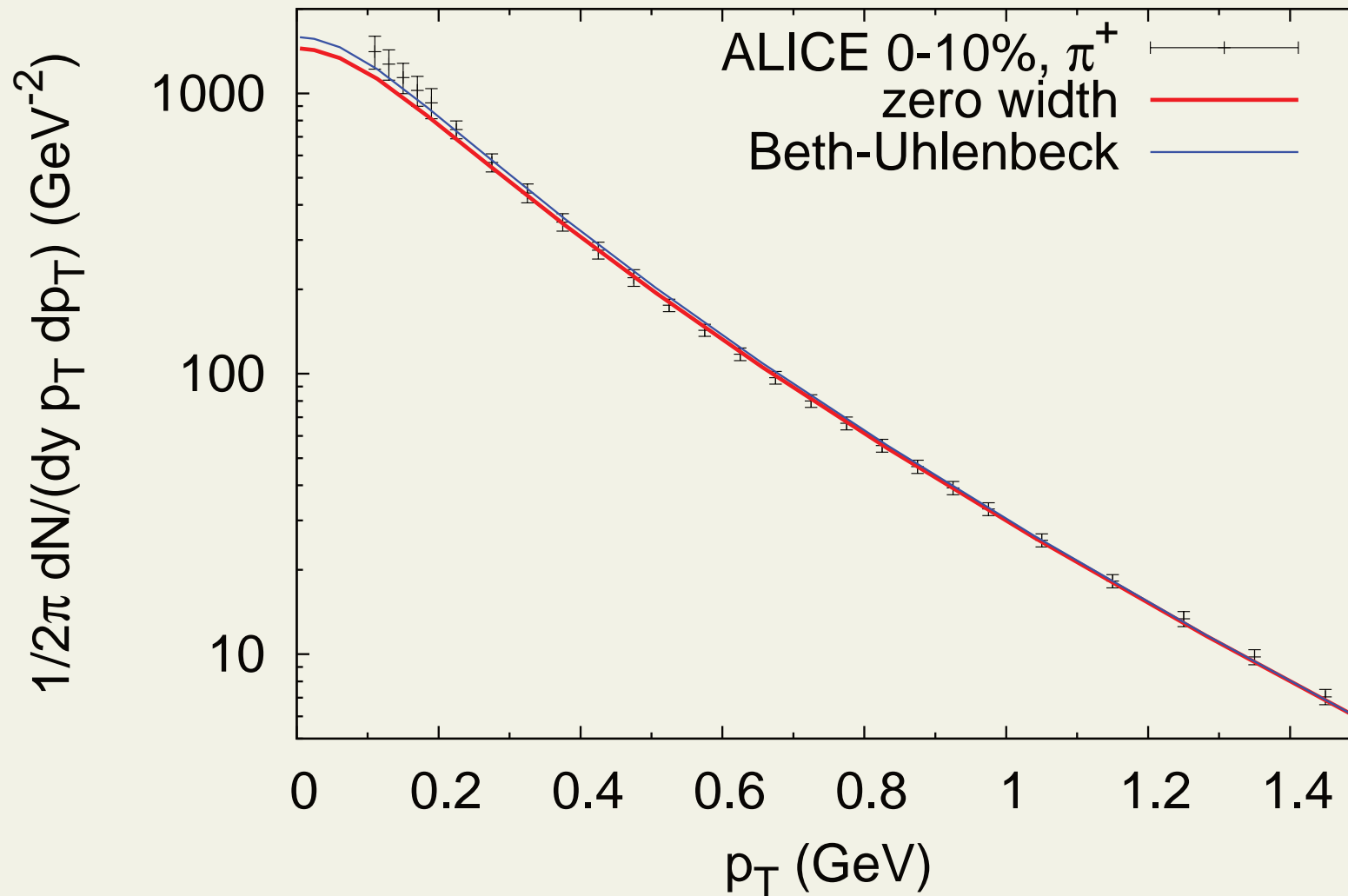
$$\tau = 13.7 \text{ fm}, \quad R = 10 \text{ fm}, \quad v_{max} = 0.78$$

# Pions from blast wave



- all resonances up to 2 GeV
- Beth-Uhlenbeck for rhos
- zero width for everything else

# Pions from blast wave



- all resonances up to 2 GeV
- Beth-Uhlenbeck for rhos
- zero width for everything else

# caveats

- so far only rho mesons



# caveats

- so far only rho mesons
- Beth-Uhlenbeck applicable to **elastic scatterings only!**

# caveats

- so far only rho mesons
- Beth-Uhlenbeck applicable to **elastic scatterings only!**
- $\rho$ ,  $K^*(892)$ ,  $f_0(980)$ ,  $\Delta(1232)$ ,  $K_0^*(1430)$ 
  - data exists

# caveats

- so far only rho mesons
- Beth-Uhlenbeck applicable to **elastic scatterings only!**
- $\rho$ ,  $K^*(892)$ ,  $f_0(980)$ ,  $\Delta(1232)$ ,  $K_0^*(1430)$ 
  - data exists
- $\Lambda(1405)$ ,  $\Xi(1530)$  applicable
  - no data

# caveats

- so far only rho mesons
- Beth-Uhlenbeck applicable to **elastic scatterings only!**
- $\rho$ ,  $K^*(892)$ ,  $f_0(980)$ ,  $\Delta(1232)$ ,  $K_0^*(1430)$ 
  - data exists
- $\Lambda(1405)$ ,  $\Xi(1530)$  applicable
  - no data
- **and everything else?**

# Summary

- **So is there anomaly. . . ?**

# Summary

- So is there anomaly. . . ?
  - **Probably not**

# Summary

- So is there anomaly. . . ?
  - **Probably not**
- **Resonance widths affect yields and distributions**

# Summary

- So is there anomaly. . . ?
  - **Probably not**
- **Resonance widths affect yields and distributions**
- **Better treatment of resonances needed**