

The experimental quest for  
in-medium effects  
Episode I

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# The standard model and QCD

FERMIONS			matter constituents spin = 1/2, 3/2, 5/2, ...		
Leptons spin = 1/2			Quarks spin = 1/2		
Flavor	Mass GeV/c <sup>2</sup>	Electric charge	Flavor	Approx. Mass GeV/c <sup>2</sup>	Electric charge
$\nu_e$ electron neutrino	$<1 \times 10^{-8}$	0	u up	0.003	2/3
e electron	0.000511	-1	d down	0.006	-1/3
$\nu_\mu$ muon neutrino	$<0.0002$	0	c charm	1.3	2/3
$\mu$ muon	0.106	-1	s strange	0.1	-1/3
$\nu_\tau$ tau neutrino	$<0.02$	0	t top	175	2/3
$\tau$ tau	1.7771	-1	b bottom	4.3	-1/3

BOSONS			force carriers spin = 0, 1, 2, ...		
Unified Electroweak spin = 1			Strong (color) spin = 1		
Name	Mass GeV/c <sup>2</sup>	Electric charge	Name	Mass GeV/c <sup>2</sup>	Electric charge
$\gamma$ photon	0	0	g gluon	0	0
W <sup>-</sup>	80.4	-1			
W <sup>+</sup>	80.4	+1			
Z <sup>0</sup>	91.187	0			

- Strong interaction:
  - binds quarks into hadrons
  - binds nucleons into nuclei
- Described by QCD:
  - interaction between particles carrying color charge (quarks, gluons)
- Mediated by strong force carriers (gluons)
- Very successful theory
  - jet production
  - particle production at high  $p_T$
  - heavy flavor production
  - ...
- ... *but with outstanding puzzles*

# Two puzzles in QCD: confinement

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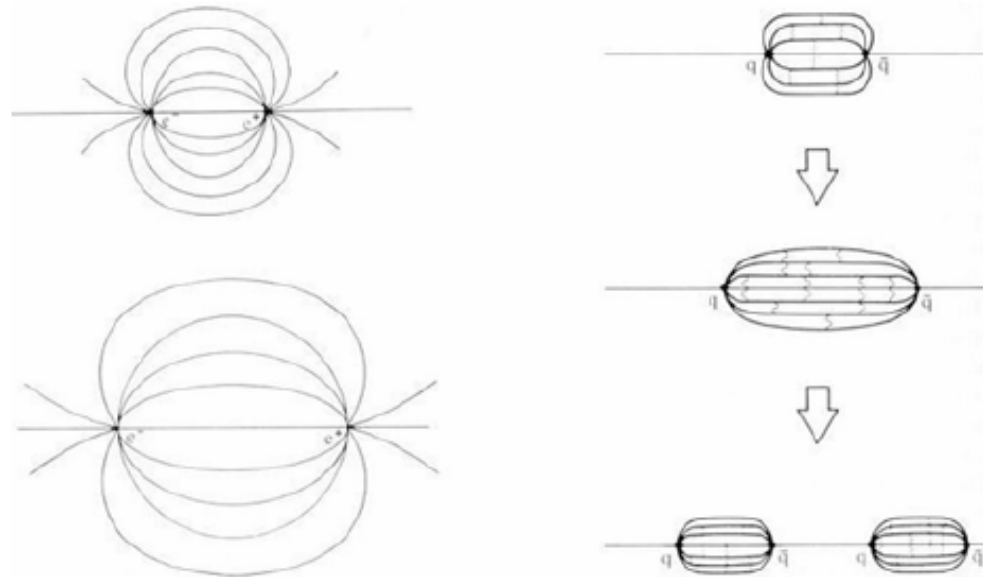
- Nobody ever succeeded in detecting an isolated quark
- Quarks seem to be permanently confined within protons, neutrons, pions and other hadrons
- It looks like one half of the fundamental fermions are not directly observable...

... *how does this come about?*

# QCD confinement

- If the distance between two quarks gets larger, more and more gluons contribute to the interaction between the quarks.
- Hence the potential energy grows with increasing distance.
- At some point, enough energy is stored in the field to produce a pair of quarks out of the vacuum (observed as jet).

$$V(r) \propto -\frac{\alpha_s(r)}{r} + Kr$$



## Two puzzles in QCD: hadron masses

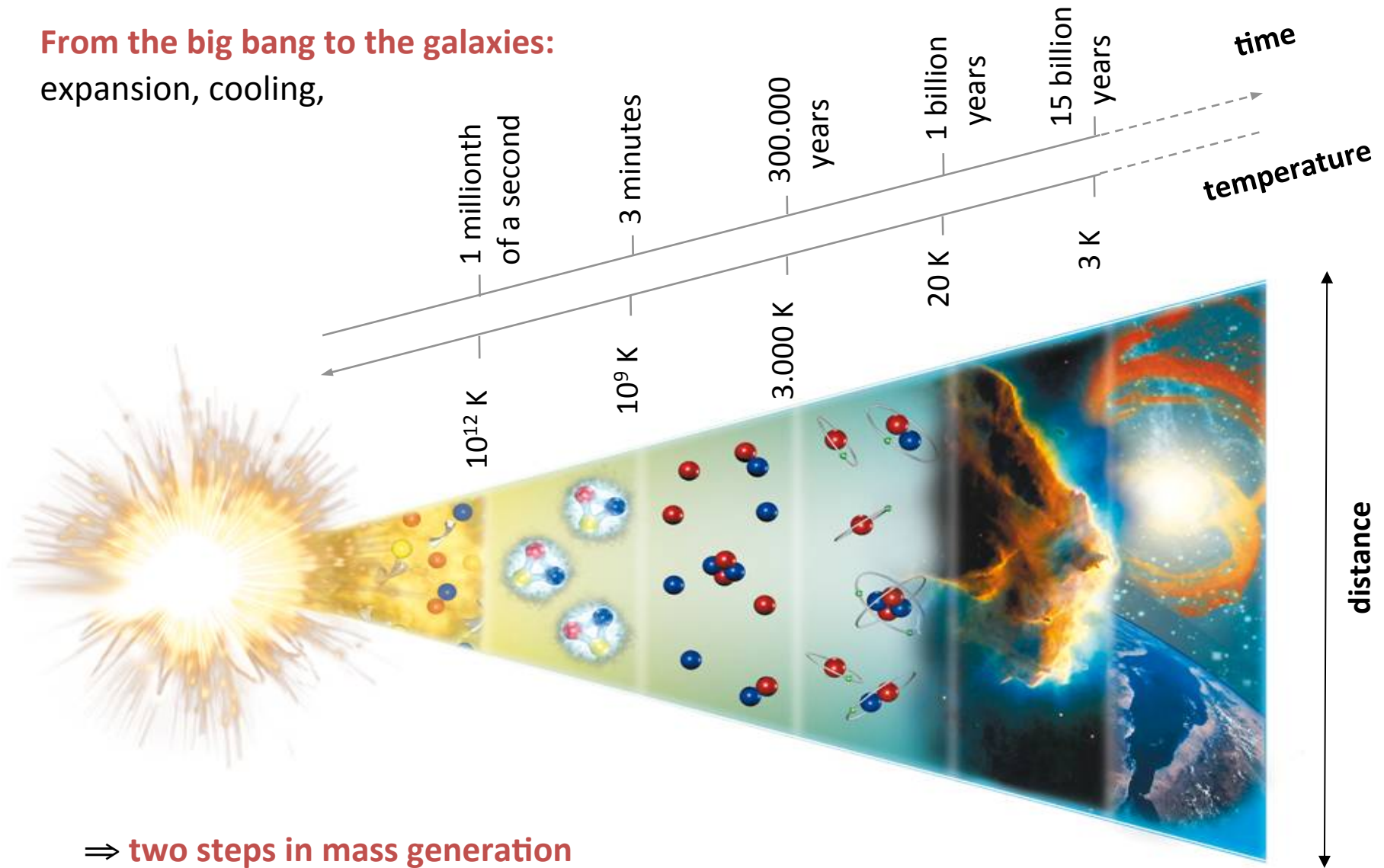
- A proton is thought to be made of two u and one d quarks
- The sum of their masses is around 12 MeV
- *... but the proton mass is 938 MeV!*



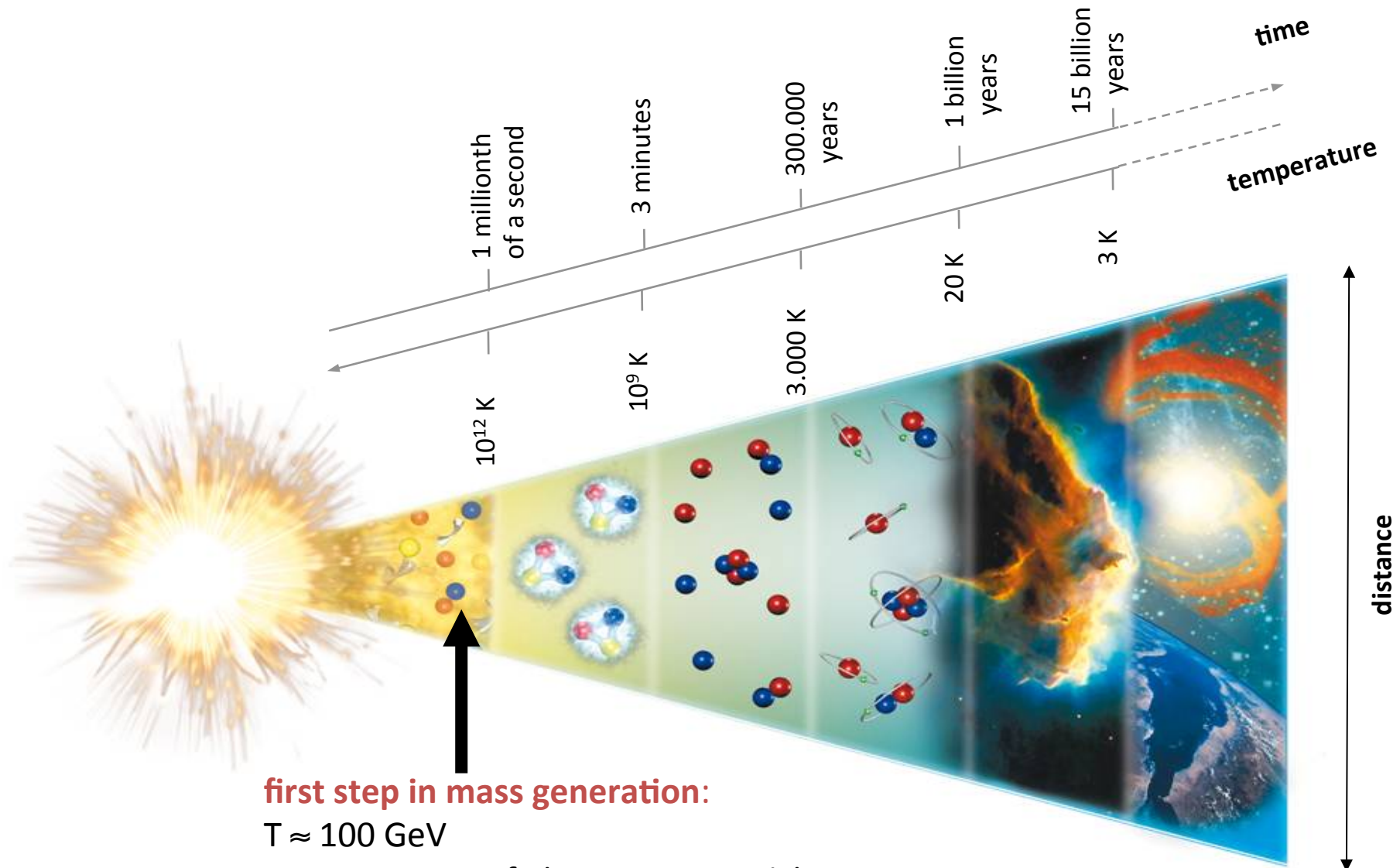
***How does nature generate massive hadrons from nearly mass-less quarks?***

# Evolution of the Universe

**From the big bang to the galaxies:**  
expansion, cooling,



# Evolution of the Universe

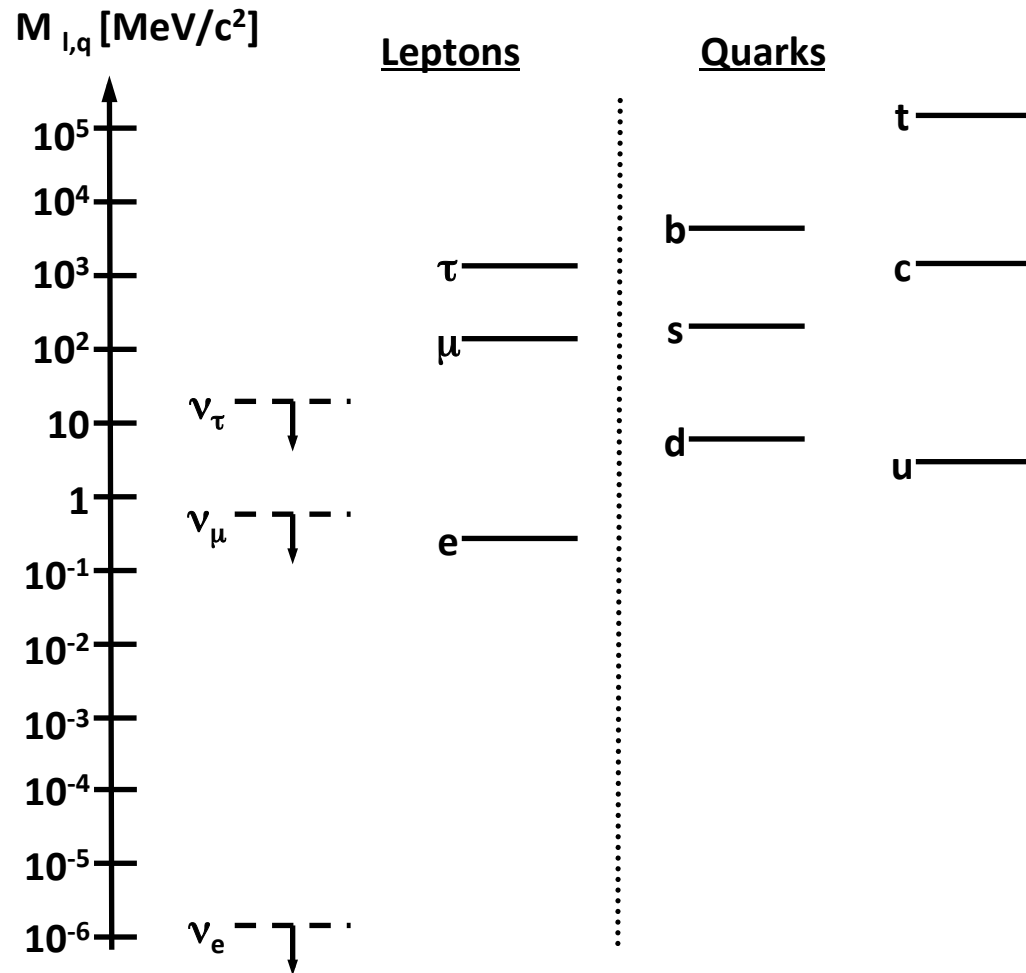


**first step in mass generation:**

$T \approx 100 \text{ GeV}$

Higgs  $\Leftrightarrow$  mass of elementary particles

# Masses of quarks and leptons

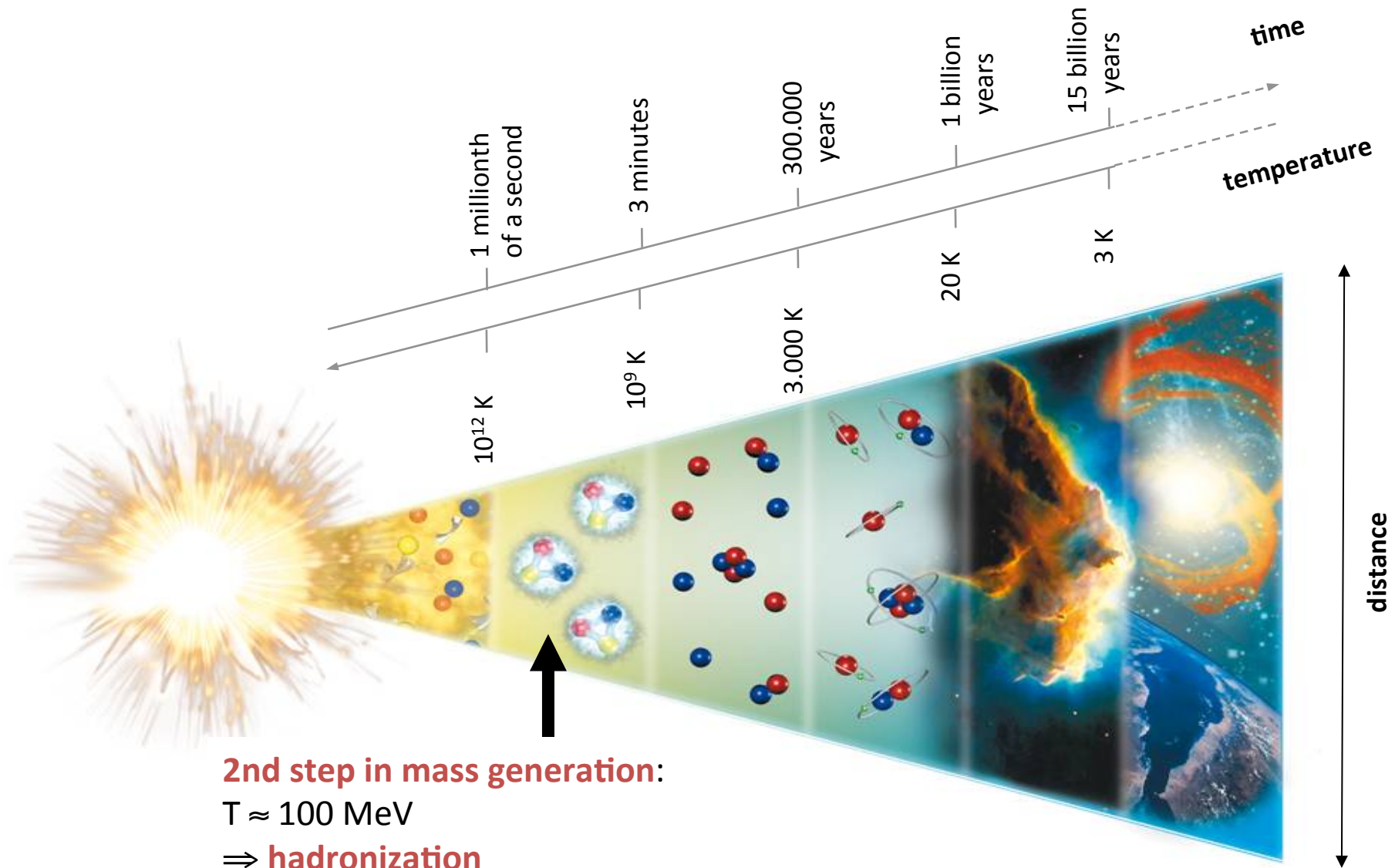


masses of elementary particles  
(quarks, leptons) generated  
by interaction with Higgs-field

⇒ search for **Higgs-particle** (LHC)



# Evolution of the Universe



**2nd step in mass generation:**

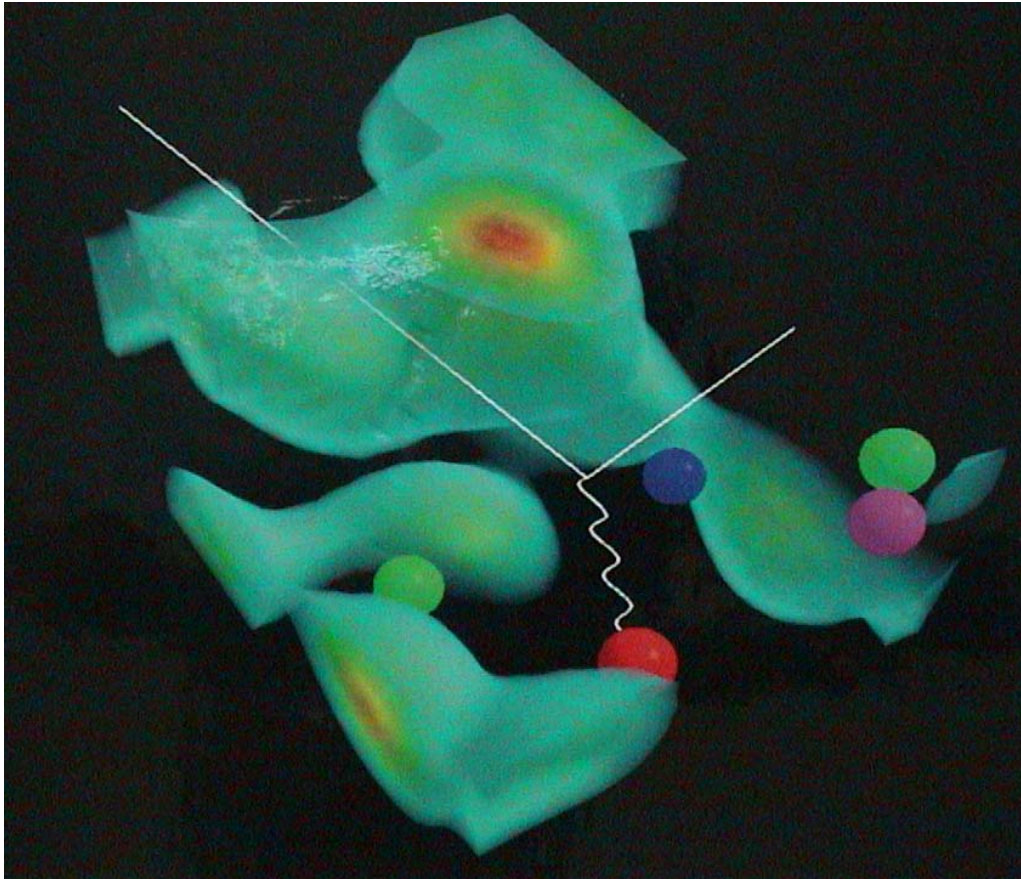
$T \approx 100 \text{ MeV}$

$\Rightarrow$  **hadronization**

## Condensation in vacuum

- **Higgs** generates **~2%** and  
**QCD** generates **98%** of the mass of ordinary matter !!!
- **How can we experimentally prove this scenario?**
- Experiments at the large hadron collider (LHC) at CERN will search the Higgs particle, the missing piece in the Standard Model.
- However, the chiral condensate cannot be studied this way, it is not an observable. Theoretical models are used to link observables to the quark-condensate.

## The nucleon is a complex object



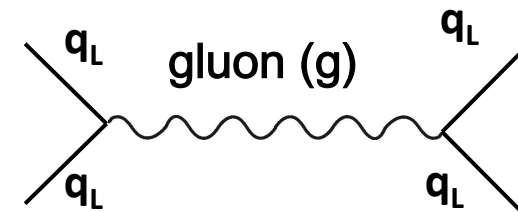
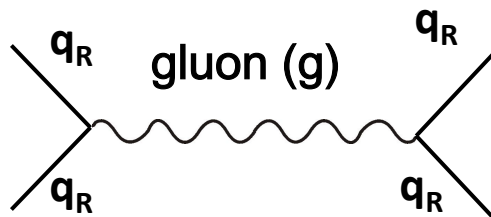
*Hadrons are very complex excitations of valence quarks in the presence of quark and gluon condensates.*

nucleon: mass not determined by sum of constituent masses  $m = E/c^2$ ; „mass without mass“ (Wilczek) mass given by energy stored in motion of quarks and by energy in color gluon fields

# Role of chiral symmetry

Chiral symmetry = fundamental symmetry of QCD for massless quarks ( $m_q=0$ )

In the interaction among quarks by gluon exchange right-handed quarks  $q_R$  (spin and momentum parallel) stay right-handed and left-handed quarks  $q_L$  stay left-handed  $\rightarrow$  chirality is conserved



**For  $m_q=0$  the QCD Lagrangian is invariant under the  $SU(3)_R \otimes SU(3)_L$  transformations**

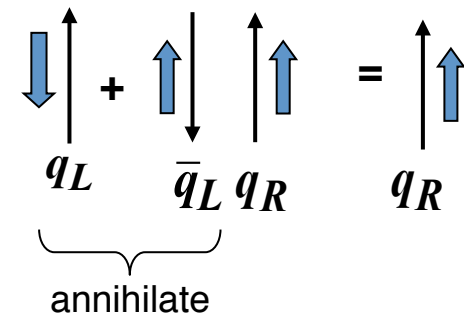
# Chiral symmetry breaking

The ground state of QCD (vacuum) is populated by quark – anti-quark pairs ( $\langle q\bar{q} \rangle$  condensate) and does not share the symmetry of the Lagrangian



A left-handed quark  $q_L$  can be converted into a right-handed quark  $q_R$  (spin and momentum parallel) by interaction with a scalar  $q - \text{anti-}q$  pair

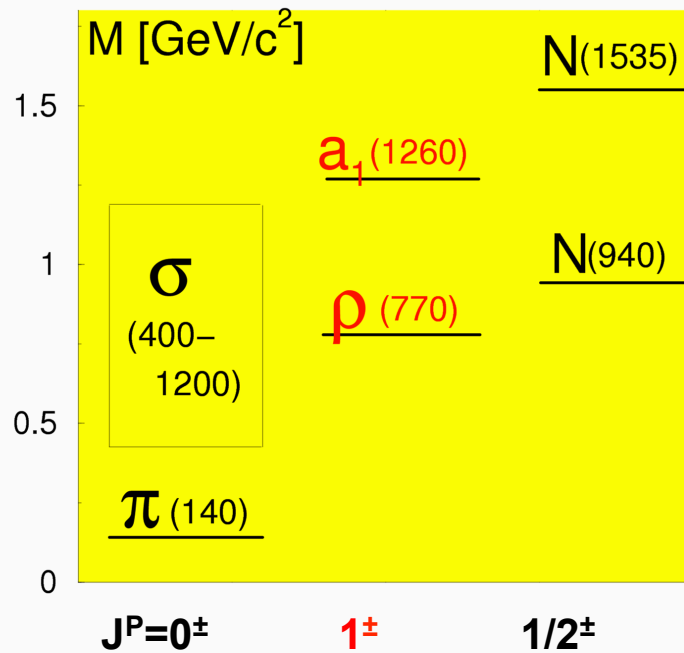
Due to the condensate  
*chiral symmetry is broken!*



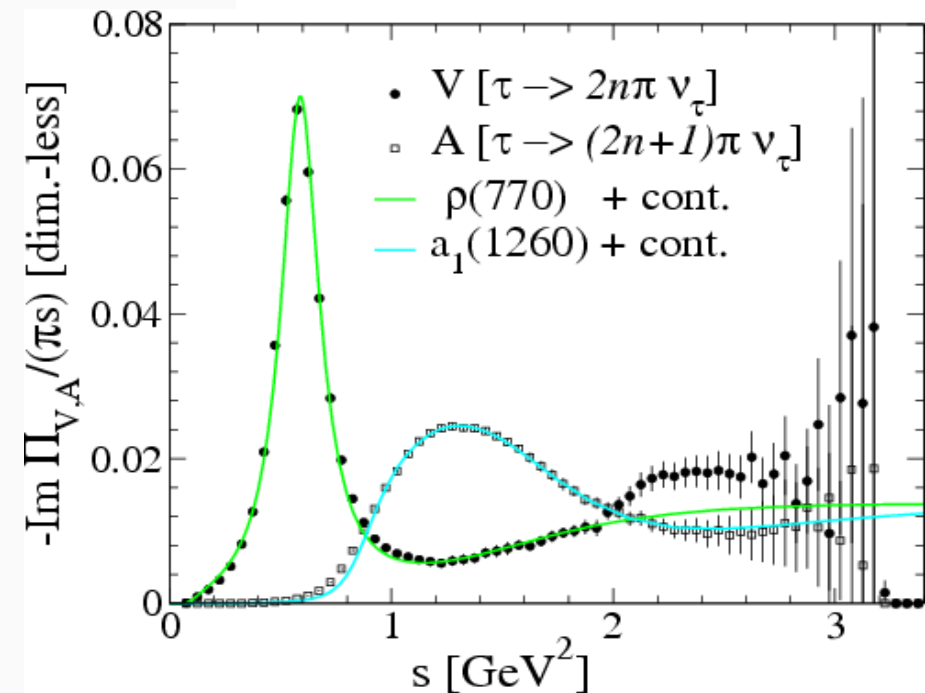
# Consequences of Spontaneous Breaking of Chiral Symmetry

- If **chiral symmetry** were to hold also in the **hadronic sector** we would expect chiral partners with same spin but opposite parity to be **degenerate in mass**:
  - e.g., nucleon N:  $J^\pi = 1/2^+$ ; chiral partner:  $J^\pi = 1/2^-$  mass degenerate??

“chiral partners” split:  $\Delta M \approx 0.5\text{GeV}$   
 mass split comparable to hadron masses



## Axial-/Vector Correlators



What happens if nuclear matter is compressed or heated?

- Compressed ( $\mu_B$ ):

- ▣ Less volume for a given number of baryons

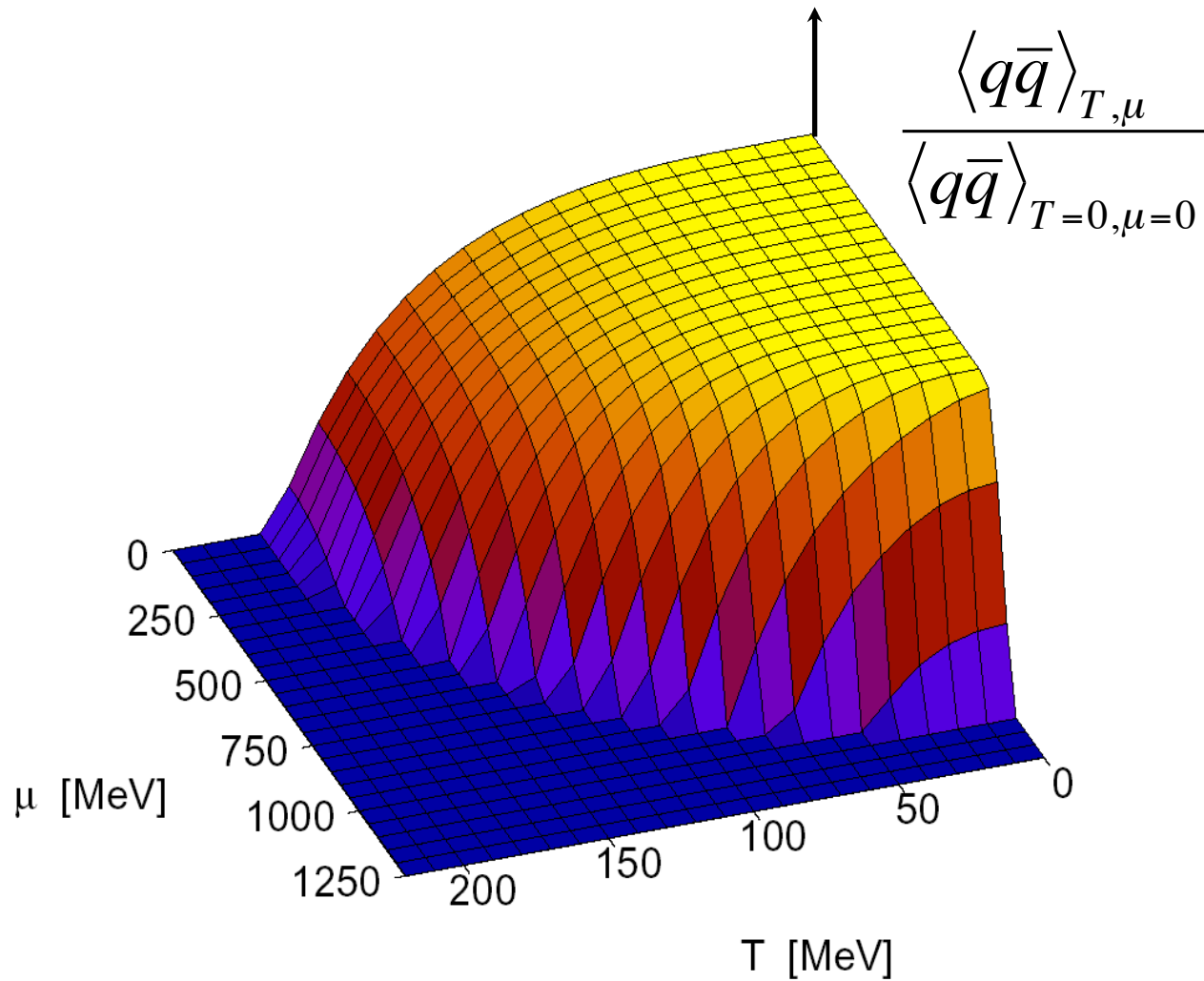
- ◆ Less condensate

- Heated (T):

- ▣ Additional pions

- ◆ Less condensate

# Properties of condensate in-medium



However  
 $\langle q\bar{q} \rangle$   
*is not an observable!!*

*B.J. Schäfer and J.Wambach*



## QCD sum rules

- However,  $\langle q\bar{q} \rangle$  is not an observable!!
- **QCD sum rules** provide a link between hadronic observables and condensates:  
(T. Hadsuda and S. Lee, PRC 46 (1992) R34; S. Leupold and U. Mosel, PRC58 (1998) 2939)

$$\frac{Q^2}{24\pi^2} \int ds \frac{R(s)}{(s+Q^2)^2} = \frac{1}{16\pi^2} \left( 1 + \frac{\alpha_s}{\pi} \right) + \frac{1}{Q^4} \left[ m_q \langle \bar{q}q \rangle + \frac{1}{24} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle \right] + \text{higher order terms}$$

hadronic spectral function: 
$$R(s) \sim F^2 \frac{1}{\pi} \frac{\sqrt{s} \Gamma(s)}{(s - M_\rho^2)^2 + s(\Gamma(s))^2}$$

- Chiral condensate related only to integral over hadronic spectral functions;  
→ spectral function are constrained, but not determined

⇒ **Hadronic models are still needed for specific predictions of hadron properties !!**

## Medium modifications of hadrons

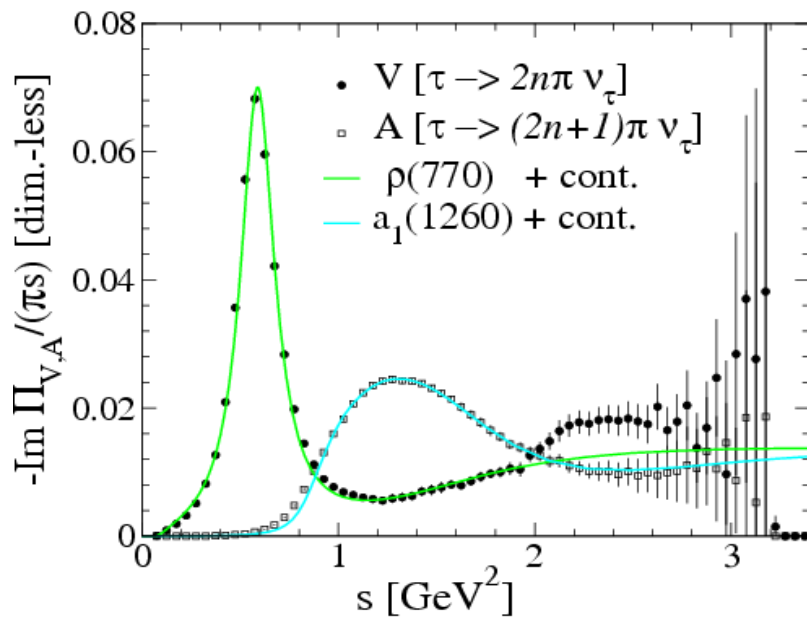
- Many models:
  - hadron mass and quark condensate are linked →
  - expect modification of hadron spectral properties (mass  $m$ , width  $\Gamma$ )
  - How is this realized?
    - Do the masses drop to zero (or simply change)?
    - Do the widths' increase (melting resonances)?
  - Good questions, without (obvious) good answers
  - ... at least chiral partners should become degenerate.

# Chiral symmetry restoration

- Light-quark sector of QCD: chiral symmetry
  - Spontaneously broken in vacuum
  - High temperature/density: restoration of chiral symmetry

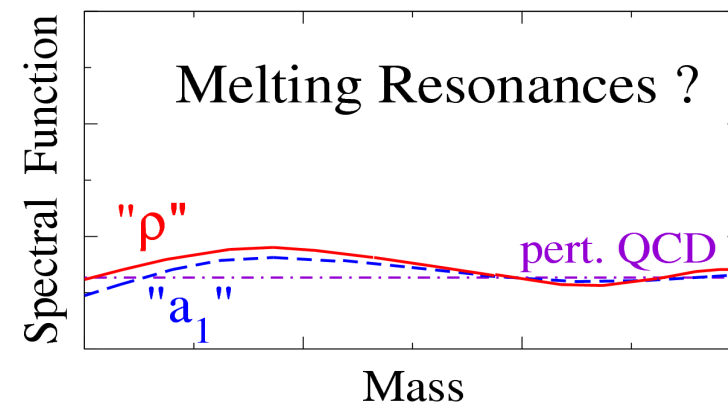
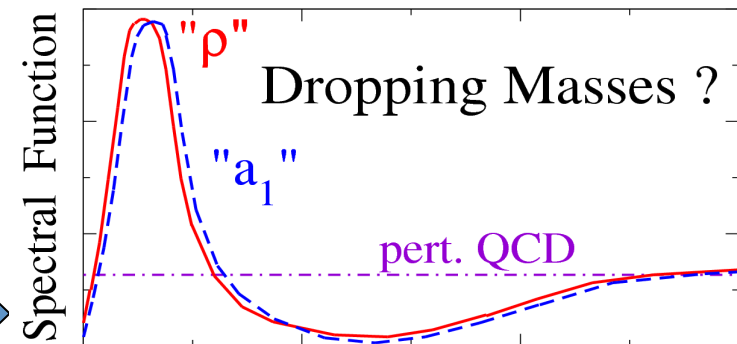
**Vacuum**

Vector and Axial vector spectral functions, hadronic  $\tau$  decays



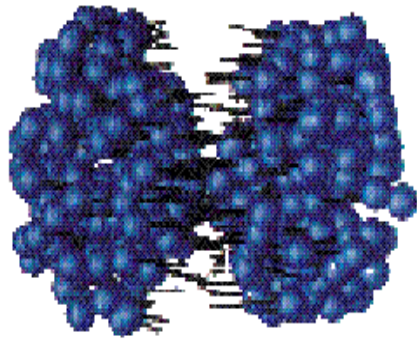
**Medium**

Schematic scenarios for chiral symmetry restoration

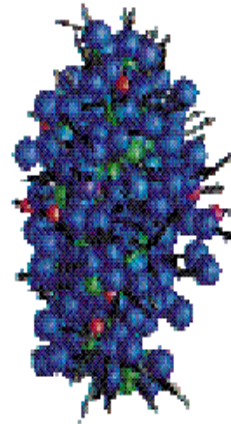


# Dileptons as probes in heavy-ion collisions

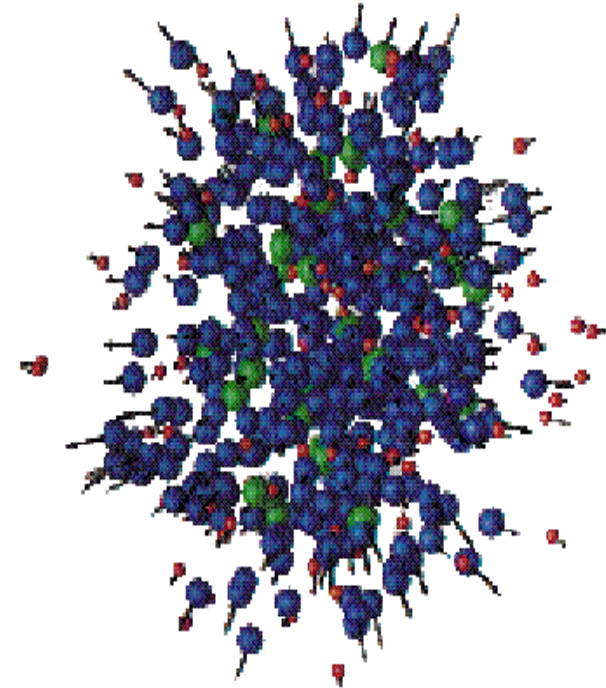
A+A at  $E_{kin} = 2 \text{ GeV/u}$



two colliding  
nuclei



formation of highly  
compressed and  
heated collision zone  
 $\tau \cong 10^{-23} \text{ s}$

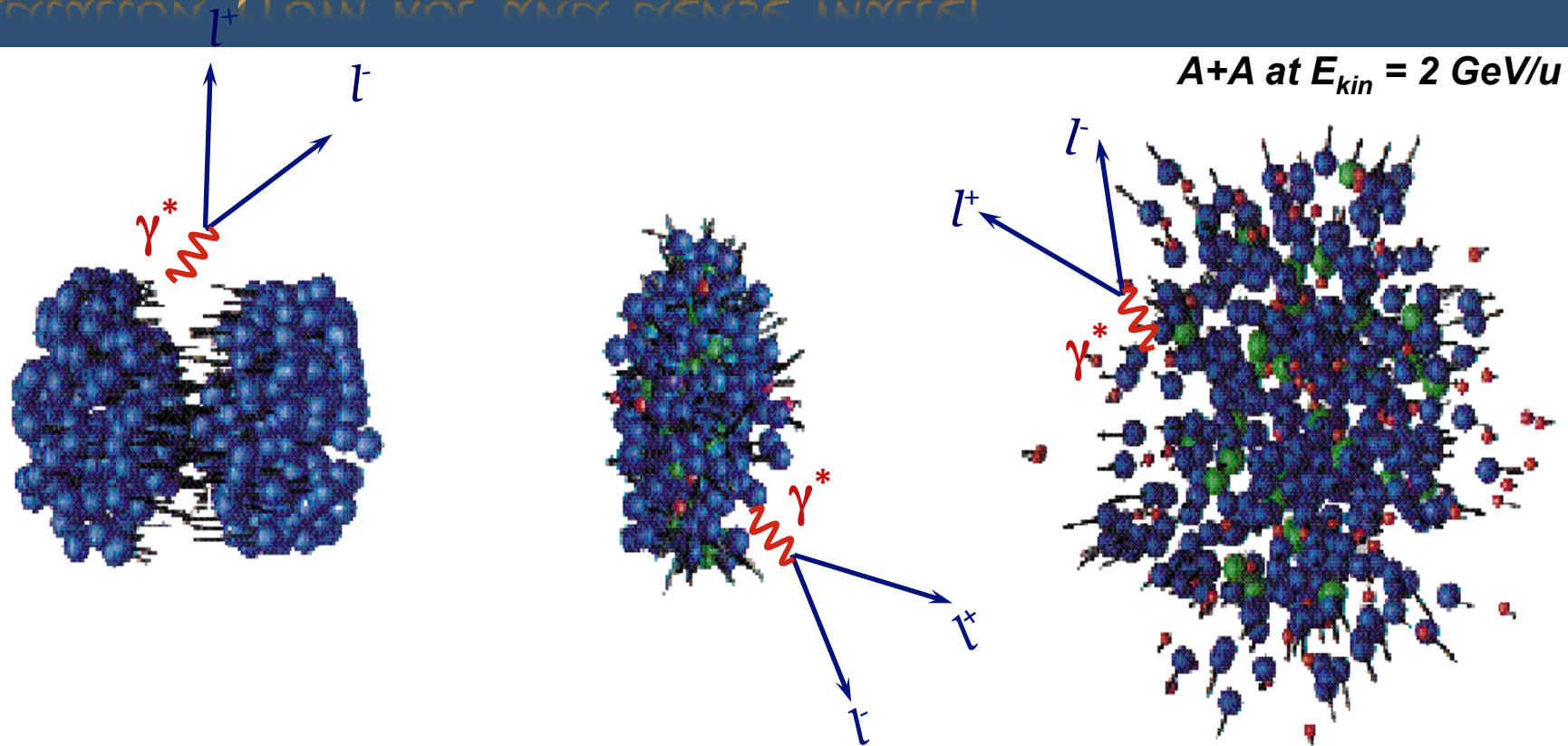


explosion of  
collision zone

**Challenge:**

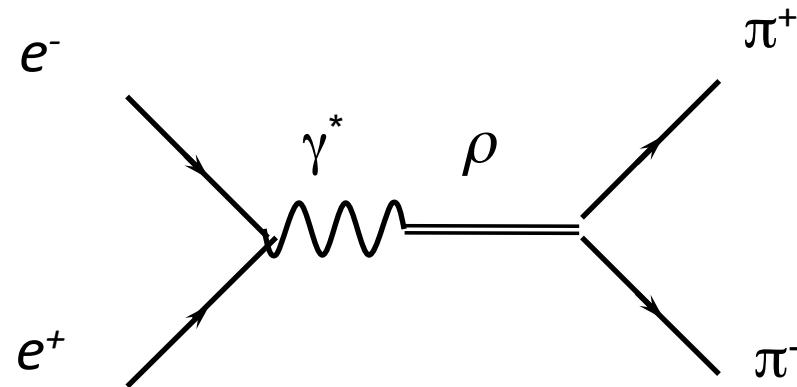
*Extract information on the high density phase*

# Radiation from hot and dense matter



- The dilepton signal contains **contributions from throughout the collision**
- No strong final state interactions  
→ **leave reaction volume undisturbed**
- Probes the **electromagnetic structure of dense/hot hadronic matter**

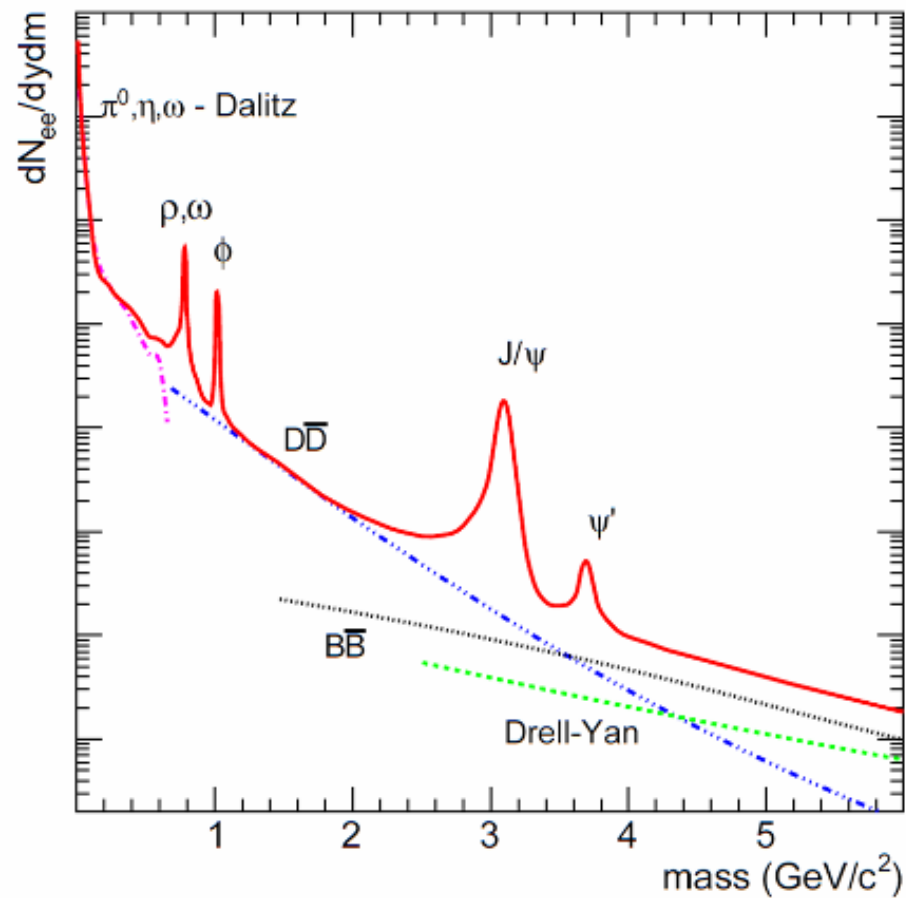
# Vector mesons Dominance Model



- $J^P = 1^-$  for both  $\gamma^*$  and Vector Meson
- Strong coupling of  $\gamma^*$  to Vector Meson  
→ Vector Meson Dominance model
- *Observable*: vector mesons ( $\rho$ ,  $\omega$ ,  $\phi$ ).

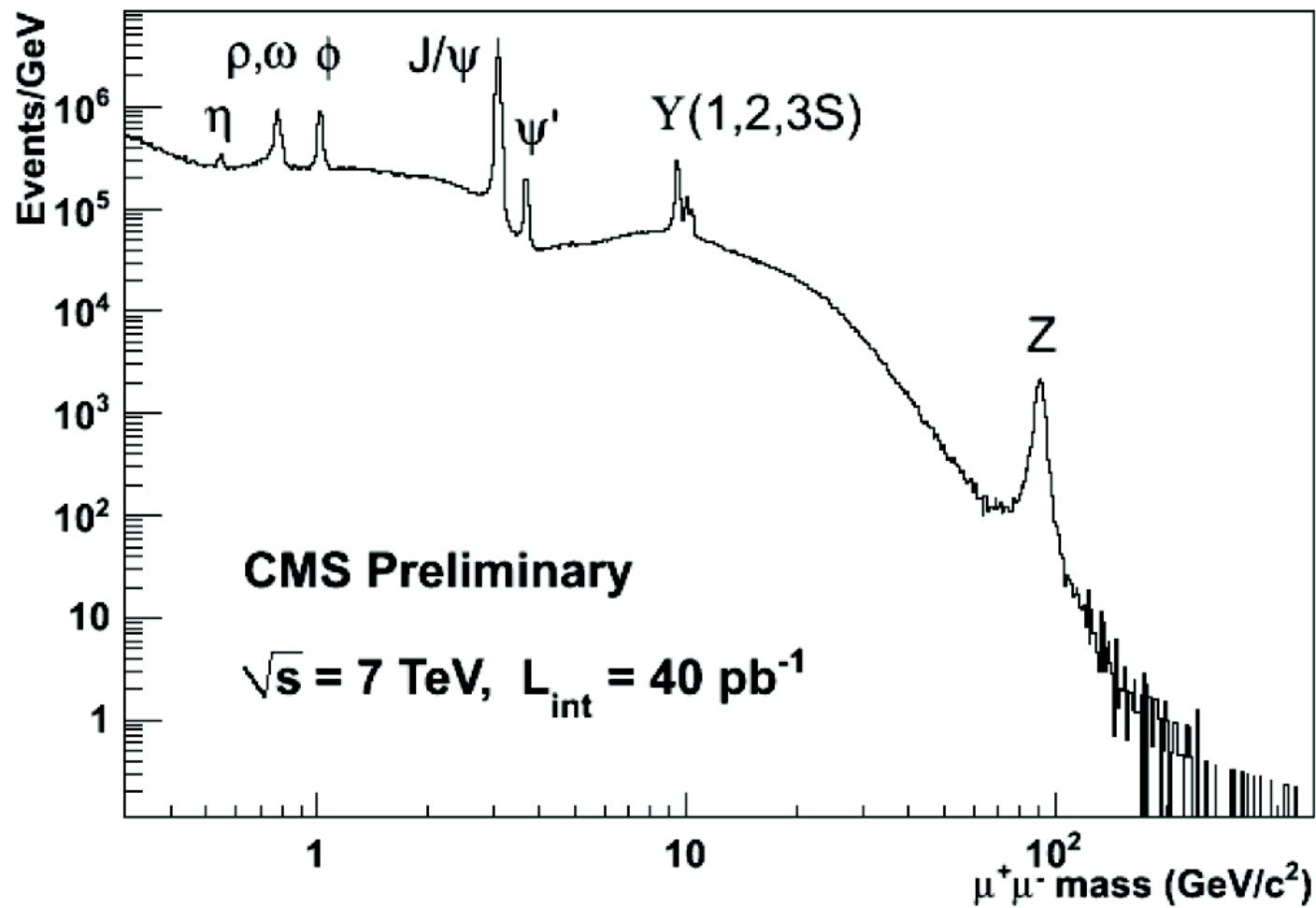
# Observable: vector mesons

*Schematical spectral distribution of lepton pairs emitted in ultra-relativistic heavy ion collisions*



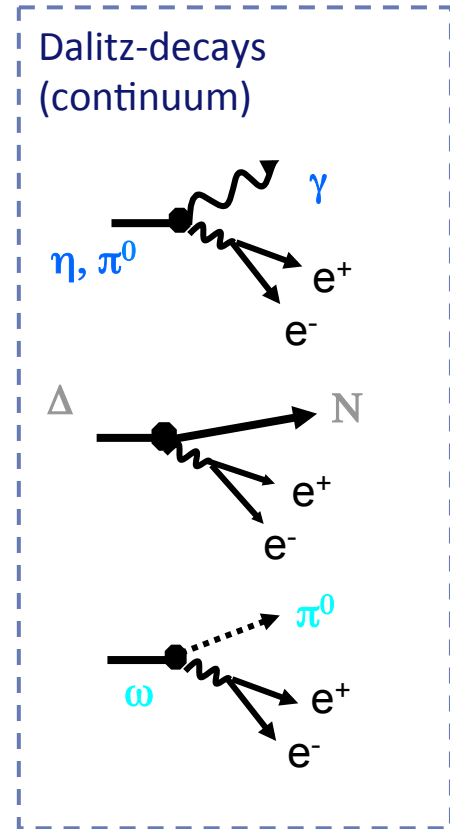
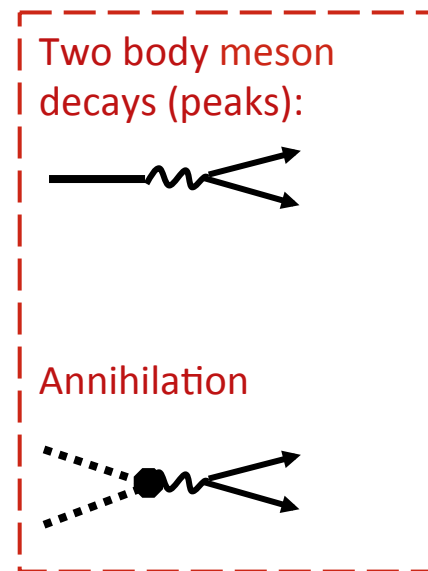
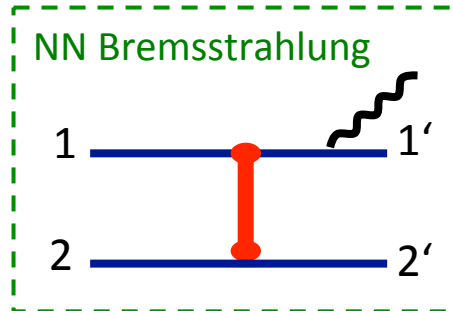
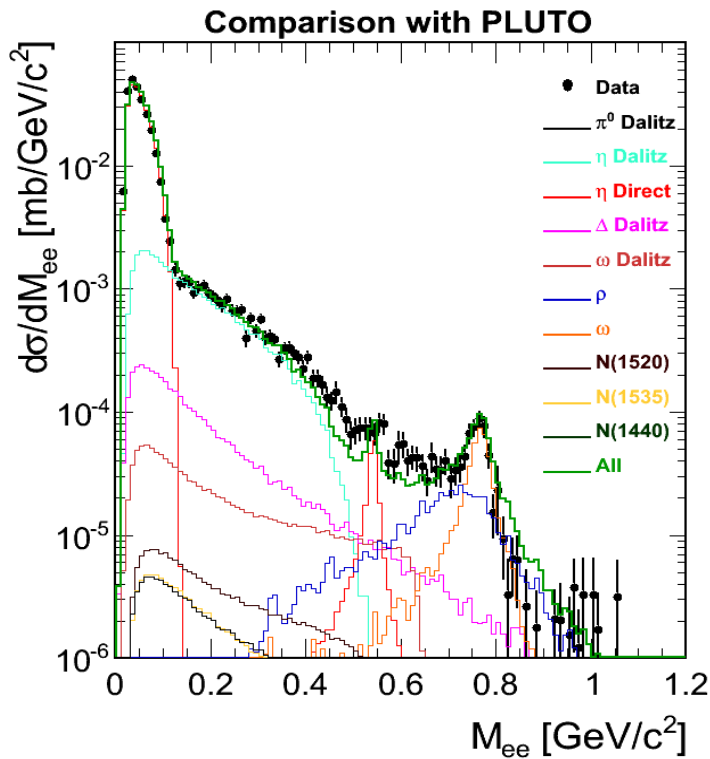
$$M_{l+l-} = 2 \cdot \sin \frac{\theta_{l+l-}}{2} \cdot \sqrt{p_{l+} \cdot p_{l-}}$$

# Dimuon spectrum from pp at LHC

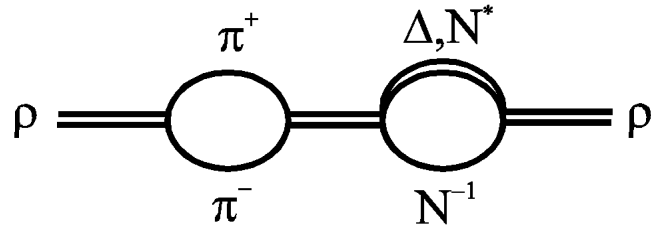




# The electron pair cocktail at low beam energies

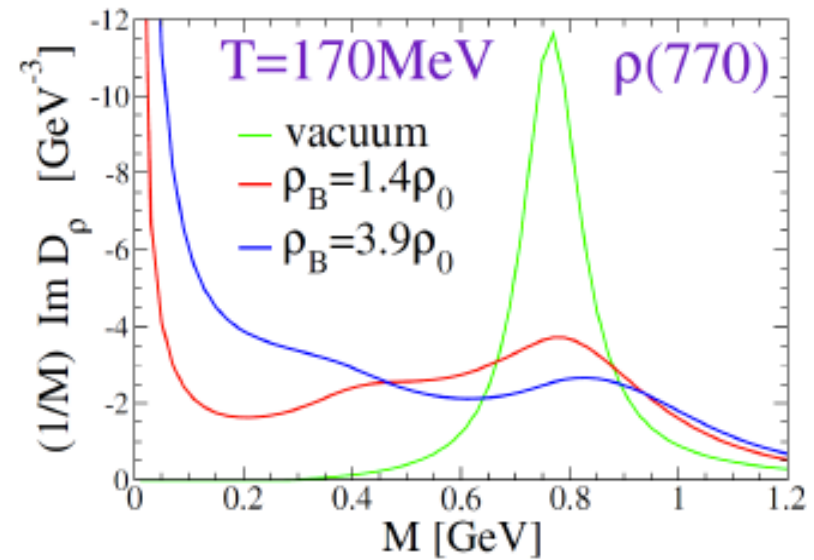
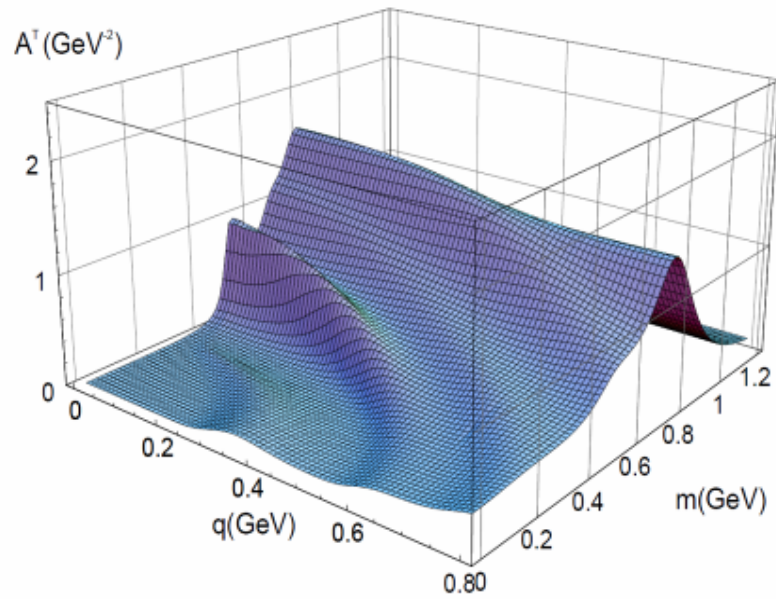


# Spectral function of the $\rho$ -meson in-medium



$$D_\rho(M, q; \mu_B, T) = \frac{1}{\left[ M^2 - m_\rho^2 - \Sigma_{\rho\pi\pi} - \Sigma_{\rho B} - \Sigma_{\rho M} \right]}$$

**Additional contributions to the  $\rho$ -meson self-energy in the medium**



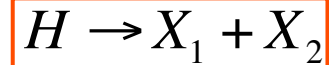
More predictions for in-medium properties of the  $\rho$  meson:

	mass of $\rho$	width of $\rho$
Pisarski 1982	↘	↗
Leutwyler et al 1990 ( $\pi, N$ )	→	↗
Brown/Rho 1991	↘	→
Hatsuda/Lee 1992	↘	→
Dominguez et. al 1993	→	↗
Pisarski 1995	↗	↗
Rapp 1996	→	↗

*One example where experiments have the potential to guide the theory*

# Experimental approach

Hadron decay in the medium:



- reconstruction of invariant mass from 4-momenta of decay products:

$$m_H(\rho, T, \vec{p}) = \sqrt{(p_1 + p_2)^2}$$

- compare  $m_H(\rho, T, \vec{p} \rightarrow 0)$  with  $m_H$  listed in PDG
- ensure that decays occur in the medium:
  - select shortlived mesons ( $c\tau = \frac{\hbar c}{\Gamma}$ ;  $\rho$ : 1.3 fm;  $\omega$ : 23 fm;  $\phi$ : 46 fm)
  - cut on low meson momenta
- avoid distortion of 4-momentum vectors by final state interaction

**⇒ dilepton spectroscopy:  $\rho, \omega, \phi \rightarrow e^+e^-$**

## Low-mass dileptons: what is been measured?

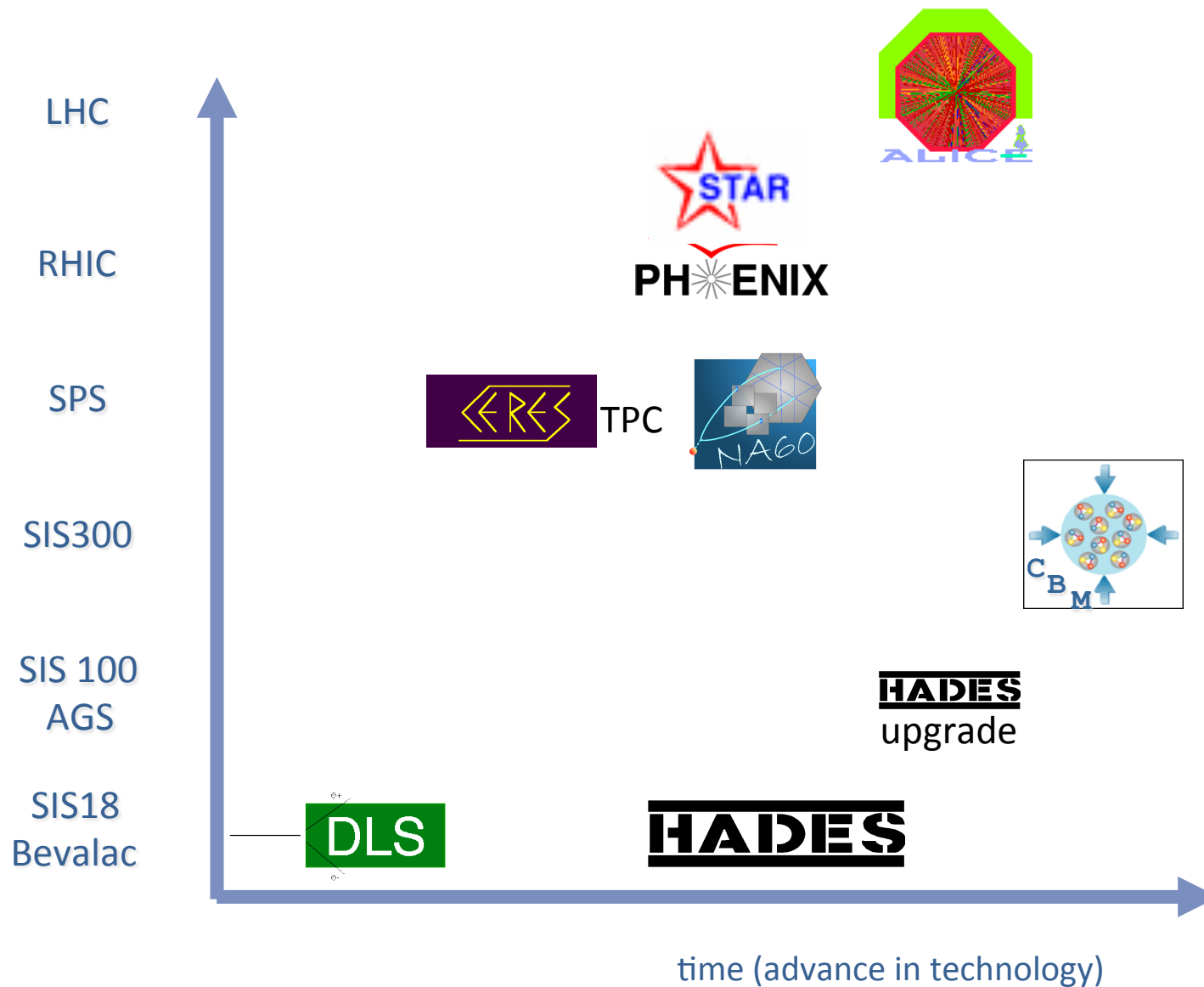
*...a needle in a haystack*

- Lepton pairs are rare probes (branching ratio  $< 10^{-4}$ )
- at SIS energies sub-threshold vector meson production
  
- Large combinatorial background in  $e^+e^-$  from:
  - Dalitz decays ( $\pi^0$ )
  - Conversion pairs
  
- Isolate the contribution to the spectrum from the dense stage

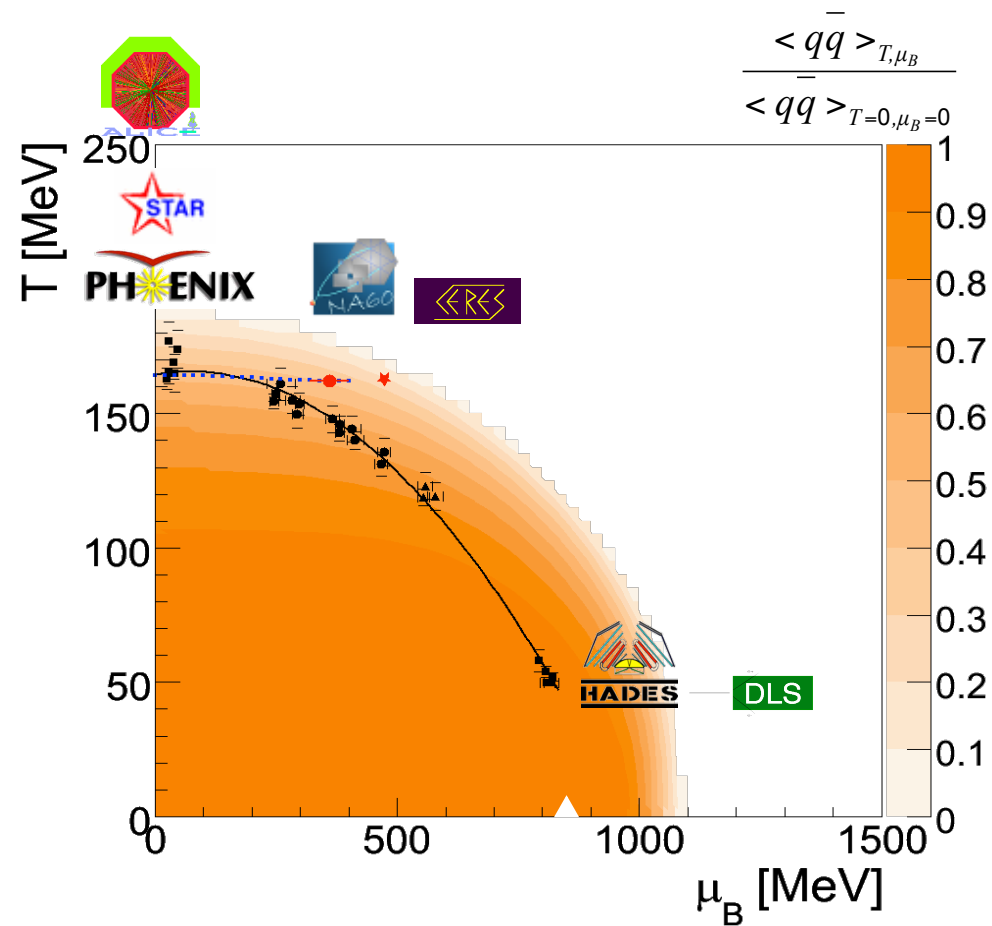
*Why not  $\rho \rightarrow \pi^+\pi^-$ ?*

*The branching ratios for hadronic decays of vector mesons are typically 4 orders of magnitude larger than for dilepton decays*

# Experiments addressing lepton pairs in HIC

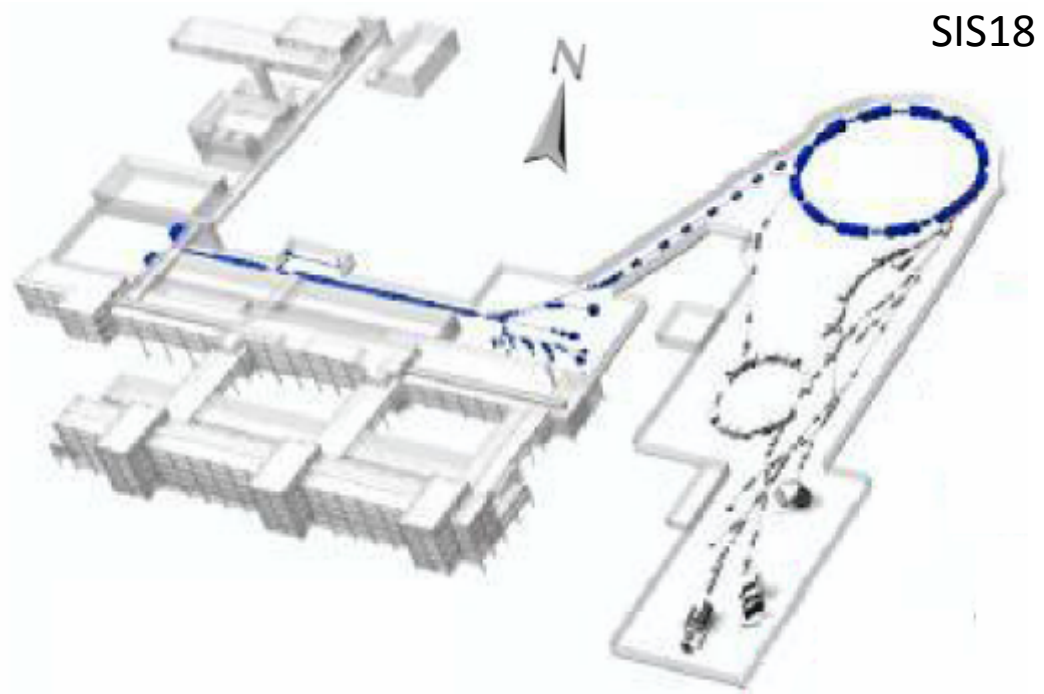


# Searching for the landmarks of the phase diagram of matter





# The HADES at GSI



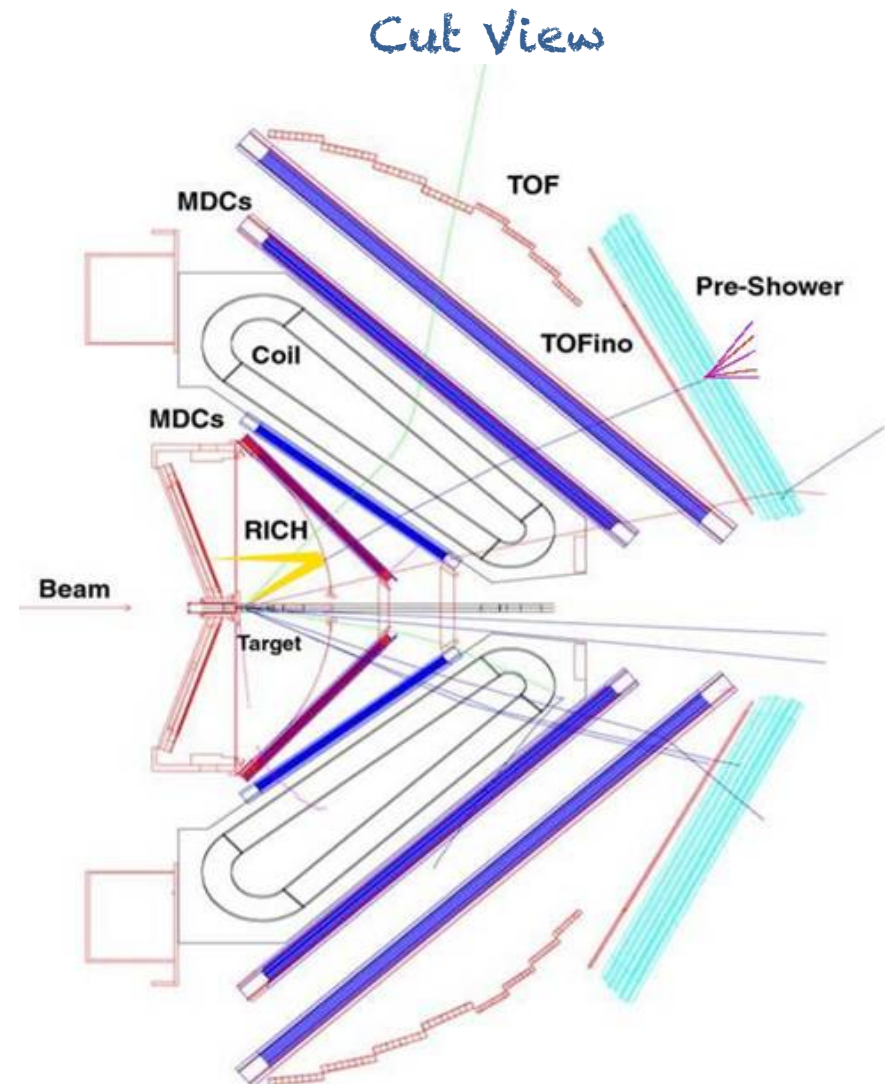
# High Acceptance DiElectron spectrometer



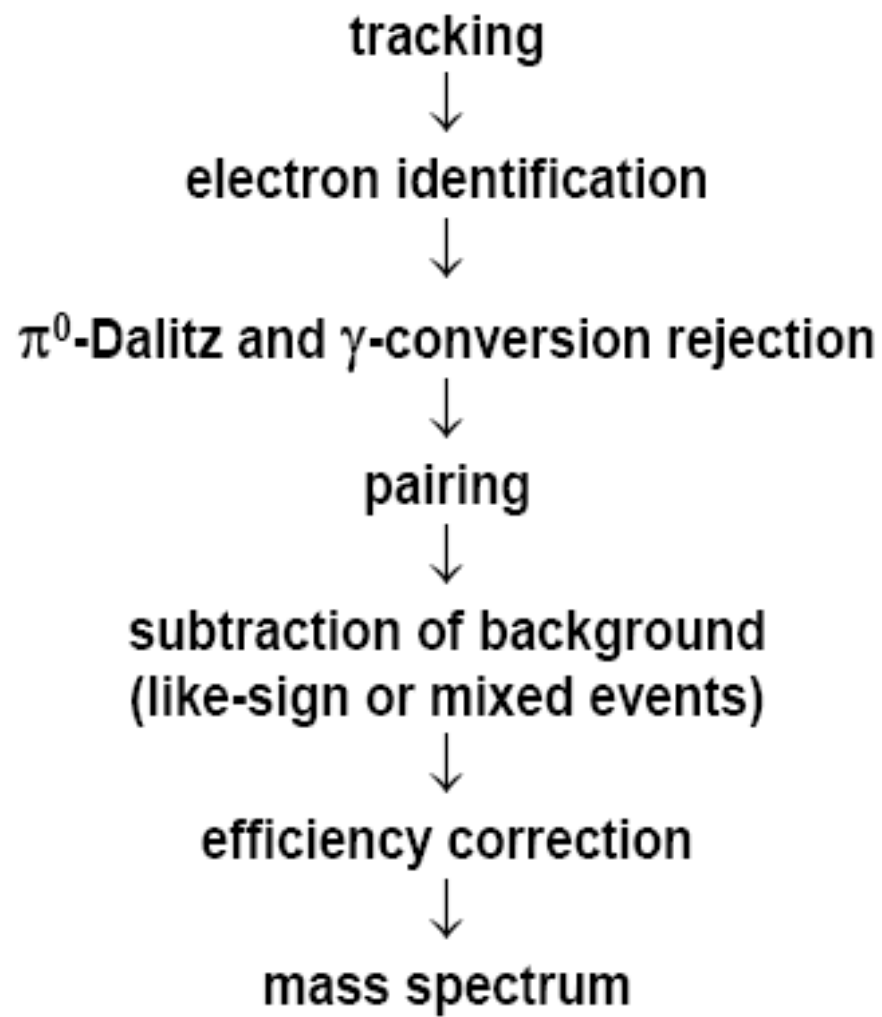
# HADES experiment

## Spectrometer with a...

- High geometrical acceptance
  - Full azimuth, polar angles  $18^\circ - 85^\circ$
  - Pair acceptance  $\approx 0.35$
- High invariant mass resolution (3% at  $\rho/\omega$  pole mass)
  - Low-mass tracking (superconducting toroidal magnet & multi-wire drift chamber (MDC), single cell resolution  $\approx 100 \mu\text{m}$ )
- Powerful PID capabilities:  $d/\pi/K/p/e$ 
  - RICH, TOF/TOFino, Pre-Shower, FW hodoscope: added 2007
- High background rejection & rate capability, dedicated LVL2 trigger:
  - LVL1: charge particle multiplicity
  - LVL2: single electron trigger



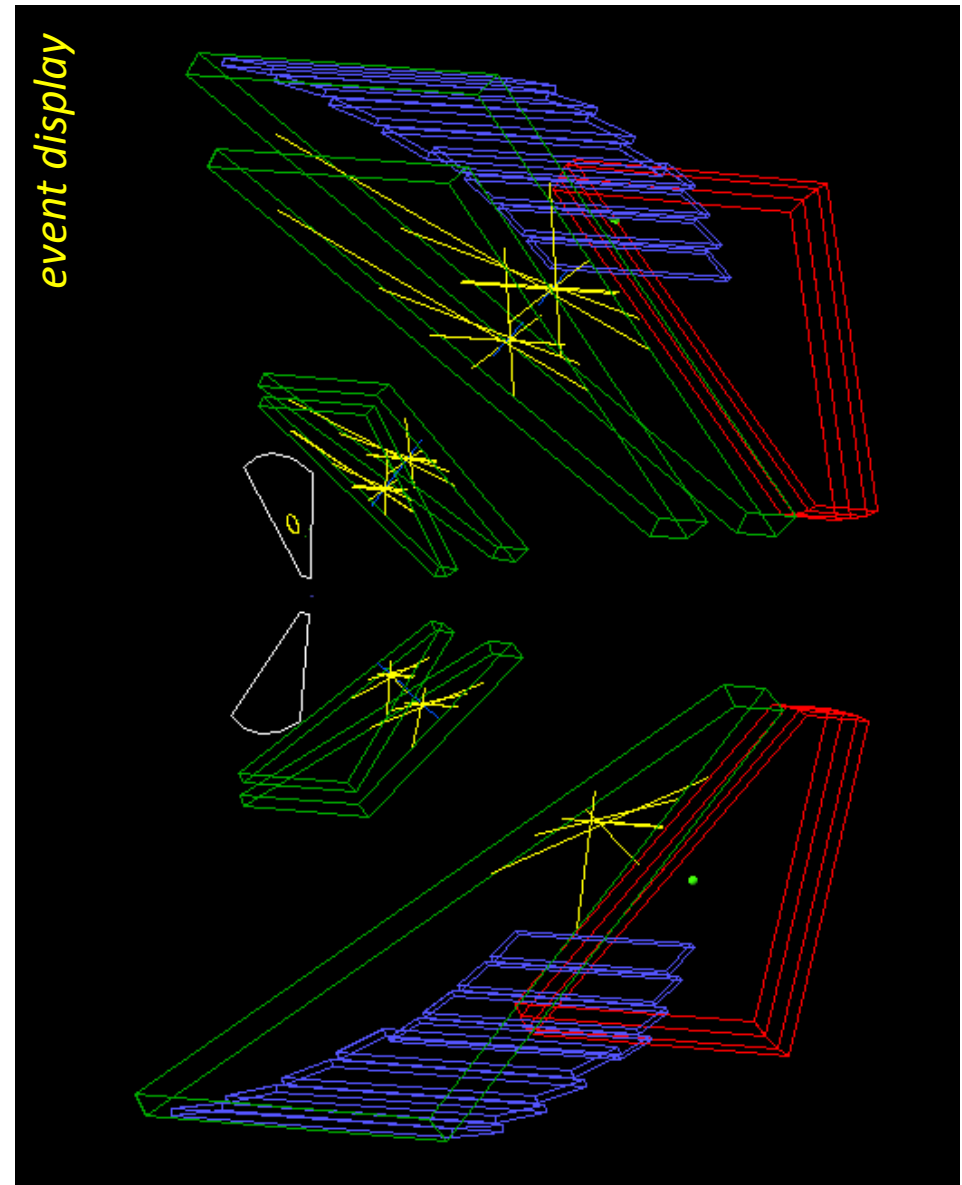
Just few steps ;)



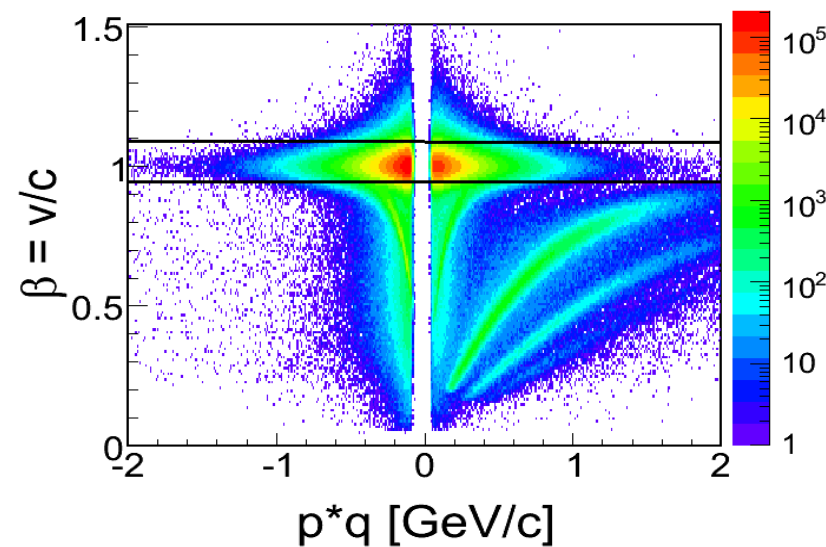
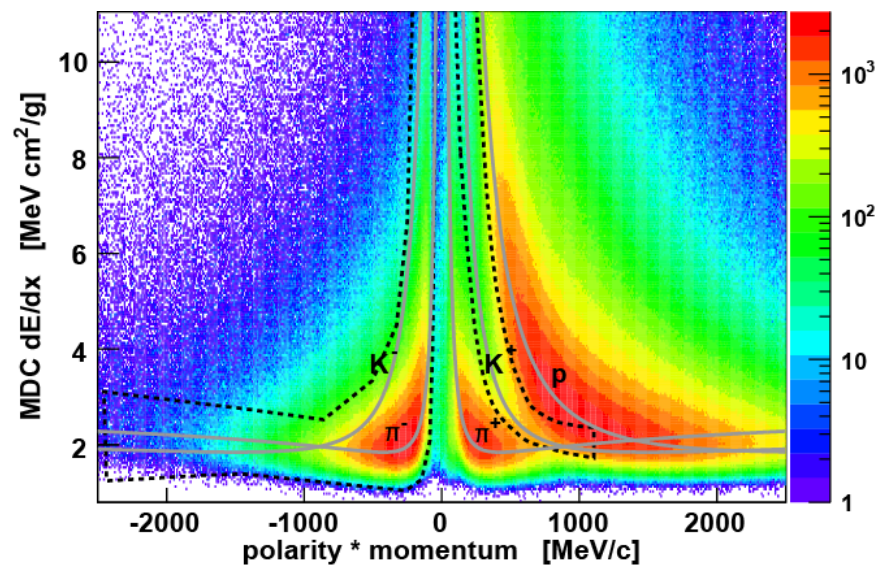
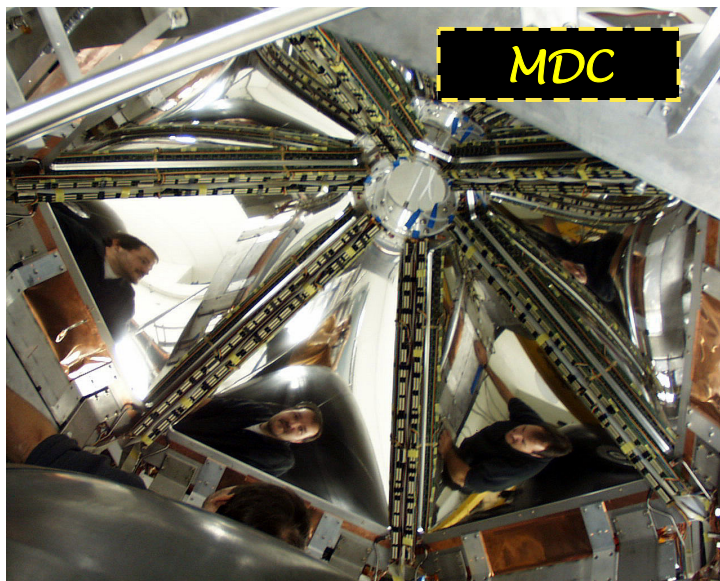
## How to reconstruct $\gamma^*$

$$M_{l+l-} = 2 \cdot \sin \frac{\theta_{l+l-}}{2} \cdot \sqrt{p_{l+} \cdot p_{l-}}$$

- Efficient track reconstruction
- Precise momentum determination
- Excellent electron/hadron identification

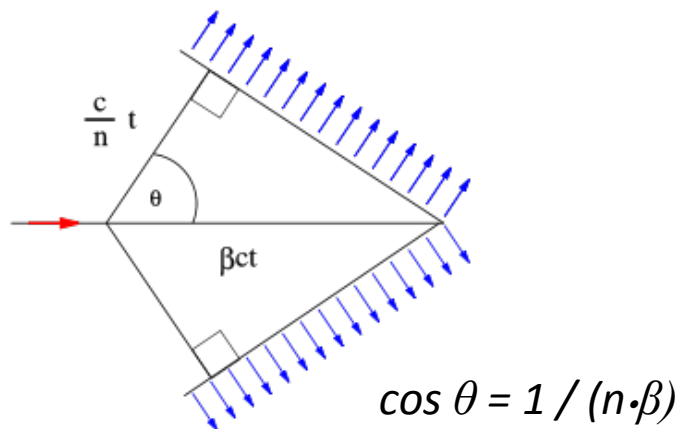


# Particle identification

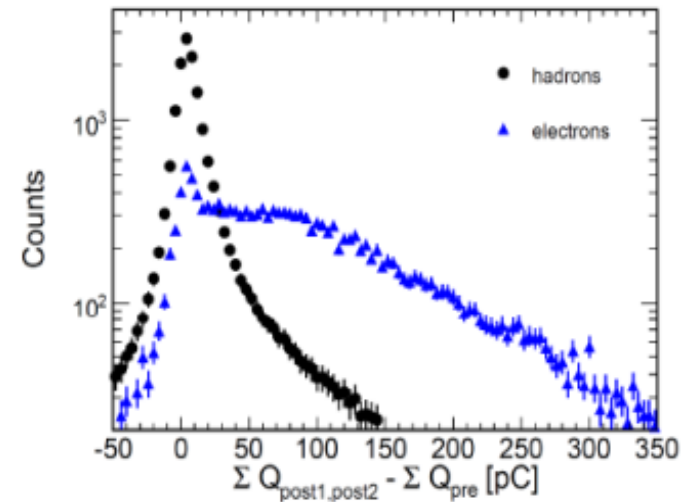
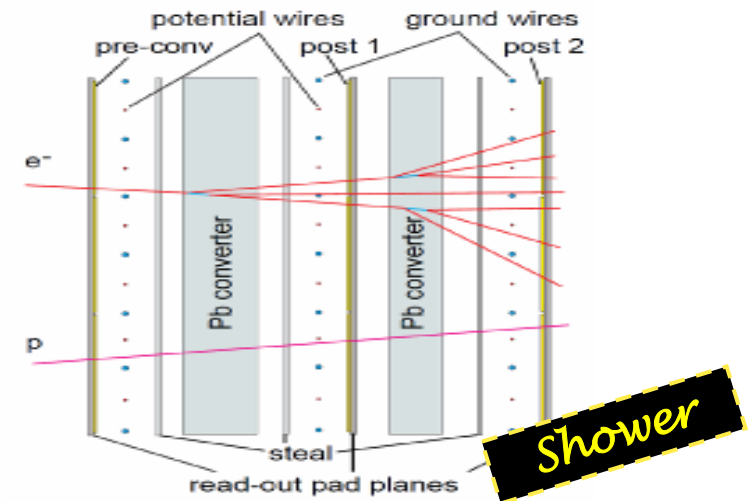


# Electron identification

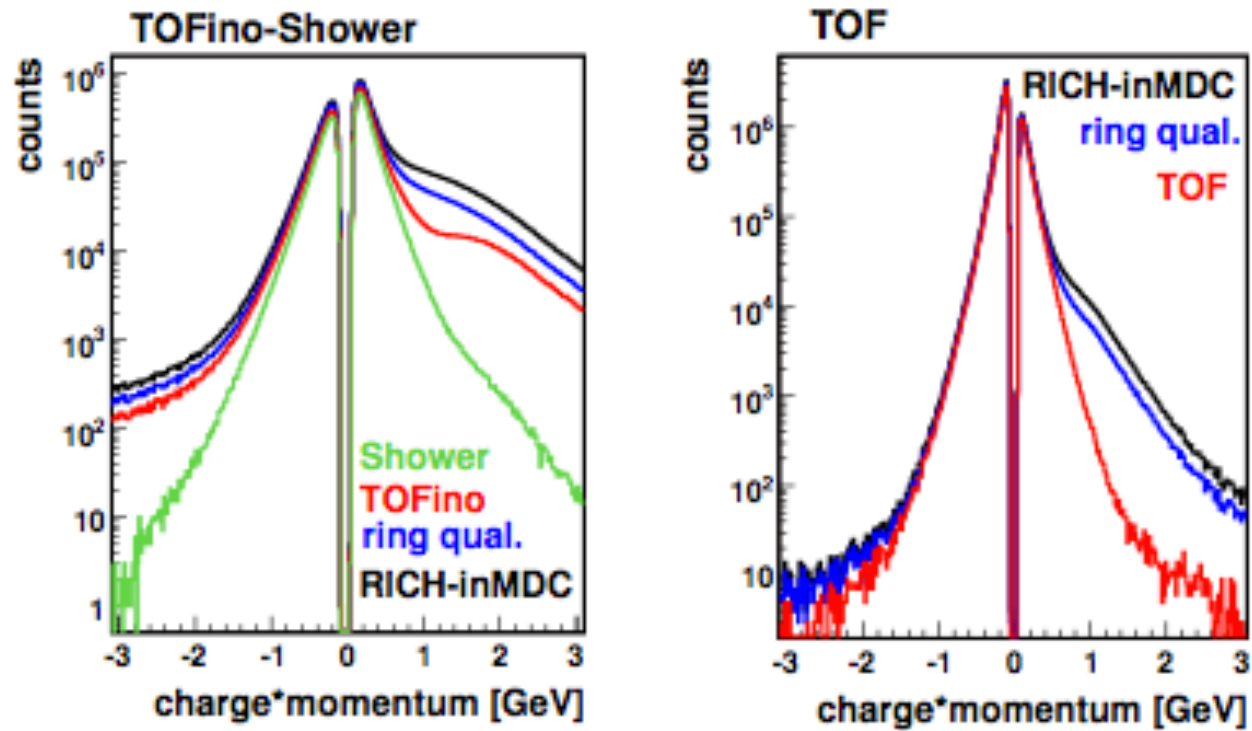
## Using Cherenkov effect



## Using Information on EM shower



# Single electron spectra

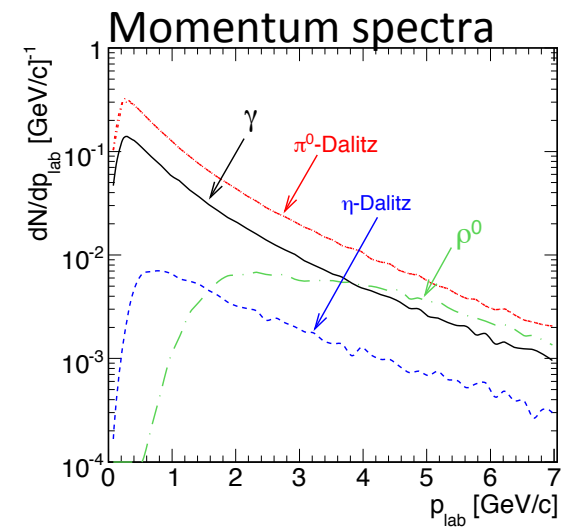
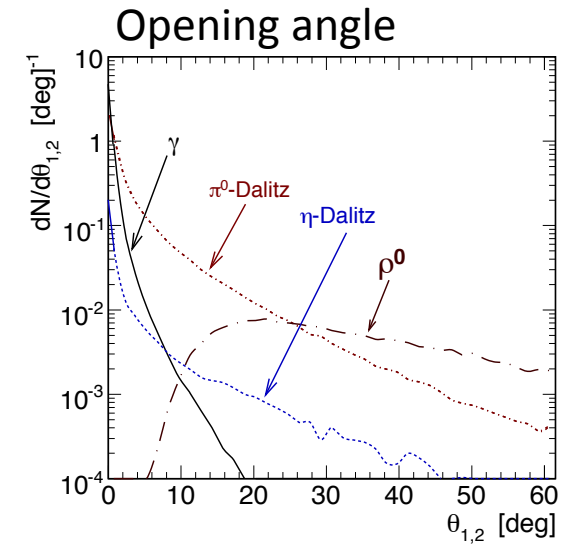
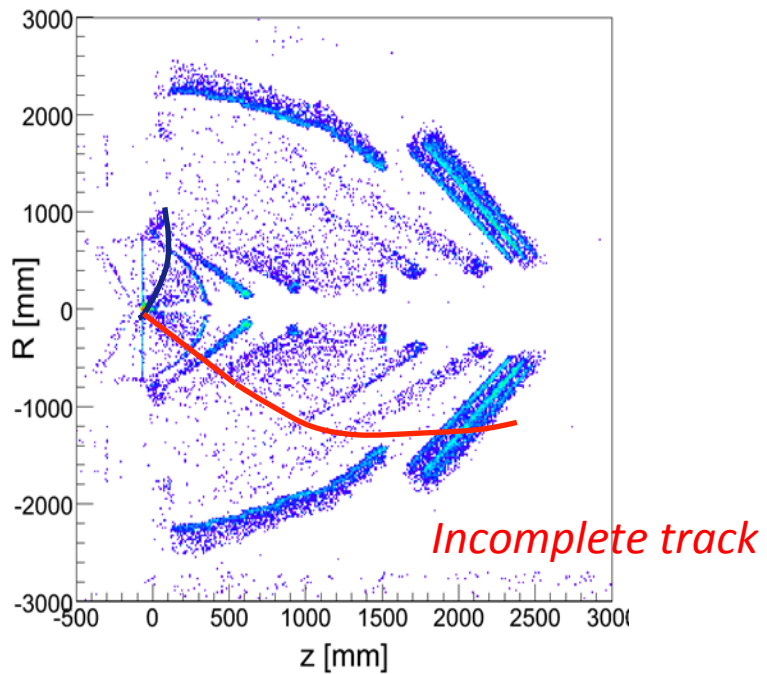
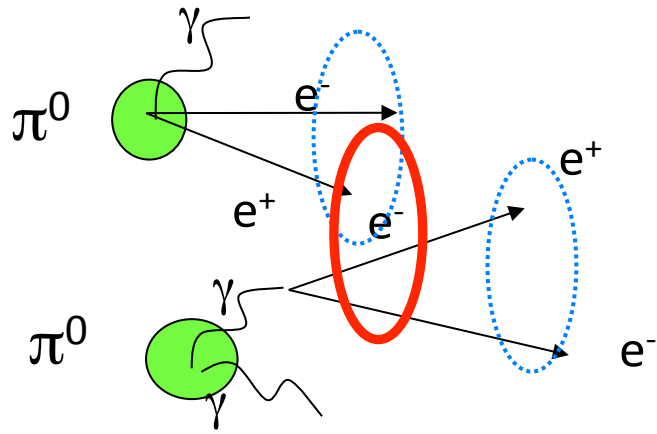


Clean electrons...

... but mainly from  $\pi^0$  Dalitz decays or from  $\gamma$  conversion



# Combinatorial background



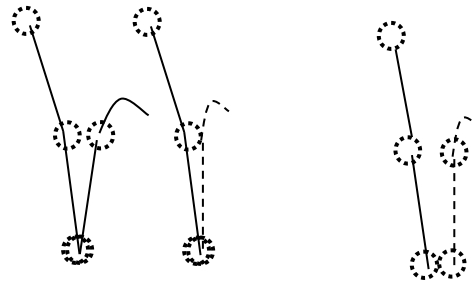
# Background rejection

TOF/Shower

B-field

MDC I-II

RICH



**C1**

**C2**

"Conversion rejection":

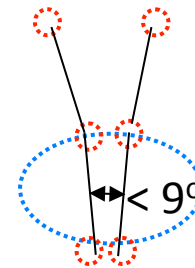
- C1, C2 = "close pair candidates" !
- Remove the track if a non-fitted track candidate is within  $10^\circ$

TOF/Shower

B-field

MDC I-II

RICH

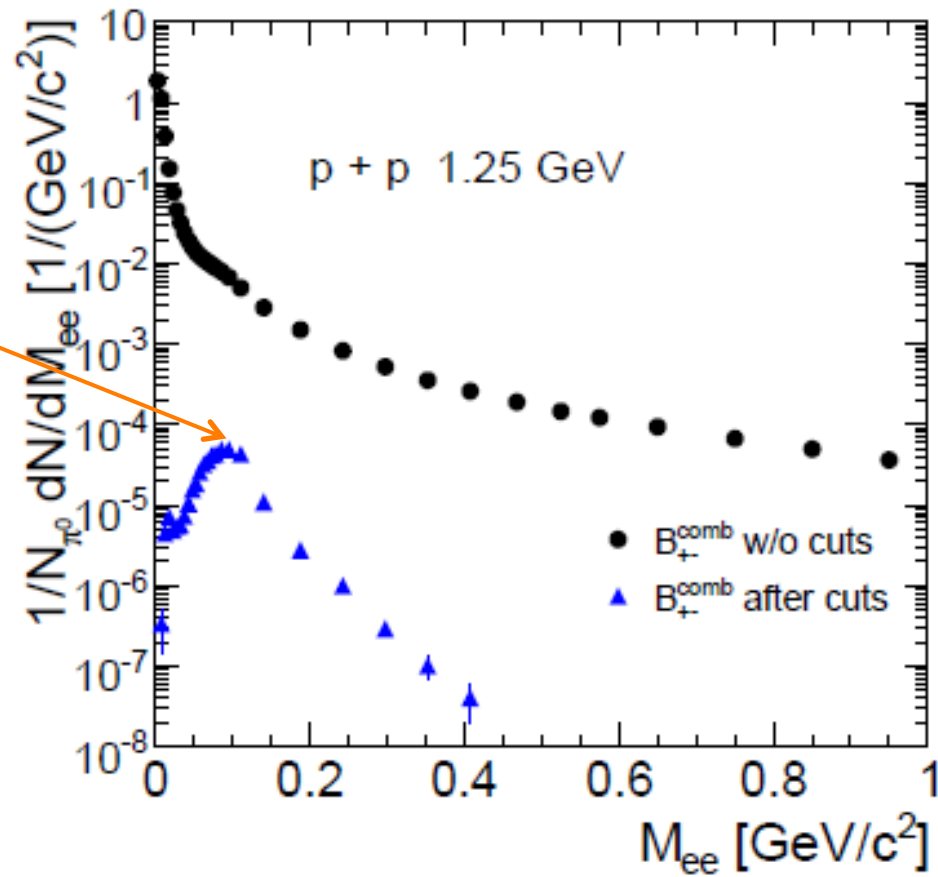


**C3**

" $\pi^0$ -Dalitz rejection":

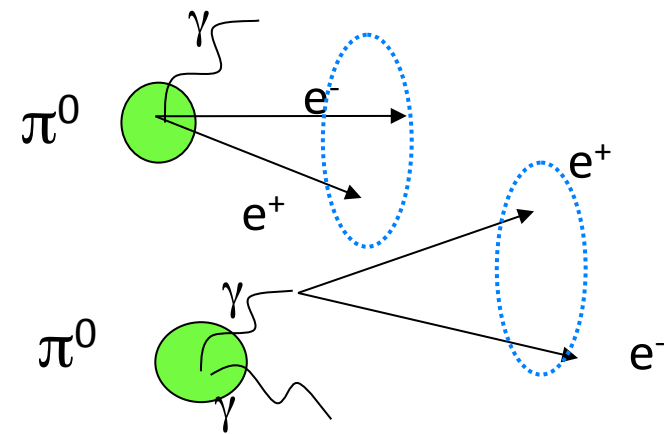
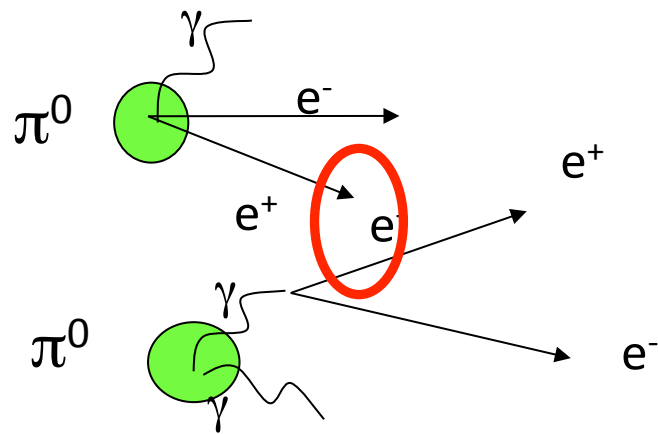
- C3: opening angle  $< 9^\circ$
- Remove both tracks from the sample

## Invariant mass spectra of the CB



# Reconstruction of the combinatorial background

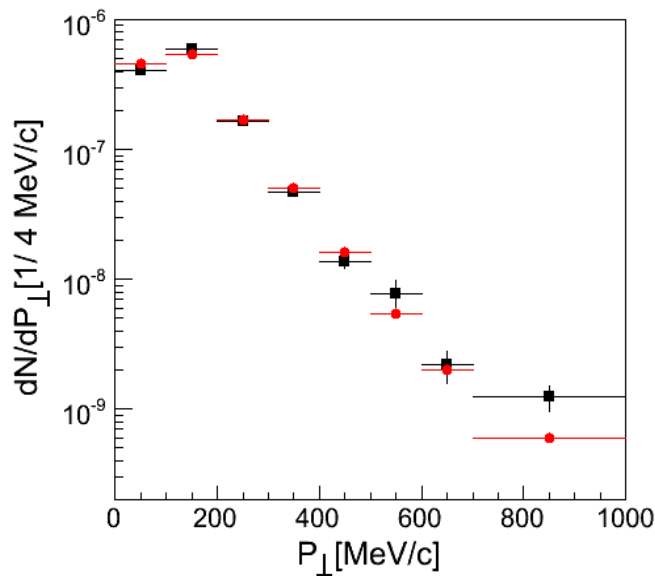
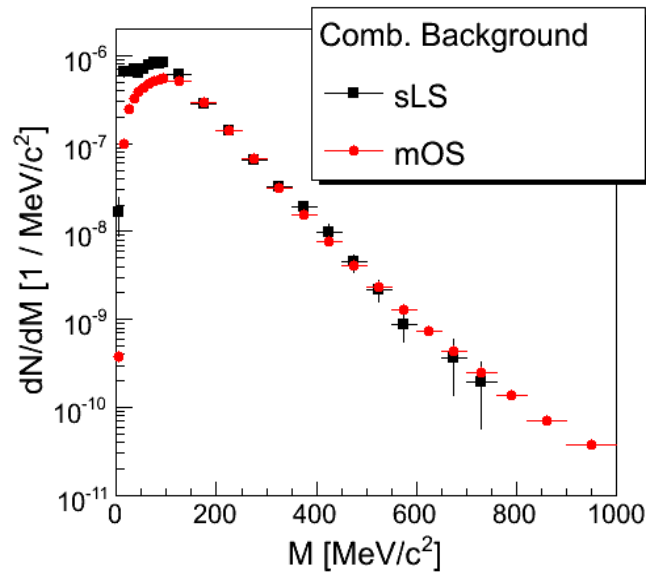
- $N_{e^+e^-}$  - all unlike-sign pairs
- CB – combinatorial background
- $S_{e^+e^-} = N_{e^+e^-} - \text{CB}$



Uncorrelated background  
Event mixing

Correlated background  
Like sign method

# Reconstruction of the combinatorial background



- Same event like-sign:

$$CB_{geom.} = 2 \cdot \sqrt{N_{e^+e^+} \cdot N_{e^-e^-}}$$

$$CB_{arith.} = N_{e^+e^+} + N_{e^-e^-}$$

- Event mixing:

- inherently independent

➔ Normalization done between 150-550 MeV/c<sup>2</sup> M<sub>ee</sub>

➔ sLS and mOS CB show same behavior for M<sub>ee</sub> > 150 MeV/c<sup>2</sup>

➔ For M<sub>ee</sub> < 150 MeV/c<sup>2</sup> deviations due to correlated background

$\pi \rightarrow \gamma\gamma \rightarrow eeX$

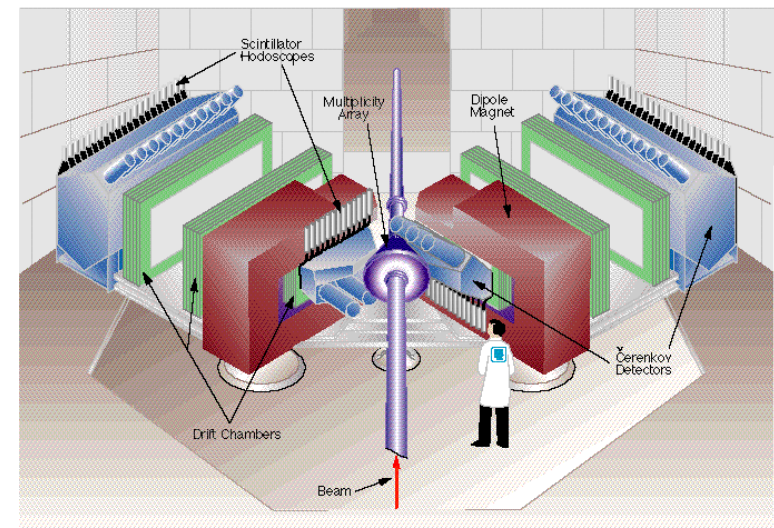


What is known at few GeV regime?

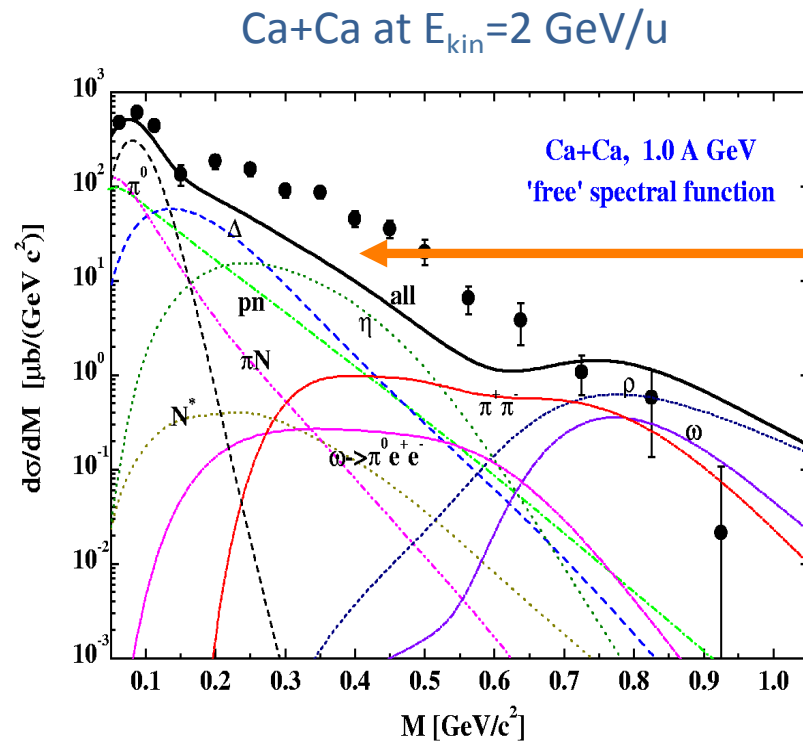
- 1988 – 1993 at Bevalac
- 2 Arm-Spectrometer
  - Minimum opening angle:  $40^\circ$
  - Each arm:  $40^\circ$  in  $\Phi$ ,  $\pm 7.5^\circ$  in  $\Theta$
  - Trigger on electron-pairs
  - Opening angle  $40^\circ$
  - Quasi-tracking:  $p > 0.05$  GeV/c
  - Mass resolution: 15% at  $\omega$  pole mass
  - 30-40% systematical error
- pp/pd, Ca+Ca, C+C



DiLepton Spectrometer



# DLS: enhanced dilepton yields in A+A



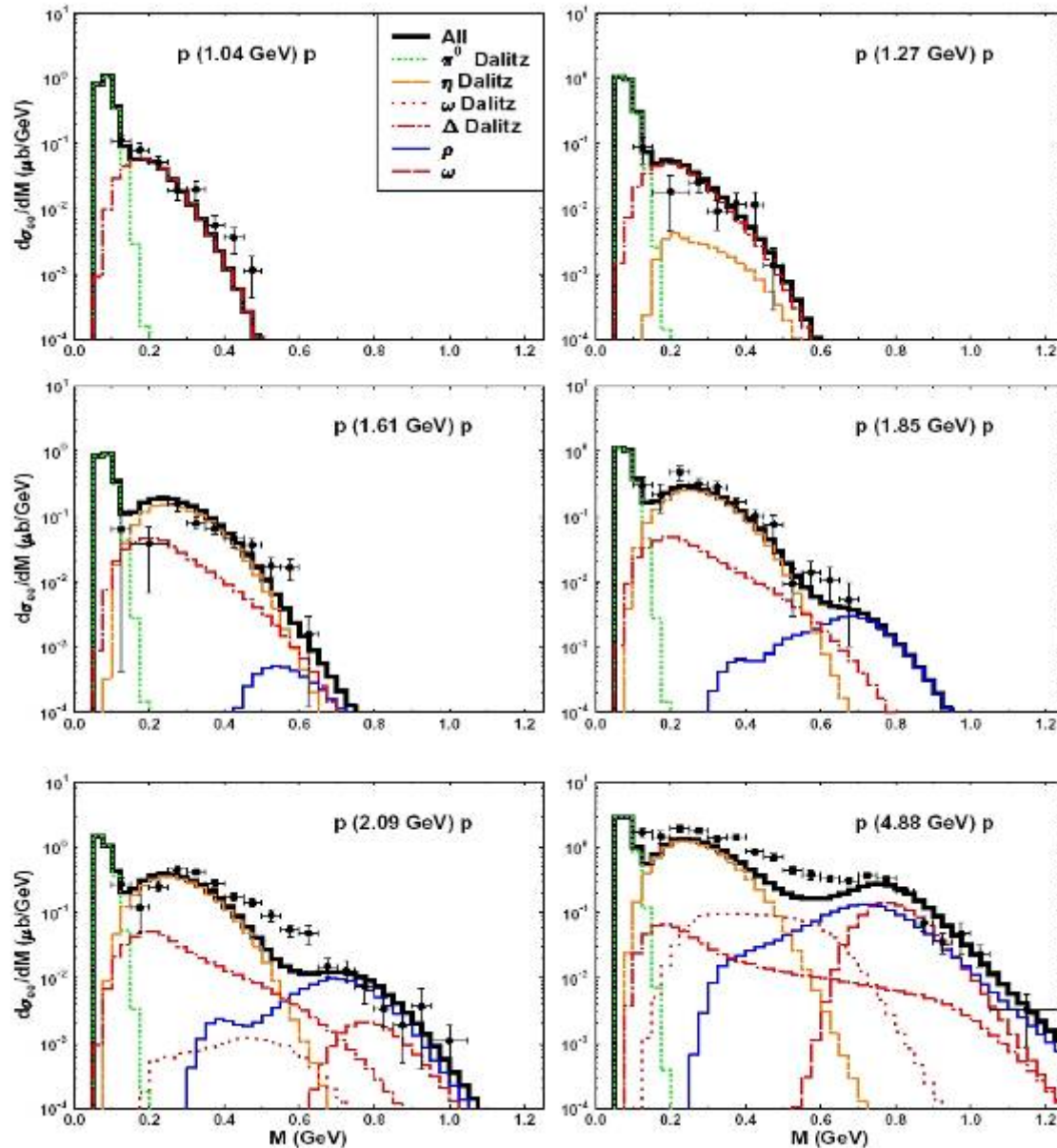
Strong dilepton  
enhancement over  
hadronic cocktails

Data: R.J. Porter et al.: PRL 79(97)1229

Model: E.L. Bratkovskaya et al.: NP A634(98)168, BUU, vacuum  
spectral function



# DLS p+p data vs. model



## Data:

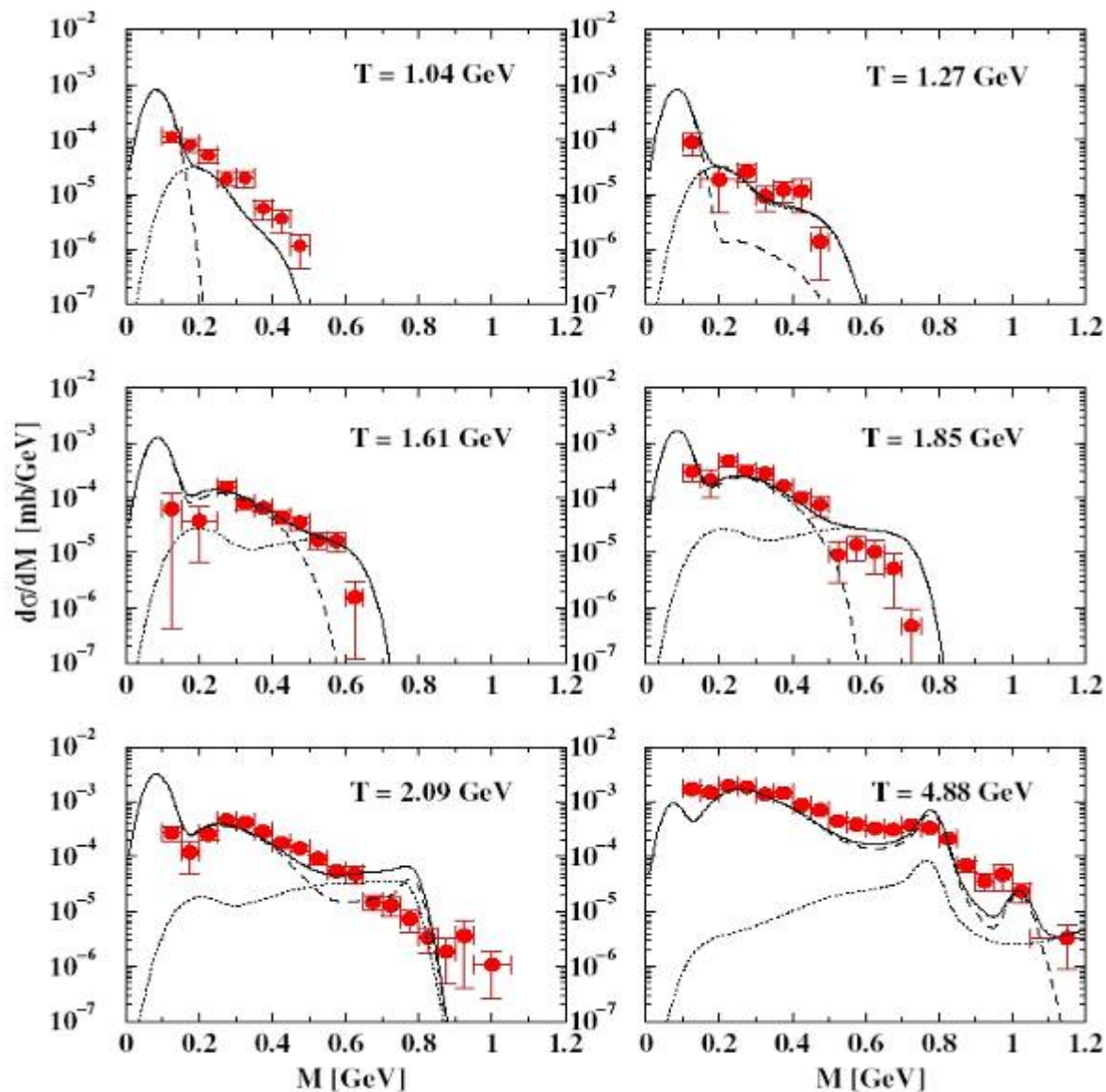
Wilson et al. PRC 57  
(1997) 1865

## Theory (folded with the DLS response):

C. Ernst et al.  
PRC 58 (1998) 447

*UrQMD 1.3*

# DLS $p+p$ data: more and different models ...



## Data:

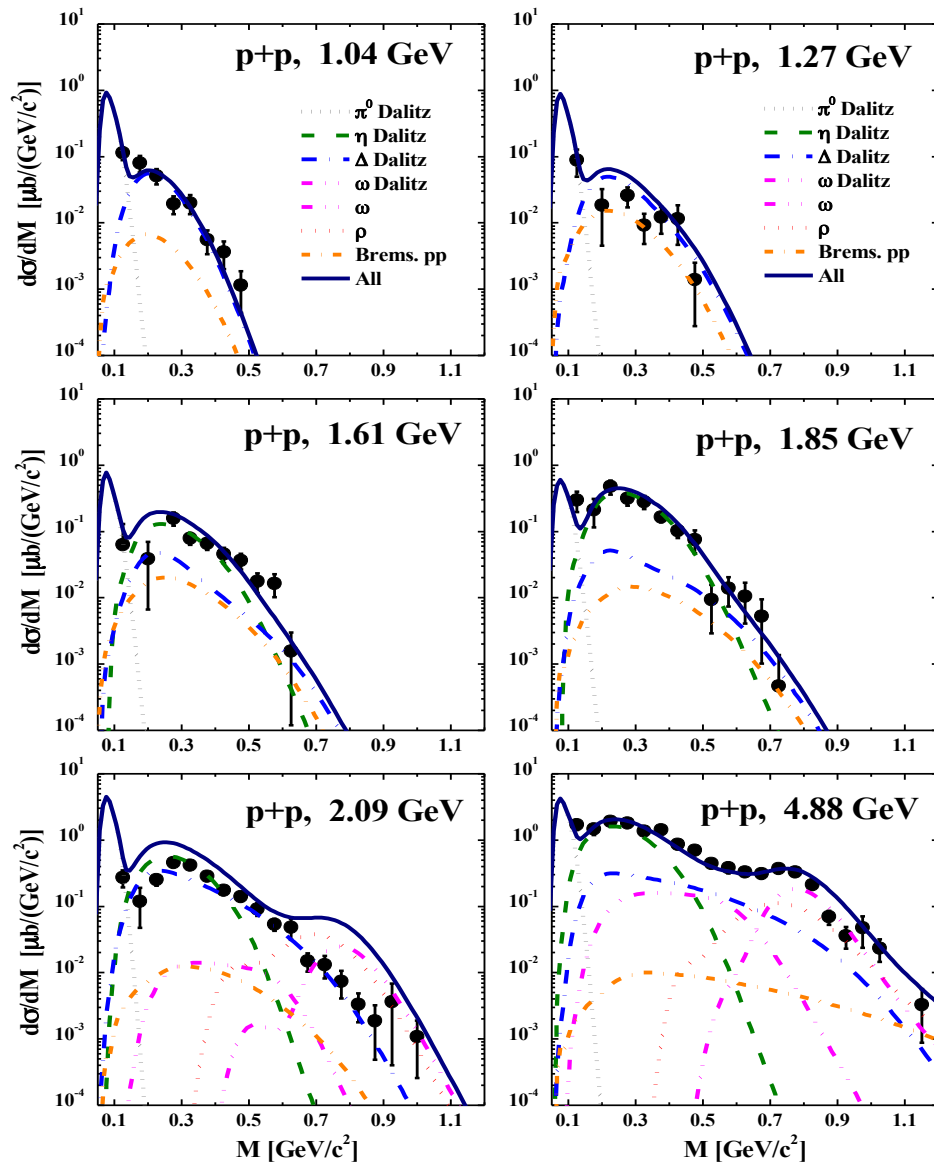
Wilson et al. PRC 57  
(1997) 1865

## Theory (folded with the DLS response):

Faessler, Fuchs et al.  
J. Phys. G29 (2003) 603  
(Resonances + decays)

*RQMD*

# DLS p+p data: ... and more models



## Data:

Wilson et al. PRC 57  
(1997) 1865

## Theory (folded with the DLS response):

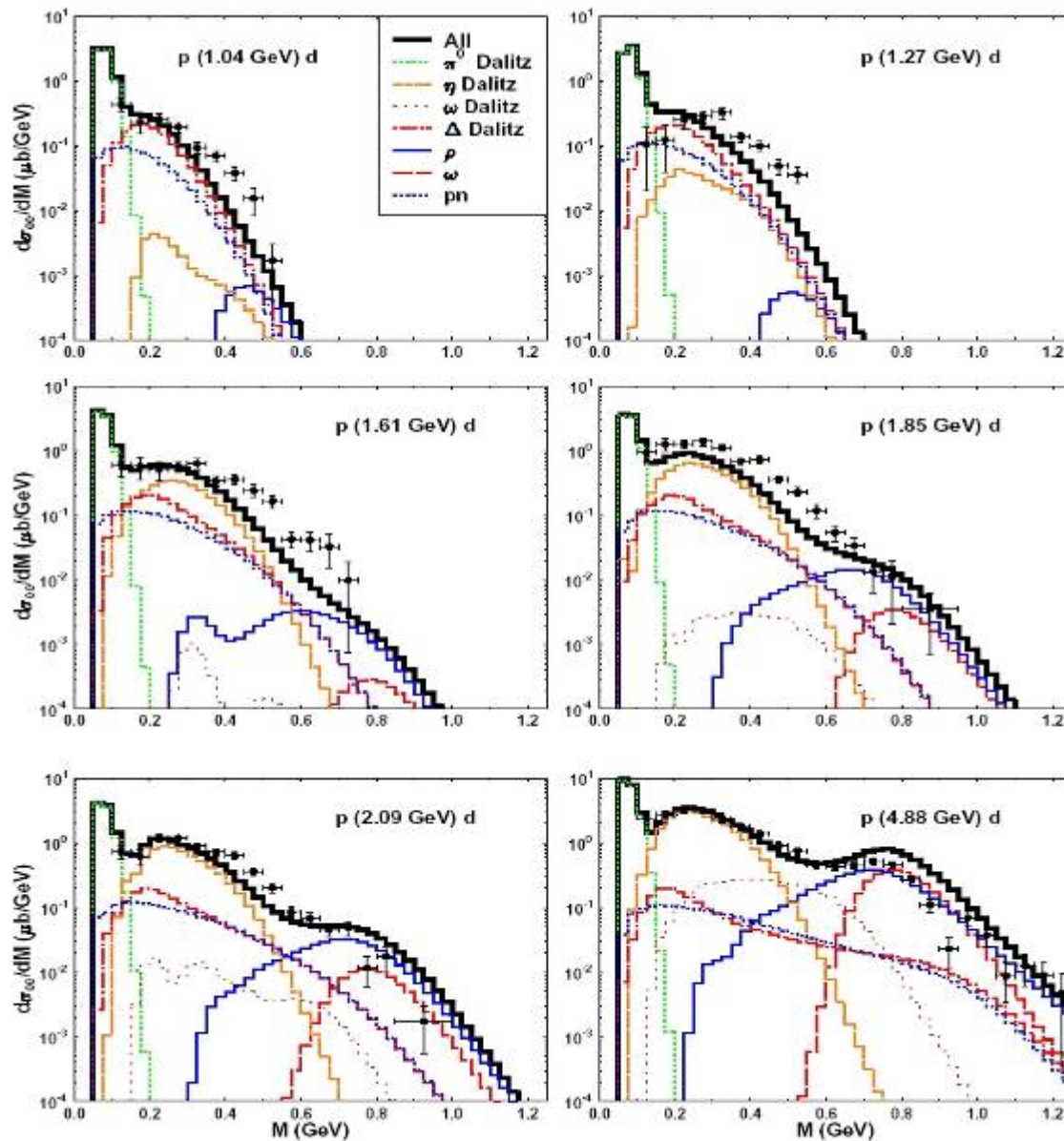
Bratkovskaya et al.,  
HSD model  
(NN Bremstrahlung a-la  
Kaptari *et al.*, 2006)

HSD

DLS p+p data: fair agreement with theory

The real trouble starts with p+d data!

# DLS p+d data vs. model



## Data:

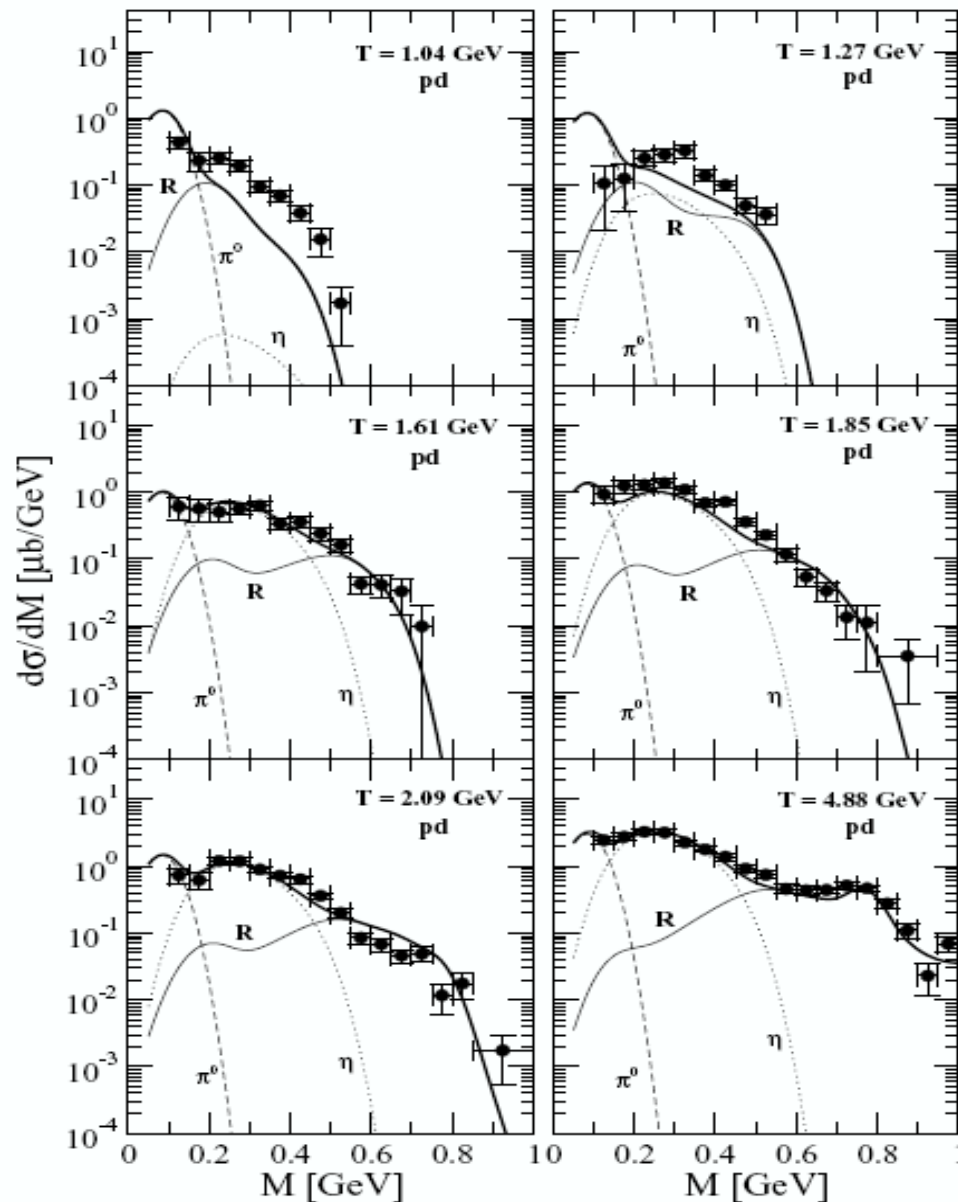
Wilson et al. PRC 57  
(1997) 1865

## Theory (folded with the DLS response):

C. Ernst et al.  
PRC 58 (1998) 447

*UrQMD 1.3*

# DLS p+d data: more and different models ...



## Data:

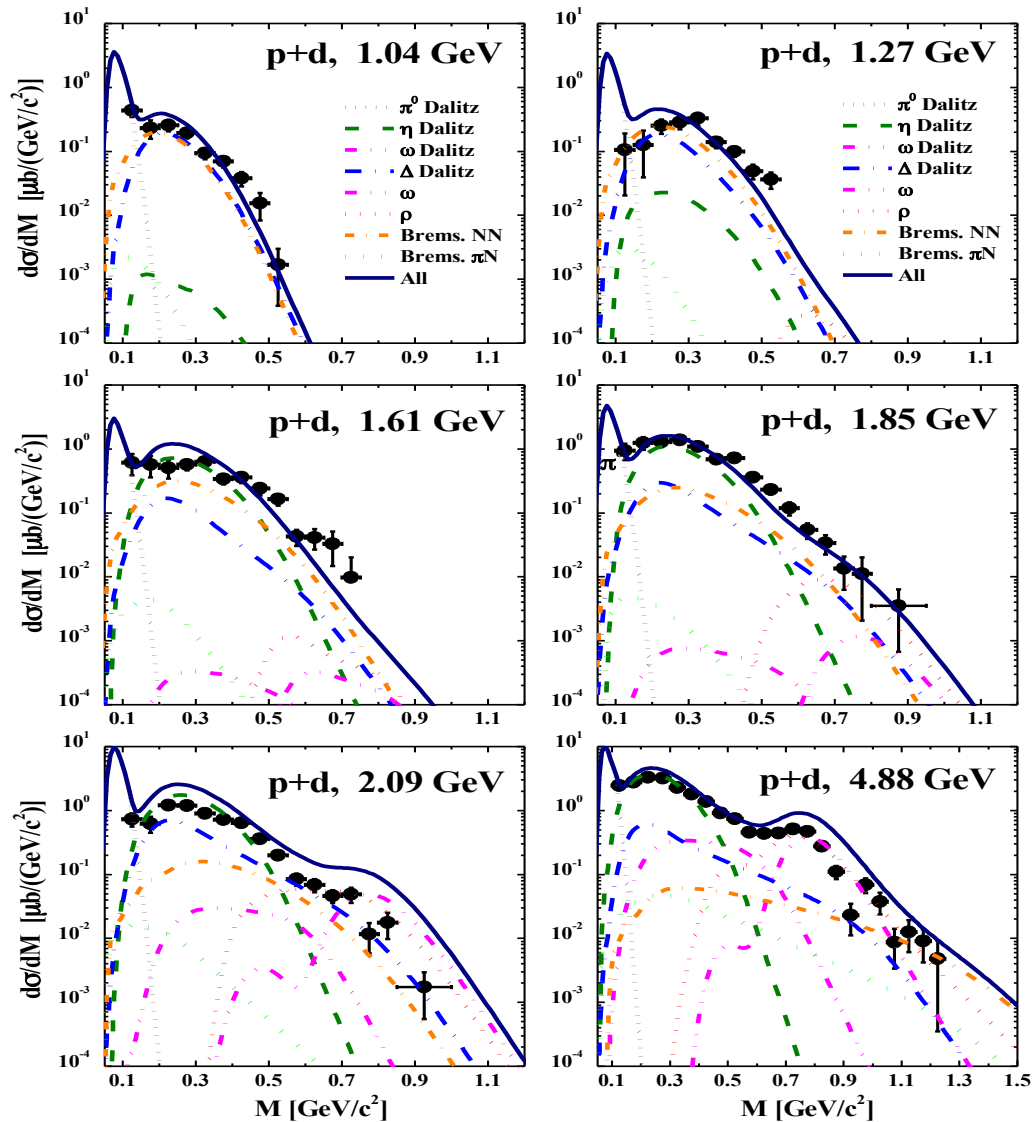
Wilson et al. PRC 57  
(1997) 1865

## Theory (folded with the DLS response):

Faessler, Fuchs et al.  
J. Phys. G29 (2003) 603  
(Resonances + decays)

*RQMD*

# DLS p+d data: ... and more models



## Data:

Wilson et al. PRC 57  
(1997) 1865

## Theory (folded with the DLS response):

Bratkovskaya et al.,  
HSD model  
(NN Bremsstrahlung a-la  
Kaptari *et al.*, 2006)

HSD

DLS p+d data: not described by theory!

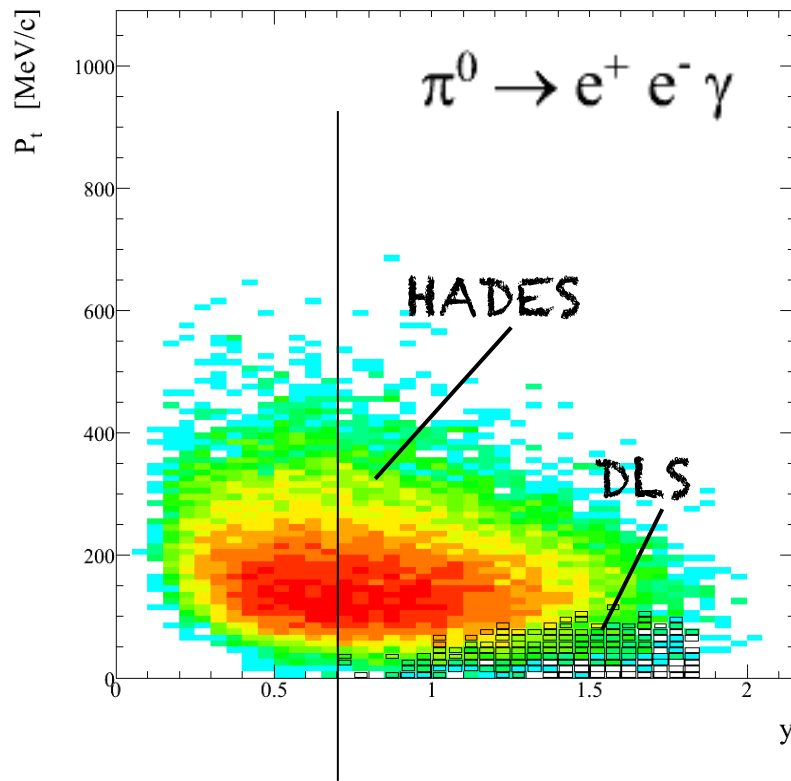
pp vs. pd : What's different?

DLS "pd Puzzle"?

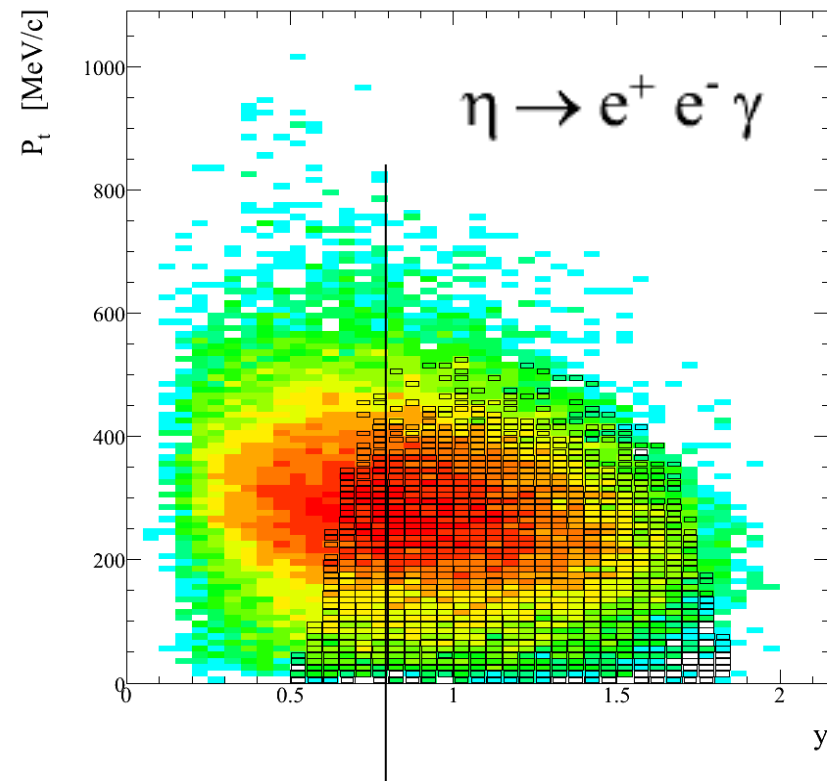


# Validate on DLS: is data correct?

## Phase space coverage: HADES vs. DLS



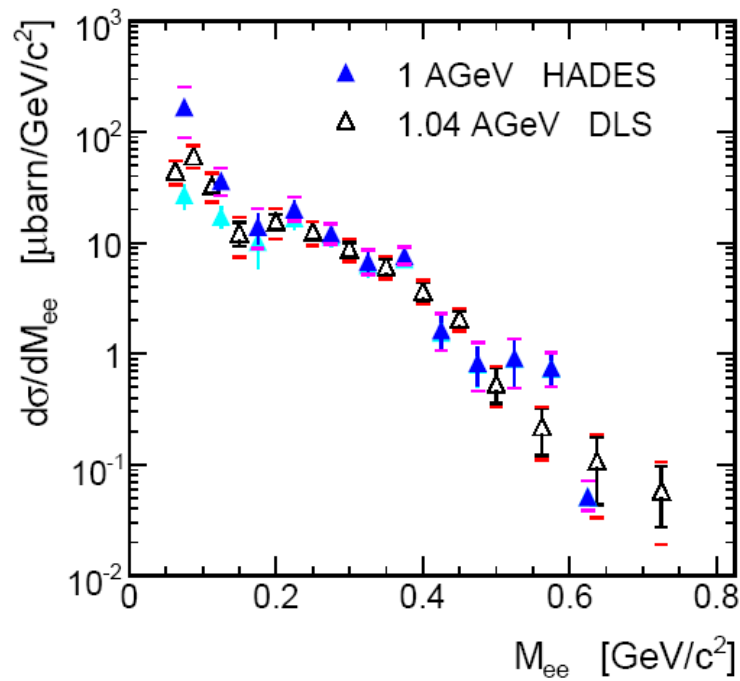
mid-rapidity



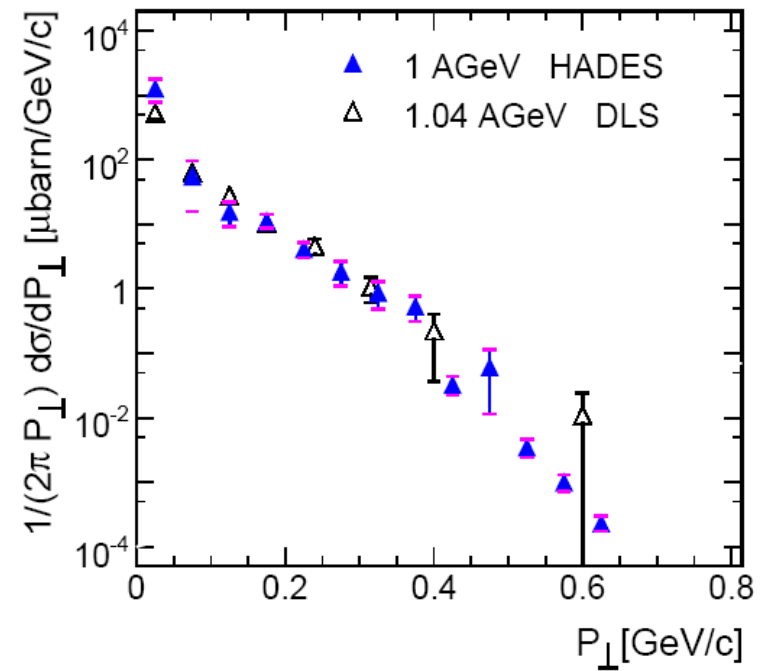
mid-rapidity

For a comparison of HADES and DLS results the HADES yield has to be extrapolated to full phase space

# Direct comparison



DLS Data: R.J. Porter et al.: PRL 79(97)1229



J. Carroll – presentation

International Workshop on Soft Dilepton Production  
August 20-22, 1997, LBNL

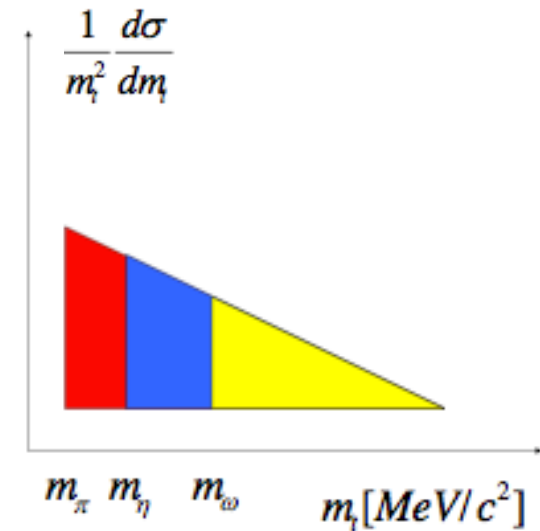
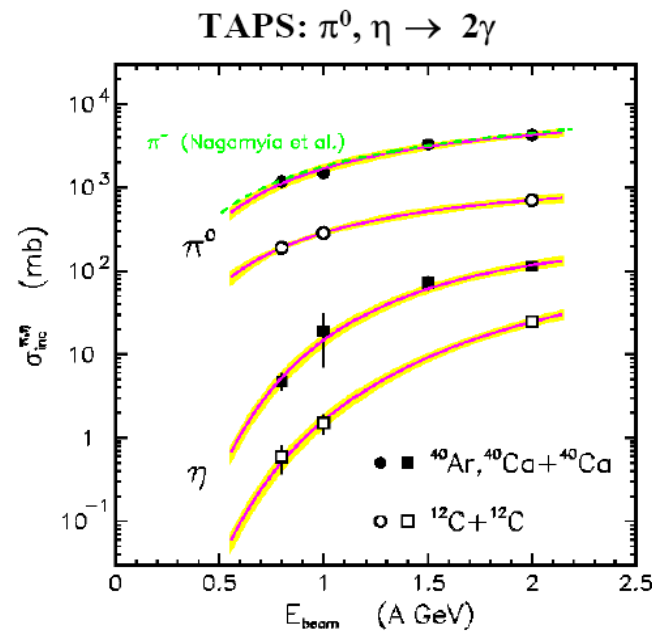
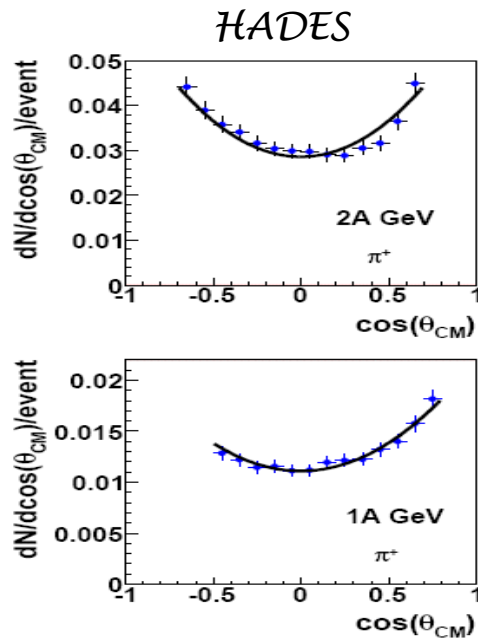
→ HADES and DLS data agree

# Hadronic cocktail

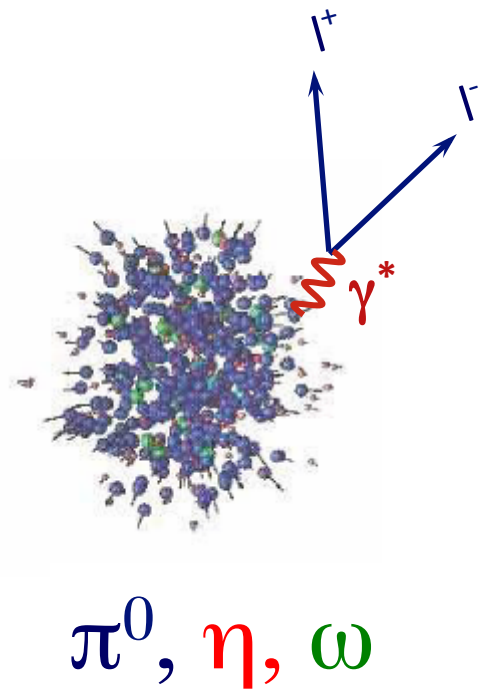
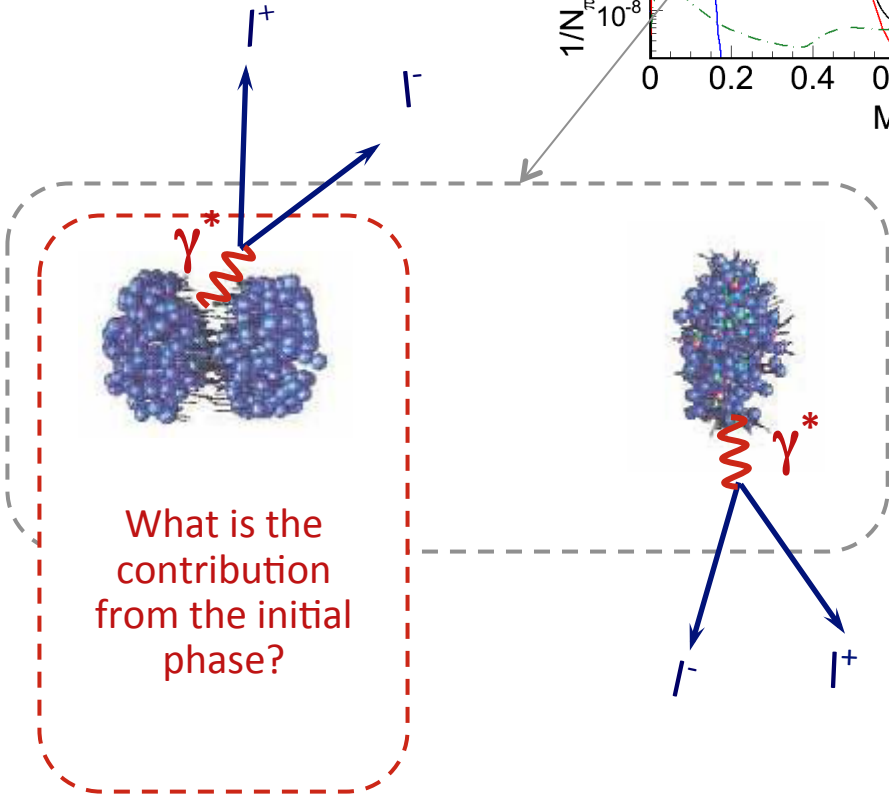
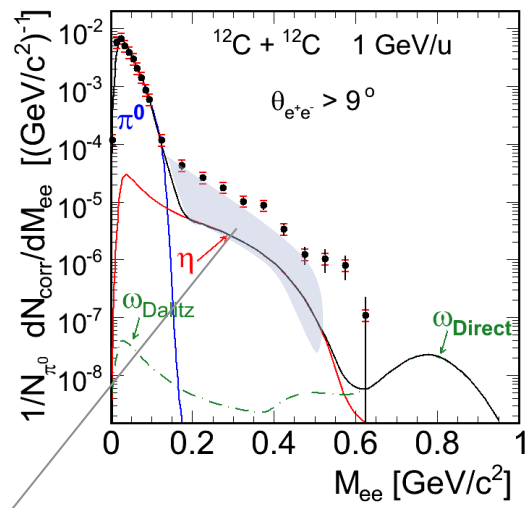


HADES Cocktail = long lived mesonic components

- $\pi^0$  thermal source, anisotropic angular distribution according to measured  $\pi^{+/-}$
- $\eta$  isotropic
- $\omega$   $m_\tau$  scaling isotropic decay pattern



# What is missing?



# Elementary reactions

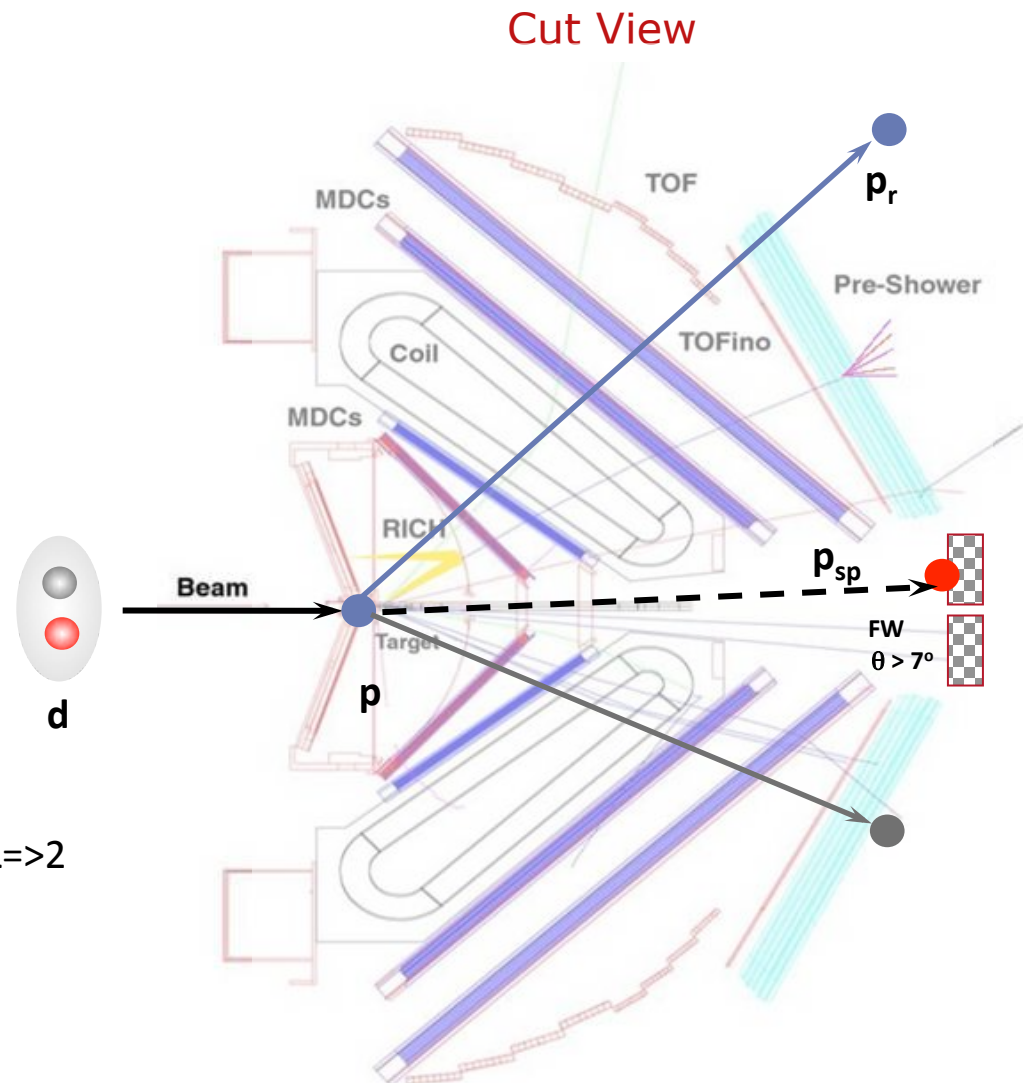
- Beam energy  $E_{\text{beam}} = 1.25 \text{ GeV}$   
( $s < s_{\text{thres}}$  for  $\eta$  production)
- LH2 target

p+p:

- × one week of running in April 2006
- ×  $\sim 2.6 \cdot 10^9$  LVL1 events collected  
(MUL=>3 trigger)

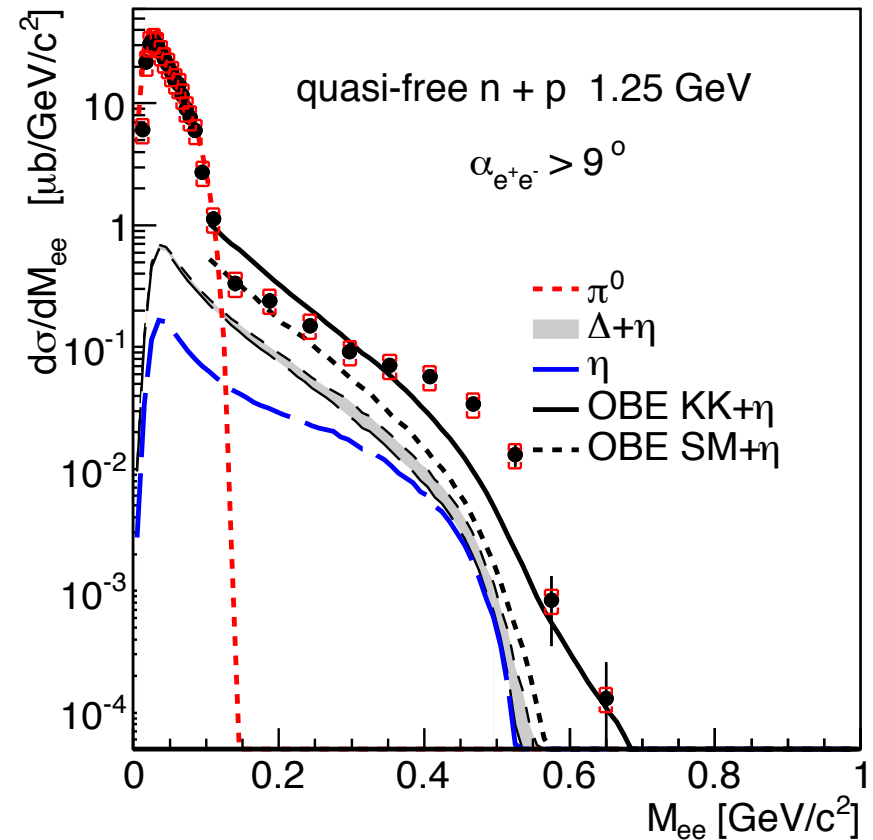
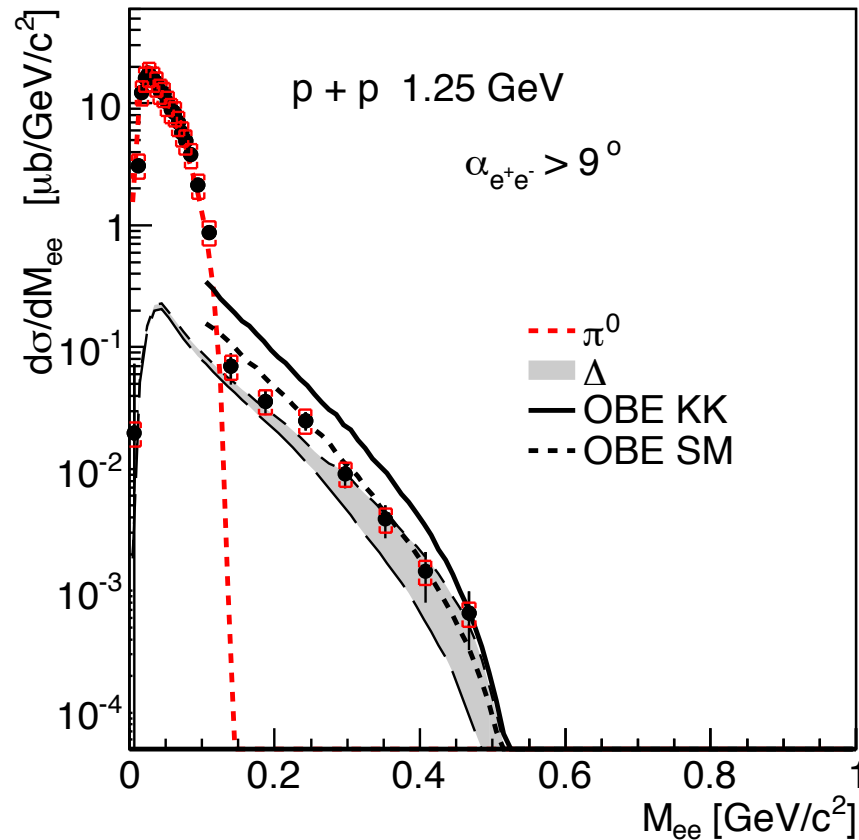
d+p:

- × two weeks of running in April 2007
- ×  $\sim 4.8 \cdot 10^9$  LVL1 events collected (MUL=>2  
&& FW "p spectator")
- tag on np  $\rightarrow e^+e^- X$  reactions



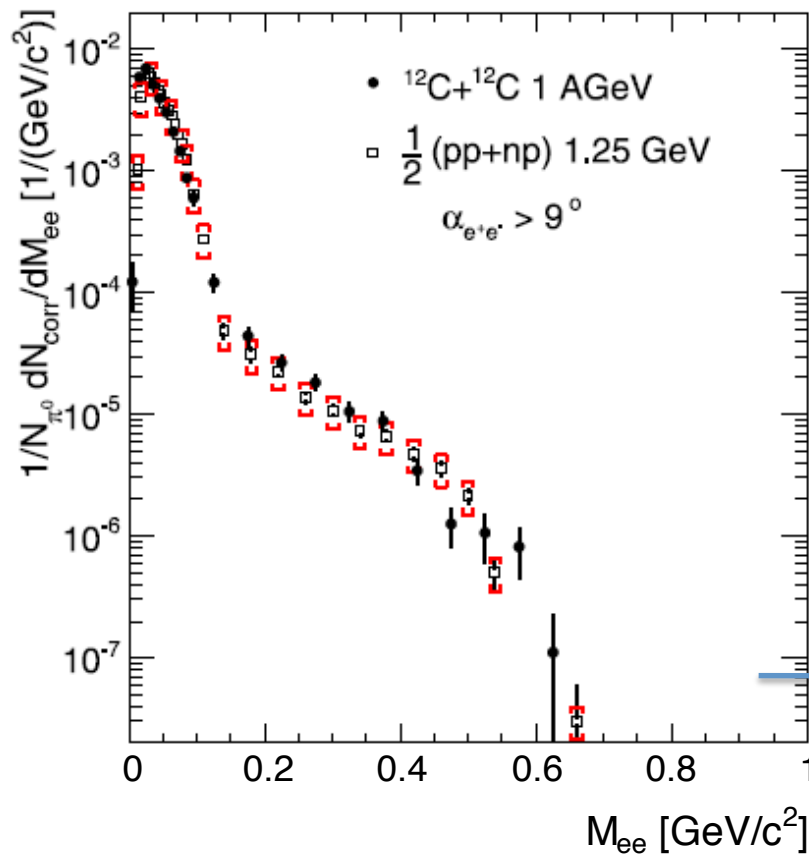
# HADES pp and dp (tagged n) data vs. models

"If you are out to describe the truth, leave elegance to the tailor" + A. Einstein



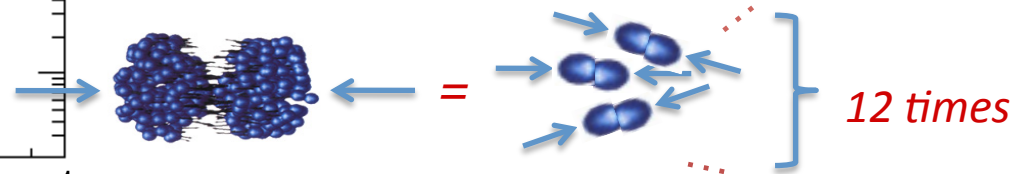
# Origin of the low-mass pair excess in C+C collisions

Comparison of C+C to N+N collisions



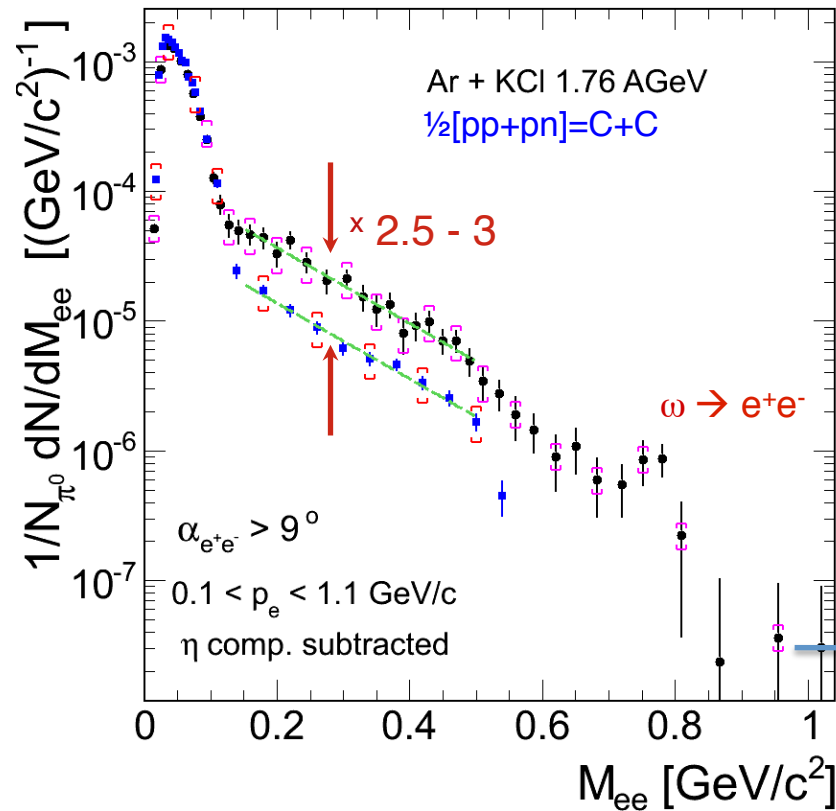
- C+C data reproduced (within 20%) by superposition of N+N interactions
- Pair excess observed in C+C data has been traced back to anomalous pair production in n+p collisions

*C+C = 12 \* (Nukelon+Nukleon)?*



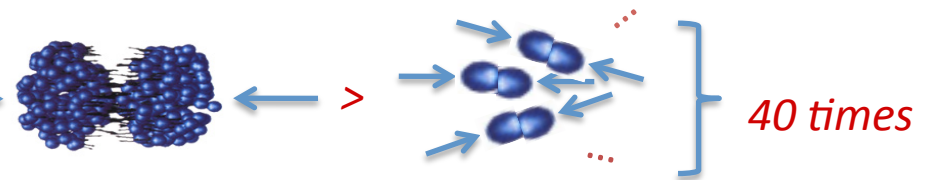
# Is there physics beyond free NN?

Efficiency corrected dielectron spectra  
from Ar+KCl at  $E_{kin} = 1.76$  GeV/u



- First observation of  $w \rightarrow e^+e^-$  peak in heavy-ion collisions at SIS energies
- First evidence for „true“ excess (strong excess above the “reference cocktail”)

*Is Ar+KCl = 40 \* (Nukelon+Nukleon)?*





## Summary: HADES and DLS

- Origin of the low-mass dielectron pair excess in nucleus-nucleus collisions at 1-2 GeV/ $u$  established
- $p+p$  and  $n+p$  data are critical test for theoretical input
- light systems (i.e  $C+C$ ) can be described by superposition of  $NN$  interactions
- “DLS puzzle”?
  - experimentally – solved
  - theoretically – only after  $n+p$  data is consistently explained

LESON: know your reference!

