# Heavy-Quark Energy Loss in the QGP and non-photonic Single-Electron Observables

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#### Outline

#### Heavy quarks in the QGP

Radiative energy loss Collisional energy loss

Dissipation and fluctuation: Fokker-Planck approach

Non-perturbative Effects

# Motivation

- Measured  $p_T$  spectra and  $v_2$  of non-photonic single electrons
- coalescence model describes data under assumption of c quarks flowing with the bulk medium [Greco, Ko, Rapp 04]

# Motivation

- Measured  $p_T$  spectra and  $v_2$  of non-photonic single electrons
- coalescence model describes data under assumption of c quarks flowing with the bulk medium [Greco, Ko, Rapp 04]
- ▶ What is the underlying microscopic mechanism for thermalization?
  - Radiative energy loss
  - +pQCD collisional energy loss
  - elastic three-body pQCD processes
- Additional problem: consistency between  $R_{AA}$  and  $v_2$ 
  - importance of thermal fluctations
  - ► Fokker-Planck approach to HQ rescattering ⇔ thermalization
  - Langevin simulation to include (anisotropic) flow of sQGP
- ▶ non-perturbative processes ⇔ resonances in sQGP

Radiative energy loss Collisional energy loss

# Heavy quarks in the QGP





HQ rescattering in QGP radiative/collisional energy loss non-perturbative effects (sQGP)



Hadronization to D, B mesons Fragmentation Coalescence



Semileptonic decay  $\Rightarrow$  "non-photonic" electron observables

Radiative energy loss Collisional energy loss

# Radiative energy loss

- medium modelled by static scattering centers [GW 94]
  radiative energy loss only!
- $\Delta E \simeq \hat{q}L^2$  [BDMPS 96]
- generalized to "thin plasmas" in [GLV 00] and heavy-quark jets



#### Heavy quarks in the QGP

Dissipation and fluctuation: Fokker-Planck approach Non-perturbative Effects Radiative energy loss Collisional energy loss

## Radiative energy loss



- Calculation: [Armesto et al 06] (static medium + geometry + BDMPS rad energy loss)
- ▶ need to tune up  $\hat{q} \rightarrow 14 \text{ GeV}^2/\text{fm}$  (pQCD prediction:  $\sim 1...3 \text{ GeV}^2/\text{fm}$ )
- ▶ *R*<sub>AA</sub> near to data but *v*<sub>2</sub> not described!

Radiative energy loss Collisional energy loss

#### Collisional vs. radiative energy loss

 for heavy quarks: elastic pQCD scattering as important as radiative [Mustafa 05]



► calculation [Djordjevic '06]: t-channel gluon exchange dressed gluon propagator  $\mu_D^2 = g^2 T^2 (1 + N_f/6), \ \alpha_s = 0.3, \ N_f = 2.5$  $dN_g/dy = 1000$ 

collisional energy loss important for light and heavy quarks!

Radiative energy loss Collisional energy loss

#### Collisional vs. radiative energy loss



Radiative energy loss Collisional energy loss

#### Three-body effects

- high densities (initially  $\gtrsim 10/{\rm fm}^3$ )
- $\Rightarrow$  three-body elastic scattering possibly relevant [Liu, Ko 06]



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- ⇒ Use Fokker-Planck equation [Svetitsky 87; Mustafa, Thoma 98; HvH, Rapp 04; Moore, Teaney 04,...] ⇔ Langevin simulations
  - can we understand heavy-quark flow properties better?
  - consistency of  $e^{\pm}$ - $R_{AA}$  with  $e^{\pm}$ - $v_2$ ?

#### The Fokker-Planck Equation

heavy particle (c,b quarks) in a heat bath of light particles (QGP)

$$\frac{\partial f(t,\vec{p})}{\partial t} = \frac{\partial}{\partial p_i} \left[ p_i A(t,p) + \frac{\partial}{\partial p_j} B_{ij}(t,\vec{p}) \right] f(t,\vec{p})$$

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- ►  $A(t, \vec{p})$  friction (drag) coefficient =  $1/\tau_{eq}$  $\langle p_i - p'_i \rangle = p_i A(t, \vec{p})$

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► to ensure correct equilibrium limit: B<sub>||</sub>(t, p) = T(t)E<sub>p</sub>A(t, p) (Einstein dissipation-fluctuation relation)

#### Langevin Study with pQCD elastic scattering

 pQCD elastic cross sections for charm-quark scattering in QGP [Moore, Teaney 04]



- hydro dynamics for bulk medium
- Langevin simulation for charm quarks
- ▶ have to increase  $\alpha_s$  in cross sections (but set  $\mu_D = 1.5 T = \text{const!}$ )

#### Non-perturbative Effects

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- ▶ from Lattice QCD: survival of mesonic bound states/resonances above T<sub>c</sub> [Karsch, Laermann 03], [Asakawa, Hatsuda 03]
- also from IQCD based potential models [Shuryak, Zahed 04], [Wong 05], [Mannarelli, Rapp 05]

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- $\Rightarrow$  assumption:

survival of D- and B-like resonance states up to  $T\lesssim 2T_c$ 

- here: use "quasi-particle" model based on chiral symmetry and heavy-quark effective theory
- ▶ states included: D, D\*+chiral partners, D<sub>s</sub> (analogous for B) [HvH, Ralf Rapp, Phys. Rev. C **71**, 034907 (2005)]

## **Resonance Scattering**

elastic heavy-light-(anti-)quark scattering



▶ *D*- and *B*-meson like resonances in sQGP



parameters

- $m_c = 1.5 \text{ GeV}, m_D = 2 \text{ GeV}, \Gamma_D = 0.4 \dots 0.75 \text{ GeV}$
- $m_b = 4.5 \text{ GeV}, m_B = 5 \text{ GeV}, \Gamma_B = 0.4 \dots 0.75 \text{ GeV}$
- Bethe-Salpeter calculations in NJL model [Blaschke et al 03]

# Contributions from pQCD

Lowest-order matrix elements [Combridge 79]



► In-medium Debye-screening mass for *t*-channel gluon exchange:  $\mu_g = gT$ ,  $\alpha_s = 0.4$ 

#### Cross sections



- total pQCD and resonance cross sections: comparable in size
- ► BUT pQCD forward peaked ↔ resonance isotropic
- resonance scattering more effective for friction and diffusion

### The Coefficients: pQCD vs. resonance scattering

 Temperature dependence of thermalization rate

- charm-quark diffusion coefficient
- microscopic properties of sQGP  $\Leftrightarrow e^{\pm}$  observables



## Initial conditions

#### Langevin simulation:

need initial  $p_T$ -spectra of charm and bottom quarks

- fit D-meson spectra from pp and dAu@RHIC
- exp. non-photonic single- $e^{\pm}$  spectra: Fix bottom/charm ratio



#### Spectra and elliptic flow for heavy quarks

- use Langevin simulation to solve Fokker-Planck equation
- expanding-fireball model to describe the sQGP medium



• 
$$\mu_D = gT$$
,  $\alpha_s = g^2/(4\pi) = 0.4$ 

- ▶ resonances ⇒ HQ thermalization without upscaling of cross sections
- Fireball parametrization consistent with hydro

# Observables: $p_T$ -spectra ( $R_{AA}$ ), $v_2$

- Hadronization: Coalescence with light quarks + fragmentation  $\Leftrightarrow c\bar{c}, b\bar{b}$  conserved
- ▶ single electrons from decay of *D* and *B*-mesons



 Without further adjustments: data quite well described [HvH, V. Greco, R. Rapp, Phys. Rev. C 73, 034913 (2006)]

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- Hadronization: Fragmentation only
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# Observables: $p_T$ -spectra ( $R_{AA}$ ), $v_2$

- Central Collisions
- ▶ single electrons from decay of *D* and *B*-mesons

#### ${\sf Coalescence}{+}{\sf Fragmentation}$

Fragmentation only



#### How to check resonance assumption?

- scattering mechanism via resonances at  $T > T_c$ ?
- dominant channel: quark-anti-c-quark s channel



energy scan@RHIC: quark dominated ⇒ c̄ quarks most affected
 thermalization effects more pronounced for D̄ (D<sup>-</sup>) than for D (D<sup>+</sup>) mesons!

#### Implementation of radiative energy loss

#### including gluon radiation work in progress [Vitev, HvH, Rapp 06]



# Conclusions and Outlook

- ▶ non-photonic  $e^{\pm}$  observables  $\Leftrightarrow$  HQ interactions in sQGP
- HQ energy loss from pQCD
  - ▶ radiative energy loss  $\Leftrightarrow$  upscaling of energy loss  $\hat{q} \rightarrow 14$  or gluon density to explain strong effects in  $e^{\pm}$ - $R_{AA}$
  - collisional (elastic) energy loss
  - ▶ high density of plasma ⇔ elastic 3-body collisions
- proper implementation of thermalization (Fokker-Planck Eq.)
  - need thermal fluctuations to describe thermalization
  - explains consistency between small  $R_{AA}$  and large  $v_2$
- non-perturbative interactions
  - survival of D- and B-meson like resonances above  $T_c$
  - ► isotropic elastic-scattering cross sections ⇒ efficient for thermalization

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  - ► isotropic elastic-scattering cross sections ⇒ efficient for thermalization
- Further investigations (work in progress)
  - microscopic models for HQ scattering [Mannarelli, HvH, Rapp 06]
  - implementation of gluon-radiation processes [Vitev, HvH, Rapp 06]
  - consequences for heavy quarkonia

### Thermalization rate (p dependence)



#### Spectra and elliptic flow for heavy quarks

#### With form-factor vertices instead of point vertices ( $\Lambda = 1 \text{ GeV}$ )

