

# The Constant-Sound-Speed parameterization of the quark matter EoS

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Alford, Han, Prakash, [arXiv:1302.4732](#)

Alford, Burgio, Han, Taranto, Zappalà, [arXiv:1501.07902](#)

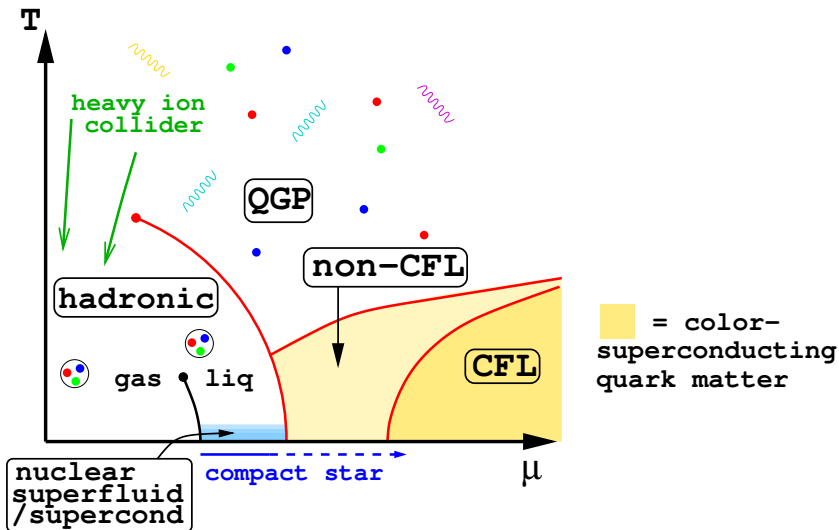
Ranea-Sandoval, Han, Orsaria, Contrera, Weber, Alford,  
[arXiv:1512.09183](#)



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# Schematic QCD phase diagram



M. Alford, K. Rajagopal, T. Schäfer, A. Schmitt, [arXiv:0709.4635](https://arxiv.org/abs/0709.4635) (RMP review)

A. Schmitt, [arXiv:1001.3294](https://arxiv.org/abs/1001.3294) (Springer Lecture Notes)

# Signatures of quark matter in compact stars

Observable ← Microphysical properties (and neutron star structure) ← Phases of dense matter

	Property	Nuclear phase	Quark phase
mass, radius	eqn of state $\varepsilon(\rho)$	known up to $\sim n_{\text{sat}}$	unknown; many models
spindown (spin freq, age)	bulk viscosity shear viscosity	Depends on phase:	Depends on phase:
cooling (temp, age)	heat capacity	$n p e$	unpaired
	neutrino emissivity	$n p e, \mu$	CFL
	thermal cond.	$n p e, \Lambda, \Sigma^-$	CFL- $K^0$
glitches (superfluid, crystal)	shear modulus	$n$ superfluid	2SC
	vortex pinning	$p$ supercond	CSL
	energy	$\pi$ condensate	LOFF
		$K$ condensate	1SC
			...

# Constraining QM EoS by observing $M(R)$

There is lots of literature about specific models of quark matter, e.g.

- ▶ MIT Bag Model; (Alford, Braby, Paris, Reddy, nucl-th/0411016)
- ▶ NJL models; (Paoli, Menezes, arXiv:1009.2906; Bonanno, Sedrakian, arXiv:1108.0559)
- ▶ PNJL models (Blaschke et. al, arXiv:1302.6275; Orsaria et. al.; arXiv:1212.4213)
- ▶ hadron-quark  $NL\sigma$  model (Negreiros et. al., arXiv:1006.0380)
- ▶ 2-loop perturbation theory (Kurkela et. al., arXiv:1006.4062)
- ▶ MIT bag, NJL, CDM, FCM, DSM (Burgio et. al., arXiv:1301.4060)

We need a **model-independent parameterization** of the quark matter EoS:

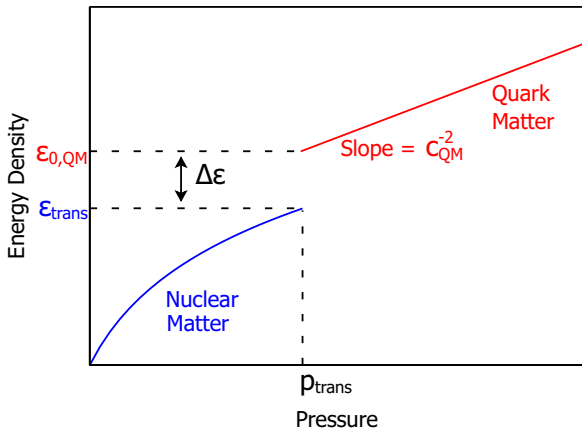
- ▶ framework for relating different models to each other
- ▶ observational constraints can be expressed in universal terms

# CSS: a fairly generic QM EoS

Model-independent parameterization with

- Sharp 1st-order transition
- Constant [density-indp] Speed of Sound (CSS)

$$\varepsilon(p) = \varepsilon_{\text{trans}} + \Delta\varepsilon + c_{\text{QM}}^{-2}(p - p_{\text{trans}})$$



QM EoS params:

$$p_{\text{trans}}/\varepsilon_{\text{trans}}$$

$$\Delta\varepsilon/\varepsilon_{\text{trans}}$$

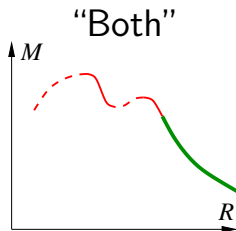
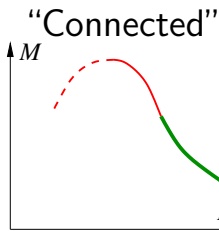
$$c_{\text{QM}}^2$$

# Hybrid star $M(R)$

Hybrid star branch in  $M(R)$  relation has 4 typical forms

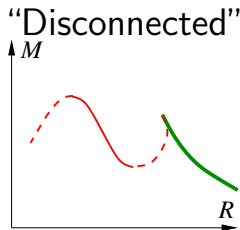
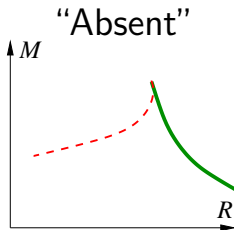
$$\Delta\varepsilon < \Delta\varepsilon_{\text{crit}}$$

small energy density jump at phase transition



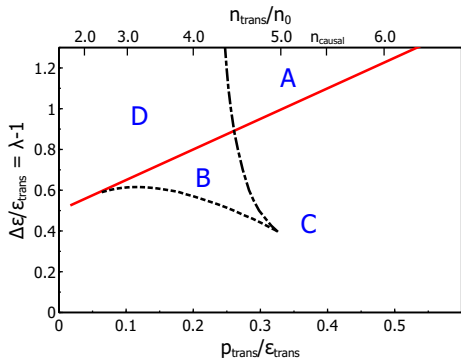
$$\Delta\varepsilon > \Delta\varepsilon_{\text{crit}}$$

large energy density jump at phase transition

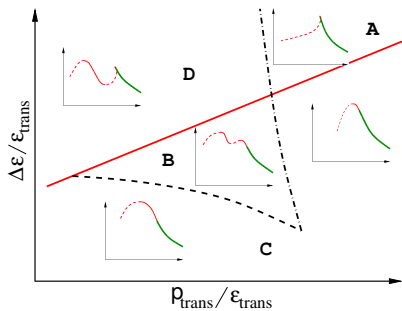


# CSS “Phase diagram” of hybrid star $M(R)$

Soft NM + CSS ( $c_{\text{QM}}^2 = 1$ )



Schematic



Above the red line ( $\Delta\epsilon > \Delta\epsilon_{\text{crit}}$ ),  
connected branch disappears

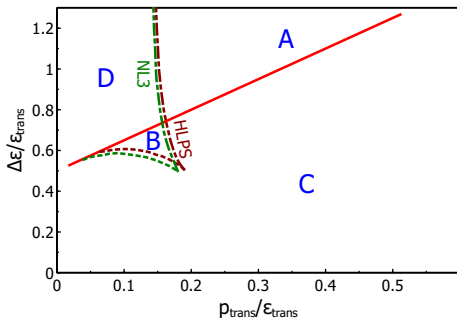
(Seidov, 1971; Schaeffer, Zduunik, Haensel, 1983; Lindblom, gr-qc/9802072)

Disconnected branch exists in regions D and B.

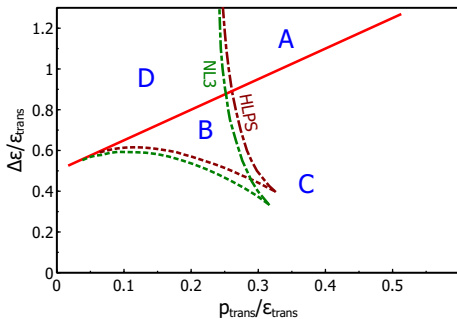
$$\frac{\Delta\epsilon_{\text{crit}}}{\epsilon_{\text{trans}}} = \frac{1}{2} + \frac{3}{2} \frac{\rho_{\text{trans}}}{\epsilon_{\text{trans}}}$$

# Sensitivity to NM EoS and $c_{QM}^2$

$$c_{QM}^2 = 1/3$$



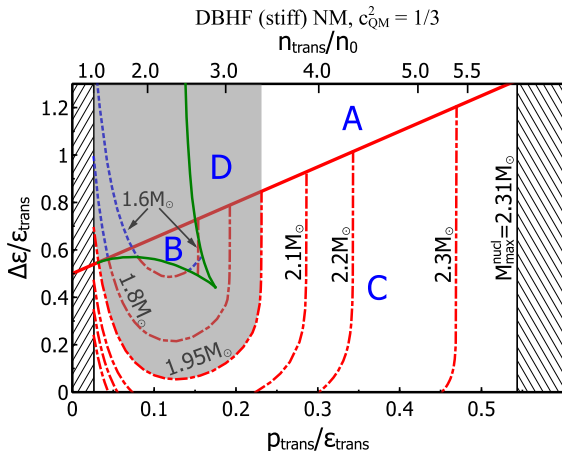
$$c_{QM}^2 = 1$$



- NM EoS (HLPS=soft, NL3=hard) does not make much difference.
- Higher  $c_{QM}^2$  favors disconnected branch.



# Constraints on QM EoS from $M_{\max}$



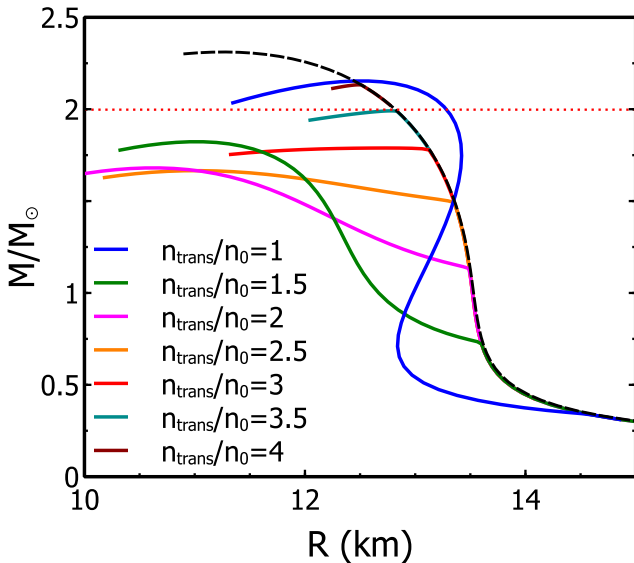
- Increasing  $\Delta\varepsilon$  reduces  $M_{\max}$
- Increasing  $p_{\text{trans}}$  at first reduces then increases  $M_{\max}$

2  $M_{\odot}$  observation allows two scenarios:

- high  $p_{\text{trans}}$ : very small connected branch
- low  $p_{\text{trans}}$ : modest  $\Delta\varepsilon$ , no disconnected branch.

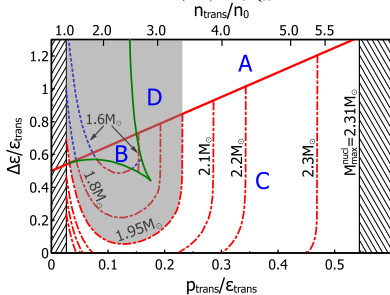
# Low $\rho_{\text{trans}}$ and high $\rho_{\text{trans}}$ windows

DBHF (stiff) NM,  $c_{\text{QM}}^2 = 1/3$ ,  $\Delta\varepsilon/\varepsilon_{\text{trans}} = 0.4$

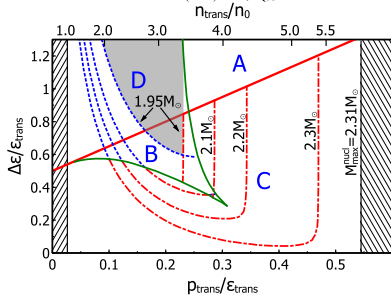


# Constraints on QM EoS from $M_{\max}$

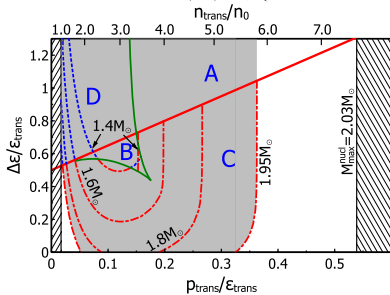
DBHF (stiff) NM,  $c_{\text{QM}}^2 = 1/3$



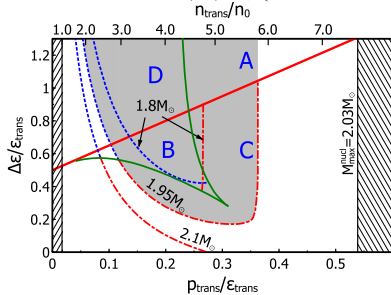
DBHF (stiff) NM,  $c_{\text{QM}}^2 = 1$



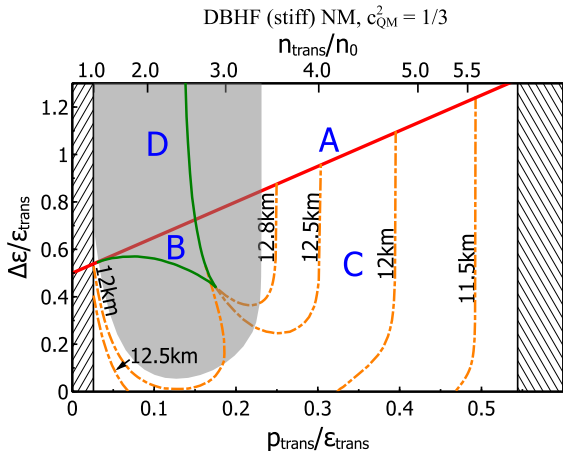
BHF (soft) NM,  $c_{\text{QM}}^2 = 1/3$



BHF (soft) NM,  $c_{\text{QM}}^2 = 1$



# Radius of heaviest star $R_{\max M}$



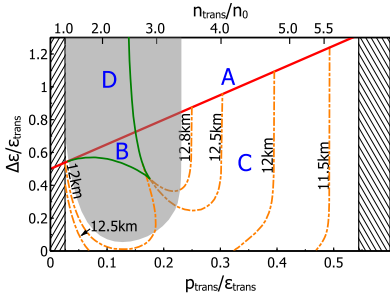
Heaviest star is typically the smallest, so lower limit on  $R_{\max M}$  is the minimum radius of compact stars.

High  $p_{\text{trans}}$ : very short connected hybrid branch, radius like that of heaviest hadronic star.

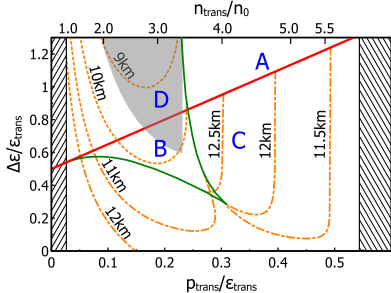
Low  $p_{\text{trans}}$ : need to zoom in.

# Constraints on QM EoS from $R_{\text{max}M}$

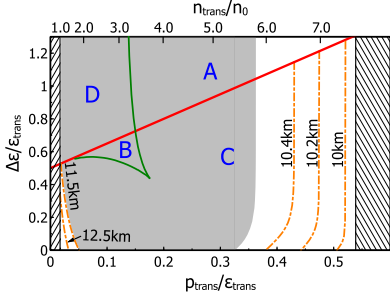
DBHF (stiff) NM,  $c_{\text{QM}}^2 = 1/3$



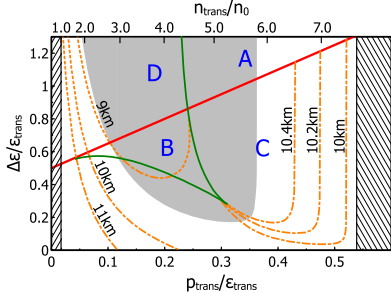
DBHF (stiff) NM,  $c_{\text{QM}}^2 = 1$



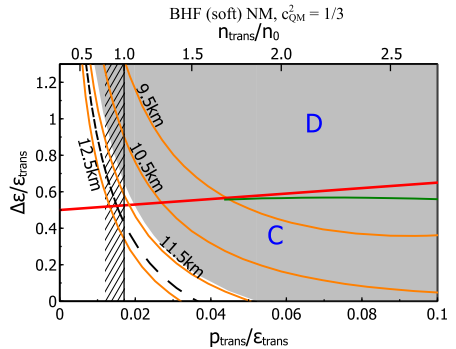
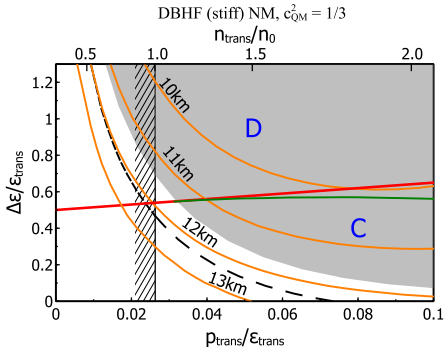
BHF (soft) NM,  $c_{\text{QM}}^2 = 1/3$



BHF (soft) NM,  $c_{\text{QM}}^2 = 1$

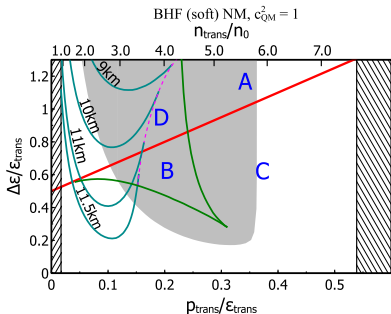
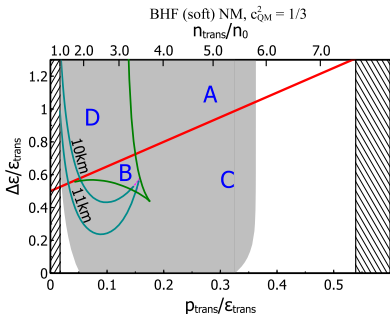
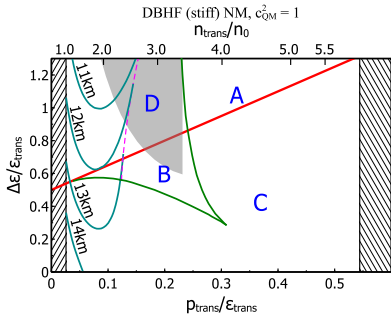
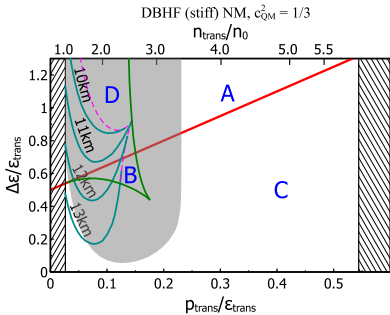


# Focus on low $\rho_{\text{trans}}$ and $c_{\text{QM}}^2 = 1/3$

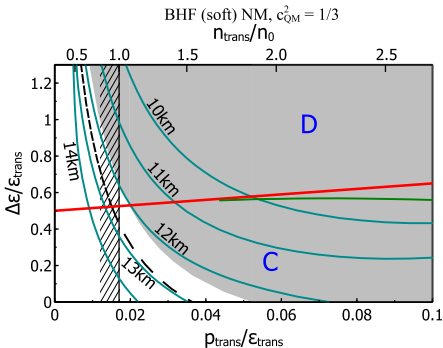
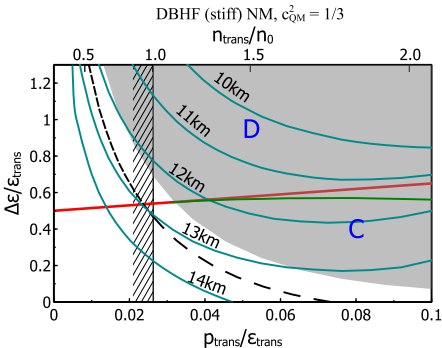


- ▶  $R_{\text{max}M}$  contours closely follow mass contours
- ▶  $M_{\text{max}} > 1.95 M_{\odot}$  requires  $R > 11.25$  km
- ▶ dashed line is  $M_{\text{max}} = 2.1 M_{\odot}$ , requires  $R > 12.1$  km
- ▶ Observation of a smaller star  $\Rightarrow$  high transition pressure or  $c_{\text{QM}}^2 > 1/3$

# Constraints on QM EoS from $R_{1.4M_{\odot}}$



# Low transition pressure and $R_{1.4 M_{\odot}}$



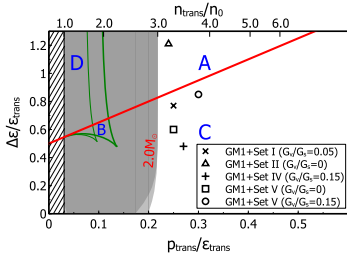
- ▶  $R_{1.4 M_{\odot}}$  contours roughly follow mass contours
- ▶  $M_{\text{max}} > 1.95 M_{\odot}$  requires  $R_{1.4 M_{\odot}} > 12 \text{ km}$  ( $n_{\text{trans}} \approx n_0$ ), rising with  $n_{\text{trans}}$ .
- ▶ dashed line is  $M_{\text{max}} = 2.1 M_{\odot}$ , requires  $R_{1.4 M_{\odot}} > 12.7 \text{ km}$
- ▶ Observation of a smaller  $1.4 M_{\odot}$  star  $\Rightarrow c_{\text{QM}}^2 > 1/3$ .
- ▶ If  $\rho_{\text{trans}}$  is high then no hybrid stars have mass  $1.4 M_{\odot}$

compare Lattimer arXiv:1305.3510:  $R > 11 \text{ km}$ .

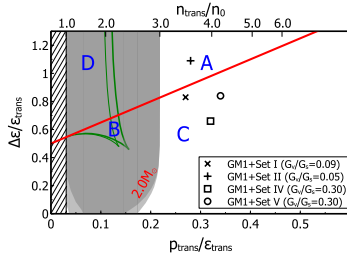


# NJL models in CSS space

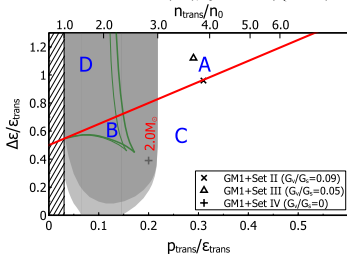
Hadronic: GM1 (soft); Quark: CSS ( $c_{QM}^2 = 0.2$ )



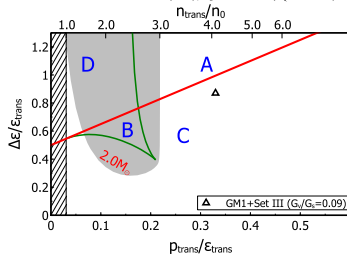
Hadronic: GM1 (soft); Quark: CSS ( $c_{QM}^2 = 0.25$ )



Hadronic: GM1 (soft); Quark: CSS ( $c_{QM}^2 = 0.3$ )



Hadronic: GM1 (soft); Quark: CSS ( $c_{QM}^2 = 0.46$ )



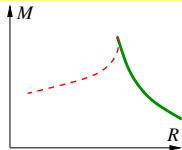
# Summary of CSS

- ▶ CSS (Constant Speed of Sound) is a generic parameterization of the EoS close to a sharp first-order transition to quark matter.
- ▶ Any specific model of quark matter with such a transition corresponds to particular values of the CSS parameters  $(p_{\text{trans}}/\varepsilon_{\text{trans}}, \Delta\varepsilon/\varepsilon_{\text{trans}}, c_{\text{QM}}^2)$ . Its predictions for hybrid star branches then follow from the generic CSS phase diagram.
- ▶ Every observation, e.g. observing a  $2M_{\odot}$  neutron star,  $\Rightarrow$  constraint on CSS parameters.  
E.g., for soft NM we need  $c_{\text{QM}}^2 \gtrsim 1/3$   
(But note that  $c_{\text{QM}}^2 = 1/3 - \mathcal{O}(\alpha_s)$  in pert QCD).
- ▶ More measurements of  $M$  and  $R$  would strengthen the constraints.
- ▶ Models of quark matter tend to have  $c_{\text{QM}}^2 \sim 1/3$  and high transition pressure  $\Rightarrow$  very short hybrid branch.

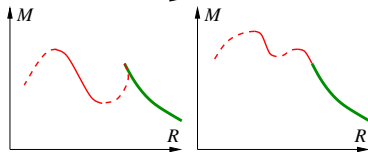
# Could we identify hybrid stars via $M(R)$ ?

We could identify a phase transition to a high-density phase

(A) Nuclear branch ends with  $dM/dR \neq 0$   
occurs if  $\Delta\varepsilon/\varepsilon_{\text{trans}}$  is large enough



(B,D) Disconnected branch  
can occur with  $M_{\text{max}} \gtrsim 2M_{\odot}$  if nuclear *and*  
quark matter are both stiff ( $c_{\text{QM}}^2 \sim 1$ )

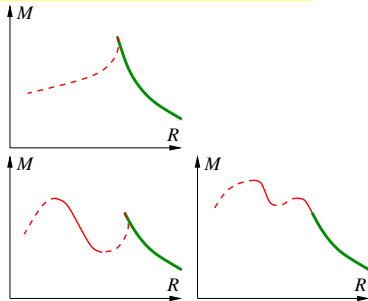


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We need:

- ▶ better measurements of  $M$  and  $R$
- ▶ knowledge of nuclear matter EoS

We could benefit from:

- ▶ theoretical constraints on parameters of QM EoS  
( $p_{\text{trans}}/\varepsilon_{\text{trans}}$ ,  $\Delta\varepsilon/\varepsilon_{\text{trans}}$ ,  $c_{\text{QM}}^2$ )